

Aerodynamics and motor control of ultrasonic vocalizations for social communication in mice and rats

Jonas Håkansson¹, Weili Jiang², Qian Xue², Xudong Zheng², Ming Ding^{3,4}, Anurag A. Agarwal⁵, Coen P.H. Elemans^{1,*}

¹ Department of Biology, University of Southern Denmark, 5230 Odense M, Denmark.

² Department of Mechanical Engineering, University of Maine, Orono, ME, 04469, USA.

³ Department of Orthopaedic Surgery and Traumatology, Odense University Hospital, 5000 Odense C, Denmark

⁴ Department of Clinical Research, University of Southern Denmark, 5000 Odense C, Denmark

⁵ Department of Engineering, University of Cambridge, UK

* **Corresponding author:** coen@biology.sdu.dk

ORCID: Jonas Håkansson, 0000-0002-1247-5586; Weili Jiang, 0000-0002-7576-9653; Qian Xue, 0000-0002-8462-1271; Xudong Zheng, 0000-0001-5403-8862; Ming Ding, 0000-0002-6610-8079; Anurag Agarwal, 0000-0002-0332-2166; Coen Elemans, 0000-0001-6306-5715.

Keywords: bioacoustics, vocal production, acoustic communication, speech, rodents

Abstract: Rodent ultrasonic vocalizations (USVs) are crucial to their social communication and a widely used translational tool for linking gene mutations to behavior. To maximize the causal interpretation of experimental treatments, we need to understand how neural control affects USV production. However, both the aerodynamics of USV production and its neural control remain poorly understood. Here we test three intralaryngeal whistle mechanisms - the wall and alar edge impingement, and shallow cavity tone - by combining in vitro larynx physiology and individual-based 3D airway reconstructions with fluid dynamics simulations. Our results show that in the mouse and rat larynx USVs are produced by a glottal jet impinging on the thyroid inner wall. Furthermore, we implemented an empirically based motor control model that predicts motor gesture trajectories of USV call types. Our work provides a quantitative neuromechanical framework to evaluate the contributions of brain and body in shaping USVs, and a first step in linking descending motor control to USV production.

34 **Introduction**

Murine rodents produce ultrasonic vocalizations (USVs) that range in frequencies from 20 to 100
36 kHz and play a crucial role in social communication behaviors, such as mating and territorial
defense (1, 2). Different USV call types strongly signal positive (3) or negative (4) emotional
38 states (5, 6) and are crucial for pups to induce maternal search and retrieval behavior, when
visual or olfactory cues are less relevant (7). USVs have been found in at least 50 rodent species
40 (8), but are probably more widespread, given that rodents comprise over 40% of all mammal
species (9) and only a fraction has been investigated (8). Furthermore, USVs have recently
42 become an increasingly used behavioral readout in mice and rats, the two most widespread
translational animal disease models in biological and medical research (10). USVs are used a
44 translational tool for linking gene mutations to behavioral changes in rodent models for speech
(11) and neuropsychiatric communication disorders, such as autism (12, 13) and Down syndrome
46 (14). The observed changes in vocalization behavior, such as altered USV occurrence (15),
sound frequency (16, 17) or aberrant USV call types (18), are attributed to changes in neural
48 control (15–18). However, linking the brain to behavior requires a causal and quantitative
understanding of the transformation from descending motor control to USV production in these
50 species that we currently lack.

Translating motor control to USV production requires both system identification of the mechanism
52 by which sound is produced and quantitative understanding of how muscles drive the control
parameters of this system. Until recently, USVs were thought to be hole-tone whistles that require
54 two orifices for producing a stable tone (19, 20), such as the human teeth-lip whistle and tea-
kettle whistle (21). However, USVs in mice were recently shown to be produced by a sound
56 production mechanism novel to mammals and previously only identified in industrial supersonic
and high-speed subsonic flows (22–24): a glottal jet impinging on a structure within the larynx
58 (25). Small instabilities in a glottal air jet that travel downstream are entrained to occur at certain
frequencies due to a feedback loop between these downstream-travelling flow structures and
60 acoustic waves travelling upstream. In the small murine larynx, where glottal jet speeds can reach
up to 10% of the speed of sound, the jet impingement mechanism can lead to stable high-

62 frequency tones from 20 to over 100 kHz (20, 25, 26). The impingement structure within the
larynx has been proposed to be either the thyroid wall (25) or a laryngeal adaptation (27) found in
64 several muroid rodents, the alar edge (27, 28) (Fig. 1). Both mechanisms constrain motor control
to the respiratory and laryngeal musculature, but the proposed aeroacoustic models for wall and
66 alar edge tones occur under distinct physiological conditions and predict very different sound
frequencies (25, 27). Thus establishing which aerodynamic mechanism is responsible for USV
68 production is critical for quantitatively linking neuromuscular control parameters to USV acoustics.

70 Here we test what aerodynamic mechanism explains USV production in rats and mice. We exploit
different predictions made by the main two proposed mechanisms – wall and alar edge
72 impingement tones – and furthermore introduce a novel mechanism, the shallow cavity tone, that
we propose as a more likely aerodynamic flow scenario than an alar edge tone. We combine a
74 series of *in vitro* excised larynx experiments with computational flow models to test three
distinguishing key physiological boundary conditions. We show that USVs are produced with
76 adducted vocal folds, that only the wall impingement model predicts anatomically correct glottal
air jet parameters, and that normal USVs are produced in absence of the alar edge and ventral
78 pouch. Together, all datasets strongly support the intralaryngeal wall impingement mechanism.
We then propose a quantitative motor control model that derives time-resolved control
80 parameters from *in vivo* USV sound recordings and provides a physiological basis for USV
syllable categorization and interpreting rat and mouse vocal behavior phenotypes. Our model
82 furthermore shows that both brain and body contribute to USV frequency traces which
emphasizes the importance of an embodied or systems approach to USV motor control.

84

Results and Discussion

86 We tested three physiological boundary conditions that are distinctive between wall and alar edge
tone models for USV production: *i*) vocal fold adduction state, *ii*) jet separation and impingement
88 locations, and *iii*) the presence of the alar edge and ventral pouch cavity.

90 The first distinctive feature between wall and alar edge tones is vocal fold adduction state (Fig.
1B). In mice, USVs are produced *in vitro* with fully adducted (opposed) vocal folds, which leaves
92 a glottal opening on the dorsal side between the arytenoid cartilages, i.e. the cartilaginous glottis,
for respiratory flow to go through (25). In contrast, the alar edge tone model predicts tones to
94 occur with vocal folds abducted (open), resulting in a much larger glottal opening that includes the
ventral opening between the vocal folds (i.e. the membranous glottis) plus the cartilaginous glottis
96 (27).

98 To test which vocal fold adduction state leads to USV production in rats, we used an excised
larynx paradigm that allowed detailed manipulation of glottal configuration (25, 29) (*Methods and*
100 *Materials, Additional File 1*). We subjected rat larynges to pressure ramps with abducted and
adducted vocal folds. With adducted vocal folds all larynges produced fictive USVs (fUSVs)
102 (N = 10), while only 1 out of 10 produced fUSVs with abducted vocal folds. fUSVs were produced
over a phonation threshold pressure of 0.8 ± 0.3 kPa (N=10), consistent with *in vivo* values of 0.4-
104 0.9 kPa (30). Flow ranged from 2.6 ± 0.6 – 3.7 ± 1.0 ml/s (N=10), which is within estimated
physiological range of 0-10 ml/s (*Methods and Materials*). Furthermore, the fUSVs peak
106 frequencies ranged from 25-61 kHz, which corresponds well to the *in vivo* range of 18-96 kHz,
including “22 kHz” (range 18-32 kHz) and “50 kHz” (range: 32 - 96 kHz) USVs (31). Thus, driving
108 excised rat larynges with physiologically realistic airflows cause fUSVs that overlap in acoustic
parameters with *in vivo* USVs, which suggest that the *in vitro* paradigm represents the *in vivo*
110 situation very well. Our data supports the hypothesis that USVs in rats are produced with
adducted vocal folds, which is consistent with a reduced airflow during USV production compared
112 to quiet respiration in rats (30, 32), earlier *in vitro* glottal adduction manipulations that lacked
sound recordings (33) and preliminary *in vivo* endoscopic observations (34). Thus, USVs in both
114 mice (25) and rats are produced with adducted vocal fold, which provides evidence against the
alar edge tone mechanism and in favor of the wall tone mechanism.

116

The second distinctive feature between wall and alar edge models is the speed, position, length,

118 and angle of a formed air jet. The wall tone model predicts jet formation at the center of the
cartilaginous glottis and impingement on the thyroid inner wall (Fig. 1B) (25). The alar edge tone
120 model predicts that "the glottal jet exits close to the ventral side of the laryngeal lumen, resulting
in a glottal jet path nearly parallel to the intralaryngeal supraglottal wall" (27). Thus, the required
122 jet is proposed to separate on the ventral side of the laryngeal lumen (at Flow Separation Point,
FSP, Fig. 1B), which implies that the jet center is located at the center of the glottis (Fig. 1B). Jet
124 impingement is constrained to the alar-edge (27). These jet location differences thus result in
different jet angles and lengths, which in turn lead to different flow-frequency transformations
126 (*Methods and Materials*). However, we think the proposed alar edge model poses an unlikely flow
scenario for the formation of a separated jet - essential to the edge tone - because the large
128 glottis leads to low flow speeds and a low flow constriction ratio. We also question the validity of
the assumption that the pouch can act as a Helmholtz resonator (27), because the anatomical
130 structure to act as the essential neck is not present. Instead, we propose a third USV production
mechanism, the shallow cavity tone, which is based on a more realistic flow scenario that does
132 not require jet formation, has FSP at the same location as the alar edge model, and leads to
stable high-frequency whistles (35). Cavity flows are produced when air flow detaches flows over
134 a cavity and reattaches downstream of the cavity (at the thyroid in Fig. 1B) and sets up a
recirculating flow inside the cavity. The flow can produce loud tonal sounds. Such flows are of
136 significant interest in aerospace applications, such as wheel wells and weapon bays of aircraft,
where the strong oscillations from the tones can lead to significant structural damage (36).

138
To estimate flow and jet conditions, we combined USV production under controlled *in vitro*
140 conditions with morphometric analysis of individual-based dice-CT scans. In all models,
frequency is set by u , the mean convection speed of the coherent flow structures, approximated
142 as the glottal exit speed, and jet or cavity length (*Methods and Materials*). While the cavity tone
model does not predict the formation of a jet, it does rely on the length of the entrance to the
144 ventral pouch and thereby, for a given frequency, also predicts a length. We measured tracheal
airflow (V) and peak frequency (f_p) during fUSV production (Fig. 3) in fresh larynges (N=5) that

146 were subsequently fixed in PFA to stabilize the geometry. Even after PFA fixation, fUSVs were
produced in all larynges and the slope of the frequency-to-flow relationship did not differ
148 significantly before and after fixation (Paired t-test, 2-tailed, $p = 0.48$). We subsequently
measured the glottal area (A_{gl}) in Dice- μ CT scans (Fig. 3A, *Methods and Materials*) of all
150 individuals to estimate jet exit speed u . The produced frequencies and jet speed predicted jet
lengths of 0.92 ± 0.21 mm, 0.46 ± 0.11 mm, and 0.46 ± 0.13 mm for the wall, alar edge, and
152 cavity tones respectively (Fig. 3C), and jet angles of $99.2 \pm 15.3^\circ$ and $62.3 \pm 11.1^\circ$ ($n=5$) for wall-
and alar edge-tone, respectively (Fig. 3D). Jet angle was not predicted by the cavity tone model.
154 To test if these predicted lengths were consistent with the physical dimensions of the larynx, we
measured minimum wall jet length (x_{wall}), alar edge jet length (x_{alar}), and cavity length (x_{cav}) on
156 Dice- μ CT scans of the individual larynges (Fig. 3A). For the wall impingement model the
predicted jet lengths were not significantly different from the minimum length (two-tailed paired
158 sample t-test, $p = 0.09$), and importantly fell within the physical range in all five cases (Fig. 3C).
However, the predicted jet length for the alar edge-tone model was significantly shorter than the
160 anatomical length (two-tailed paired sample t-test, $p = 0.003$), and fell 0.26 ± 0.07 mm too short to
reach the alar edge (Fig. 3C). The predicted cavity length for the cavity-tone model was also
162 significantly shorter than the anatomical length (two-tailed paired sample t-test, $p = 0.020$), (Fig.
3C). Therefore, these experiments support the wall-tone whistle mechanism.

164
To further test if the predicted jet length and angles were consistent with intralaryngeal flow, we
166 performed Computational Fluid Dynamics simulations (37) of airflow through a 3D-reconstructed
larynx in USV producing state (Fig. 3E, See Methods and Materials). Using the same boundary
168 conditions as under experimental settings, our CFD model showed first of all that jet formation
occurred with jet separation points at the dorsal and ventral side of the cartilaginous glottis (Fig.
170 3F,G; Additional File 2). Second, the jet impinged on the thyroid planar wall and not the alar edge.
Third, the jet was 0.76 mm long, at a 98.0° angle and had a speed of 36.5 m/s, which was in
172 excellent agreement with the predicted $x_{wall} = 0.71$ mm at 86.6° and 33.2 m/s of our aeroacoustic
model for that individual (Fig. 3E-G). The simulated jet angle is also in excellent agreement with

174 the earlier estimate in the mouse larynx (25). Taken together, the predicted jet lengths and flow
structure from CFD simulation provide evidence against both alar-edge and shallow cavity-tone
176 models, and support the intra-laryngeal planar impinging jet model of USV sound production in
rats.

178

The third distinct feature between the wall, alar edge and cavity tone models is the required
180 presence of the alar edge and a small airsac-like cavity rostral to the vocal folds, called the
ventral pouch, which is found in several muroid rodent species (27, 38, 39). The wall tone model
182 allows air circulation in the ventral pouch, but does not require its presence because the feedback
that stabilizes the tone comes from acoustic waves within the jet (22–25). The alar edge tone
184 model on the other hand evidently requires the presence of the alar edge and suggests that
pouch cavity resonance properties affect sound frequencies (27). The cavity tone model too
186 requires the presence of the ventral pouch for air circulation. To test if the alar edge and ventral
pouch are essential to fUSV production, we prevented both the presence of an edge and air
188 circulation in the pouch by filling the pouch with a small aluminum sphere in excised rat (n=7) and
mice (n=6) larynges. In rats, six out of seven larynges retained fUSV production after sphere
190 insertion (Fig. 4) and the mean, minimum, and maximum peak frequencies (f_p) did not change
significantly in rats (two-tailed paired sample t-test, N=7, mean f_p : $p = 0.89$, max f_p : $p = 0.71$, min
192 F_p : $p = 0.87$) and mice (N=6, mean f_p : $p = 0.65$, max f_p : $p = 0.48$, min f_p : $p = 0.45$). To estimate
how filling the ventral pouch affected the intralaryngeal flow, we performed CFD simulations of
194 the same experimental manipulation (Fig. 4E,F; Additional File 3). A glottal jet formed that
impinged on the thyroid planar wall slightly more rostral due to the sphere, leading to a slightly
196 increased angle (103° , +5.1%) and jet length (0.79 mm, +2.6%). Thus, neither the ventral pouch
or the alar edge is essential for USV production in rats and mice

198

Finally, we used CFD simulations to test if the proposed flow scenario (27) for the alar edge
200 model in Fig 1B is physically plausible. We ran CFD simulations on the previously published 3D
reconstructed rat vocal tract (27) that has abducted VFs and arytenoids (See methods). Driven by

202 *in vivo* tracheal pressure, our simulations show that no intralaryngeal jet is formed and no air
circulates in the ventral pouch (Additional File 4). Therefore we can conclude that the suggested
204 flow scenario for the alar edge model (27) is not physically accurate.

206 Our data conclusively shows that both in the most widely used rodent models in biological and
medical research, rats and mice, USVs are produced by an aerodynamic wall impingement
208 whistle. The three distinctive features - closed vocal fold adduction state, jet properties and non-
essential presence of edge and pouch - provide evidence against alar edge and shallow cavity
210 tones and support the wall tone. The notion that wall impingement is incongruent with laryngeal
anatomy (27) is thus incorrect. However, given the large diversity of laryngeal morphology and life
212 history found in the 1500 species of rodents (40), our data does not exclude that multiple
mechanisms contribute to USV production in other rodents species such as singing (41, 42) or
214 grasshopper mice (43). Shallow cavity tones (35) provide an alternative mechanism to explain the
loud USVs of several new world rodents with more pronounced alar and pouch structures and
216 may be a wide-spread mammalian sound production mechanism that requires further
investigation.

218

In vivo rodent USVs are characterized and classified by the time-varying frequency trajectories of
220 syllables (18, 30 31). Based on our aerodynamic model of USV production, we have implemented
a quantitative data-driven model of *in vivo* USV motor control (Methods and Materials). Our
222 aerodynamical model of USV production predicts that the frequency of pressure and flow
structure variations are set by the jet speed and jet length. The frequency of these whistles is 50-
224 100 kHz and the pressure fluctuation thus occur at the microsecond scale and are at least three
orders of magnitude faster than the millisecond laryngeal motor control (30, 44, 45) of the jet
226 parameters jet speed and jet length. As a consequence, the USV instantaneous frequency can be
considered time-invariant compared to the motor control that shapes the frequency trajectories.
228 Furthermore, in contrast to an earlier suggestion (27), the fact that USV exhibit changes *in vivo*
does thus not inform on the aerodynamical mechanism. We focused on rats where pressure, flow

230 and muscle electromyography data has been measured during USV production *in vivo* (30, 44,
45). Within correct anatomical and physiological ranges, the x , u control space produces all
232 frequencies observed *in vivo* (Fig. 5A,B). We used an orifice constriction model that accurately
estimated tracheal mass flow from pressure (Fig. 5C) to calculate how subglottal pressure and
234 glottal area affect frequency (Fig. 5D). Surprisingly, glottal area barely affects frequency because
the increase in jet speed from decreasing area is counteracted by the decrease in flow.

236

Two motor systems drive the parameters of our model; first, the respiratory muscles that control
238 subglottal pressure and second, intrinsic laryngeal muscles that control laryngeal geometry, such
as glottal area and impingement length. Because rodent laryngeal muscles share developmental
240 origin (46), location and function (39) with other vertebrates, we based their mechanical actions
on better studied mammals such as human (47) and dog (48–50). We included three muscle
242 groups; respiratory muscles (RM) that control subglottal pressure, the cricothyroid muscle (CT)
that controls impingent length, and a combination of intrinsic laryngeal muscles (Thyroarytenoid
244 (TA), posterior cricoarytenoid (PCA) and interarytenoid (IA) muscles) that set vocal fold adduction
and thereby glottal area (Fig. 5E; *Methods and Materials*). Consistent with earlier observations in
246 mice (25), with increasing CT force and decreasing x , USV frequency goes down. Interestingly,
the CT has thus the opposite function compared to vocal fold vibration driven voiced sound
248 production where CT shortening increases frequency (29, 49, 51). The laryngeal muscles affect
the jet shape and flow conditions that determine whistle stability, which gates the sound on and
250 off (*Methods and Materials*). The three muscle groups together affect USV frequency in a highly
redundant control space (Fig. 5F-H), which makes it challenging to invert the system and
252 estimate control parameters from sound alone. However, with additionally known factors such as
pressure or flow, and at higher frequencies where the jet becomes unstable, this redundancy
254 collapses (Fig. 5H).

256 We computed putative *in vivo* motor control trajectories of 22 kHz and 50kHz USV calls (6) from
acoustics and corresponding *in vivo* subglottal pressure (30, 44) (Fig. 5I-K). Our model can

258 reproduce these call types including several subtypes, such as flat and modulated trill calls
including frequency jumps (Fig. 5I), and accurately predicts that pressure increases and flow
260 decreases during USVs consistent with *in vivo* recordings (30, 32, 52). Moreover, increased TA
and CT force correlates with higher frequencies (Fig. 5K) consistent with *in vivo* recordings (44) to
262 counteract abductive forces of increased respiratory pressure and to overcome whistle instability.
Lastly, we explored the effects of changing larynx geometry with unchanged motor control
264 trajectories and show that change of only 180 μm (20%) can cause frequency shifts of 10 kHz,
which is similar magnitude observed in behavioral models (16, 17). Taken together our simple
266 model provides a physiological basis for the neuromuscular control of USVs and interpreting rat
and mouse USV call phenotypes.

268
Mice and rat USVs often contain distinct frequency jumps that play an important role in call type
270 classification (2, 53). These jumps occur on the millisecond scale and do not correlate with either
laryngeal muscle activity or pressure (30, 45). Our aerodynamic model predicts that these
272 frequency jumps are jumps between stable whistling modes which explains why they can overlap
in vivo (53). Our motor model includes jet stability criteria that predict which simultaneous modes
274 are stable, and these seem to correspond well (Fig. 5I) with *in vivo* observed jumps in rats (2).
What exact modes are finally produced *in vivo* depends on local flow conditions at the jet exit and
276 needs further investigation.

278 Detailed control of laryngeal muscles is crucial in shaping USVs and connecting spiking motor
neurons to muscle action and laryngeal biomechanics requires more complex modelling
280 approaches and additional knowledge on motor neuron and muscle properties, motor unit
organization, and mechanical effects of muscle shortening. Interestingly, laryngeal muscles are
282 typically very fast, but the cricothyroid muscle's origin is somatic and it is slower than e.g. TA in
many mammals including rats (54, 55). Because our model suggest that USV frequency is
284 predominant set by CT action and respiratory pressure, both rather slow systems compared to
the precision and speed of other vocal production system, such as birdsong (56, 57) and bat

286 echolocation (58), this may explain the slower cadence of frequency modulation in in rat and mice
USVs.

288

The brain is constrained and modulated by the biomechanics, morphology and material
290 properties of the body (59–61). Our findings show that both neural and anatomical components
contribute to USV production (Fig. 5L). Therefore, the mechanisms that drive changes in strain
292 specific USVs or USV changes in mouse and rat disease models can be both altered motor
programs and laryngeal geometry. This emphasizes the importance of an embodied approach to
294 USV motor control to provide a physiological basis for USV syllable categorization and
interpreting rat and mouse vocal behavior phenotypes.

296

298 **Materials and Methods**

Subjects

300 We used 16 male Sprague Dawley rats, 11 juveniles (51 to 78 days old) and 5 adults, and 6 adult
male C57BL/6 mice. All animals were housed at Odense University Hospital. All experiments
302 were conducted at the University of Southern Denmark and were in accordance with the Danish
Animal Experiments Inspectorate (Copenhagen, Denmark).

304 **Larynx dissection and mounting**

All animals were euthanized with fentanyl/fluanisone or carbon dioxide, and kept on ice
306 (maximally 180 min). The trachea, larynx, and surrounding tissue were dissected, flash frozen in
liquid nitrogen and stored at -80°C. Before each experiment, we thawed the tissue in a
308 refrigerator and then submerged it in refrigerated ringer's solution (62) in a dish on ice and
removed additional tissue surrounding the larynx and trachea. We then mounted the larynx in the
310 setup. For rats, we mounted the larynx on a plastic Luer connector (1.1mm inner diameter and
1.6 mm outer diameter), filed down so that the tip was a straight tube. For mice, we mounted the
312 larynx on a rounded, blunt 19G needle. The larynx was slid over the tube connector until the

caudal edge of the cricoid touched the tube exit and secured with a suture around the trachea, 6-
314 0 braided silk suture (Deknatel, USA) for rats, and 10-0 monofilament suture for mice.

Experimental setup

316 We mounted larynges in an excised larynx setup described in detail in (25). In brief, this setup
(Figure S1), allows for running humidified air through the larynx at precisely controlled pressure,
318 while simultaneously measuring volumetric flow, pressure, and sound. The position of arytenoid
flanges is controlled by micromanipulators. The rate of volumetric flow through the larynx was
320 measured using a MEMS flow sensor (PMF2103, Posifa Microsystems, San Jose, USA). Sound
was recorded using a 1/4-inch pressure microphone-pre-amplifier assembly (model 46BD,
322 frequency response $\pm 1\text{dB}$ 10 Hz – 25 kHz & $\pm 2\text{dB}$ 4 Hz – 70 kHz G.R.A.S., Denmark) located 5
cm above the larynx pointing downwards and to the side of the larynx as not to be hit by the
324 airflow leaving it (Fig. S1). The microphone signal was amplified by 40 dB for rats and 50 dB for
mice (amplifier 12AQ, G.R.A.S., Denmark). We calibrated the microphone before each
326 experiment (Calibrator 42AB, G.R.A.S., Denmark). The positions of the larynx and microphone
were fixed relative to each other (Fig. S1). The sound, pressure and flow signals were low-pass
328 filtered at 100, 10 and 10 kHz, respectively (filter model EF502 low pass filter DC – 100 kHz and
EF120 low pass filter DC – 10 kHz, Thorlabs, U.S.A.) and digitized at 166, 224 (mice), or 240 kHz
330 (USB 6259, 16 bit, National Instruments, Austin, Texas). All control and analysis software were
written in MATLAB 2018a (MathWorks).

332 We imaged laryngeal configuration during ramps with a Leica DC425 camera mounted on a
stereomicroscope (M165-FC, Leica Microsystems) or with a high-speed camera (MotionPro X4-
334 M-4, Integrated Design tools, Inc., USA) at 250 fps. The DC425 camera was controlled using LAS
(Leica Application Suite Version 4.7.0, Leica Microsystems, Switzerland), and the high-speed
336 camera was controlled using Motion Studio (x64, Version 2.10.01, Integrated Design tools, Inc.,
USA). We illuminated the larynges with Leica GLS150 lamp (photography) or Leica EL6000
338 (highspeed imaging) through a liquid light guide connected to the stereomicroscope.

340 **Experimental protocols**

We performed three experiments to study USV production in the larynx *in vitro*. In all
342 experiments, we applied air pressure ramps from 0 to up to 2 kPa.

Protocol 1 – Vocal fold adduction. We first applied a pressure ramp in resting state without any
344 vocal fold or arytenoid adduction. Because the airflow typically pushed the arytenoid flanges
apart, we next approximated the arytenoid flanges with suture (Suture: Black Polymaide
346 Monofilament USP. 10-0 (0.2 metric) 13 cm, Needle: Taper Point, 4mm, 70 μ , 90°. AROSurgical
Instruments Corporation, USA) to stabilize the glottis dorsally. Next, we applied pressure ramps
348 with 1) the vocal folds in rest position and 2) with the vocal folds adducted using two adduction
methods. First, we adducted the vocal folds using micromanipulators. Next we glued the vocal
350 folds in adducted state by applying cyanoacrylate tissue glue (3M Vetbond, TissueAdhesive –
1469-SB, 3M Animal Care, U.S.A) with a pulled glass micropipette to the rostral side of the vocal
352 folds in an adducted state. We recorded the glottal configuration using high-speed video (250 fps)
and still image camera for 6 and 4 larynges, respectively. We obtained complete datasets in 10
354 rats.

Protocol 2 – USV production in fixed larynges. After the last ramp of Protocol 1, for five
356 animals we coated the outside of the larynx in UV-glue (Loon outdoors, UV FLY clear finish, thick,
USA) and placed the larynx and mounting tube in 4% PFA. After 7 days, we mounted the fixed
358 larynx in the setup and applied a pressure ramp to test if fUSVs were produced.

360 The larynx was stained for two days in 15% Lugol solution, 1 day in 10% Lugol solution, and
1 day in 5% Lugol solution on a roller mixer (Stuart SRT6D, Cole-Parmer, UK) at 6 rpm. The
362 samples were then rinsed in distilled water for 2 times 10 minutes on the roller mixer at 12 rpm,
and scanned in a μ CT scanner (μ CT50, Scanco Medical AG, Brüttisellen, Switzerland, 8 μ m
364 resolution) at Odense University Hospital. We obtained complete datasets in 5 rats.

Protocol 3 – USV production with filled ventral pouch in rats and mice. We applied pressure
366 ramps with subsequently 1) the vocal folds and arytenoids adducted (as in Protocol 1), and 2)

368 with an aluminum sphere placed in the ventral pouch. This size sphere fitted exactly in the pouch
to fill it completely and caused the alar edge to lay on top of the sphere (Fig. 3B).

370 Based on measurements from CT scans, we used a 0.8 mm diameter sphere for rats, and a
0.5 mm sphere for mice, to fill the pouch. We then subjected the larynges to a pressure ramp. For
372 rats, we used a ramp from 0 to up to 1.5 kPa and down to 0 kPa again at a rate of 0.5 or 1.66
kPa/s. For mice, we used a ramp from 0-2 kPa at a rate of 0.5 kPa/s. The position of the sphere
374 in the ventral pouch was confirmed with a photo before and after the pressure ramp We obtained
complete datasets for 7 rats and 6 mice.

376

Signal conditioning and USV extraction

378 Digitized sound signals were bandpass filtered (3rd order Butterworth filter; 2.5-100 kHz (Protocol
1, 2, and 4 rats and all mice in protocol 3) or 2.5 - 83.3 kHz (3 rats in Protocol 3) with zero-phase
380 shift implementation *filtfilt* function). We calculated spectrograms (nfft = 2048 or 1024 (for 3 rats in
Protocol 3), overlap: 0%, Hamming window). For each time bin, we calculated mean flow (V) and
382 Shannon's entropy (63) scaled to $\log_2(nfft/2)$ of the spectrogram's power distribution between
15-100 kHz (All but 3 rat individuals in protocol 3) or 15-82.3 kHz (3 rat individuals in Protocol 3).
384 Entropy was averaged over six time bins for rats, and three time bins for mice. Because turbulent
air flow at high flow rates produces white noise up to 100 kHz, we designed a subjective detector
386 for USV whistles over flow-induced noise. We used the pressure ramps recorded from completely
unadducted larynges (N = 10, *Protocol 1*), calculated the mean Shannon's entropy during
388 maximum flow and used this value (0.7) for all other rat ramps to detect USVs. For mice, a
Shannon's entropy limit of 0.8 was used (based on visual inspection). Any continuous period of
390 sound below these threshold levels, allowing for breaks of 1 time slice, were considered USVs.
For time bins with USVs we extracted the peak frequency (f_p) using the *tfridge* function.

392

394

Comparison between model predictions of jet length and laryngeal geometry

396 **Laryngeal geometry reconstruction and quantification.** The diceCT scans were analyzed in
Amira (Amira 5.2.1, 2009, Konrad-Zuse-Zentrum Berlin (ZIB), Visage Imaging Inc.). An oblique
398 slice was placed in the sagittal plane, and another oblique slice was placed perpendicular to the
first one, and overlapping the glottal opening (Fig. 1A). The slices were exported as TIF-images
400 and imported into ImageJ (Version 1.52a, Wayne Rasband, National Institute of Health, USA) for
measuring the following laryngeal dimensions: On the cross-section overlapping the glottal
402 opening, we measured the glottal area, A_{gl} , as the area of a polygon manually fit into the glottal
constriction on the corresponding cross section (Fig. 3A, right). On the midline cross section, we
404 measured the shortest and longest distances between the point of jet formation and the ventral
intralaryngeal wall, x_{alar} and x_{max} (Fig. 3A, left), i.e. the range of jet lengths that could possibly fit in
406 the larynx. Here we also measure the length of the opening of the ventral pouch, x_{cav} . The point of
jet formation was approximated as the point in the middle of the distance between the adducted
408 arytenoids and adducted vocal folds (Fig. 3A, left). We assumed bilateral/axial symmetry for the
jet, i.e. that its direction was parallel to the sagittal plane.

410

Mode analysis. In order to compare the jet length predictions based on the aerodynamic models
412 corresponded with internal laryngeal geometry during USV production, we needed to identify
which mode was extracted from the fUSVs. Both the jet impingement model and the alar edge-
414 tone model predict the frequencies of several modes and therefore it was paramount to identify
the mode numbers of fUSVs. We manually selected fUSVs where multiple modes were visible
416 and compared the frequencies of other modes to the dominant frequency, f_p , over time (Fig. 3B)
using the tfridge MATLAB method on the spectrogram. The frequencies of the modes above the
418 first one, f_0 , are given as $f_n = n \cdot f_0$ (where $n = 2, 3, 4, \dots$). The difference in frequency of two
adjacent modes is thus equal to f_0 and the mode of the dominant frequency can be calculated as
420 $n = \frac{f_p}{\Delta f}$, where Δf is the difference between the frequency of the dominant mode and the closest
mode, equal to f_0 if the modes are of adjacent mode number. The frequency of the first mode was

422 then calculated as $f_0 = \frac{f_p}{n}$.

Jet geometry predictions. The models predict different jet lengths:

424 $x_{wall} = \frac{u}{f_0}$, for the wall impingement model, (25)

$x_{alar} = \frac{u}{2 \cdot f_0}$, for the alar edge-tone model, (27)

426 $x_{cav} = \frac{u \cdot (n - \gamma) \cdot \kappa}{f_n}$, for the cavity-tone model, (35)

where f_0 is the fundamental frequency, u is the mean convection speed of downstream

428 moving coherent structures, approximated as the jet exit speed $u = \frac{V}{A_{gl}}$ (Fig. 3B, right), V is

volumetric flow rate and A_{gl} is glottal constriction area.

430 Jet angle was determined by first fitting the predicted jet length between jet exit midpoint
and the ventral intralaryngeal wall on the midline cross section. We then measured the angle
432 between the resulting line and the midline of the cartilaginous glottis in ImageJ (Fig. 3A). As the
jet length predictions for the edge-tone model were too short to reach the alar edge, we were
434 unable to measure jet angle resulting from fitting x_{alar} between the jet exit midpoint and a point on
the alar edge, but in theory, the alar edge-tone model predicts jet angles similar to α_{min} (Fig. 3A,
436 left). For the cavity tone model, we did not investigate jet angle, as the model does not rely on jet
formation.

438

In vivo threshold flow estimate

440 We estimated tracheal air flow (V) during USV production in rats based on *in vivo* data. During
quiet respiration V is 15-20 ml/s (64). However during USV production V reduces, which is seen
442 in measurements of tracheal mass flow (30, 32, 52). We approximated V to be below 4 ml/s
during USV production (Fig. 3A in (52)) for a 250-300 gram animal. We then linearly corrected for
444 body size, which suggested that V during USV production for the animals used in this study was
below 10 ml/s.

446

Computational Fluid Dynamic simulations

448 We performed CFD (Computational Fluid Dynamic) simulations of air flowing through 3D-
reconstruction of intra-laryngeal rat airways with and without a sphere digitally added to the
450 ventral pouch (Fig. 3E-G & Fig. 4E-F). From the μ CT scan of one of the larynges, the laryngeal
airway was labeled in Amira. Under the experimental subglottal pressure condition, the mean jet
452 speed is estimated to be about 40 m/s. The according Mach number (defined as $Ma=u/c$, where u
is the mean jet speed, $c=346$ m/s is the speed of sound at 25°C) would be about 0.12. Therefore,
454 the flow is modeled as an incompressible flow. The governing equations are the three-
dimensional, unsteady, viscous, incompressible Navier-Stokes equation as below

$$\nabla \cdot \vec{U} = 0$$

$$\frac{\partial \vec{U}}{\partial t} + (\vec{U} \cdot \nabla) \vec{U} = -\frac{1}{\rho_0} \nabla P + \nu_0 \nabla^2 \vec{U}$$

456 where \vec{U} , ρ_0 , P , ν_0 are the incompressible flow velocity, density, pressure and kinematic
viscosity, respectively. $\nu_0=1.562 \times 10^{-5}$ m²/s and $\rho_0=1.184$ kg/m³ at 25°C. The computational solver
458 employs the sharp-interface immersed-boundary method (37). The laryngeal wall is represented
by triangular elements exported from Amira and smoothed through Meshlab. The governing
460 equation is solved on non-uniform Cartesian grids, with finest grids at the glottal jet region. A
1.0 kPa subglottal pressure is applied at the subglottal entrance. A non-penetration non-slip wall
462 boundary condition is applied at the laryngeal wall. Sensitivity studies on domain and grid size
showed that the numerical solution converges with a minimum grid of 10 μ m with a domain size
464 of 26 x 32 x 28 mm with a 0.8% difference in the jet speed.

466 To test if the proposed flow scenario for the alar edge model is physically plausible, we performed
CFD simulations of air flowing through a previously published intra-laryngeal rat airway (27). We
468 obtained the 3D geometry from Morphobank (www.morphobank.org, project ID 2686,
Morphobank media number M451228) and smoothed it in Meshlab. Simulation conditions were
470 as listed above. We applied a 0.9 kPa subglottal pressure at the subglottal entrance.

472 **Statistics**

All statistical testing was performed in MATLAB (MATLAB 2018a, MathWorks, USA). A two-tailed
474 t-test was performed to test if the slope of the frequency-to-flow relationship differed before and
after fixation of the larynges. As the core of the experiment was to predict jet length from peak
476 frequency, mode number, and jet speed, we compared x_{wall} predictions between fixed and fresh
larynges. For two of the fresh larynges, the mode number was difficult to determine and we
478 decided to perform the tests with mode numbers that fell within the range we saw for the fixed
larynges and that gave the predictions that best matched the corresponding predictions in the
480 fixed larynges. We performed a two tailed t-test comparing the difference between the x_{wall}
predictions from the fixed and fresh larynges to zero.

482 To compare the wall impingement model and alar edge-tone models' jet length predictions
to each other and to the internal laryngeal geometry during USV production, we performed two-
484 tailed paired sample t-tests to compare predicted and measured x_{wall} , x_{alar} , and x_{cav} , respectively.

We compared the peak frequencies of the fUSVs produced with and without a sphere filling
486 the ventral pouch and blocking the alar edge using two-tailed paired sample t-tests on mean,
maximum, and minimum frequencies of the fUSVs produced in the two treatments and for the two
488 sets of animals (rats and mice).

490 **Quantitative motor control model**

We constructed a quantitative data-driven model to capture how the activity of respiratory
492 muscles (RM, mainly the diaphragm muscle) and a combination of intrinsic laryngeal muscles
affect the main control parameters of our aerodynamic model: jet speed (u) and jet length (x).
494 Subglottal pressure increases linearly from 0 to 5 kPa with RM activity. Tracheal flow (V) was
predicted as a function of glottal area, tracheal diameter (measured from dice-CT scans), and
496 subglottal pressure by assuming the glottal constriction to constitute a tube with an obstruction
(65). We compared this obstruction model a ramp for a fixed larynx where glottal area, flow and
498 pressure were known. The model prediction aligned well with experimental data (Fig. 5C).

State-of-the-art measurements and 3D models of vocal fold adduction on canine larynges
500 (48-50) show how shortening of the adductor and abductor muscles sets glottal area. Based on
these insights, we modeled glottal area as sum of the membranous glottis (area between the
502 vocal folds) and cartilaginous glottis (areas between the arytenoid). The area of the membranous
glottis was set by Thyroarytenoid (TA) activity and the cartilaginous glottis is set by TA and a
504 combination of posterior cricoarytenoid (PCA) and interarytenoid (IA) muscles:

$$A_{gl} = (1 - TA)A_{max} + (1 - PCA \cdot IA)A_{max}$$

506 where A_{max} was measured from dice-CT scans (Fig. 3A). Because we lacked data on interaction
between the TA and PCA.IA parameters we assumed them to be coupled.

508 The jet speed was defined as tracheal flow divided by glottal area. Contraction of the CT
muscle rotates the thyroid wall away from the glottal opening (50), thereby increasing jet
510 impingement length, x (Fig. 5E). TA weakly counteracts this rotational action of the CT by
shortening the vocal folds (50), thereby decreasing the impingement length:

$$512 \quad x = x_{min} + (CT - 0.24 \cdot TA)x_{max}$$

where x_{min} was defined as 50% of the minimum predicted impingement length and x_{max} 150% of
514 the maximum predicted impingement length (see jet length prediction). The constant 0.24 was
chosen as it results in an impingement length of zero at 100% TA activation and 0% CT activation.

516 By chaining these functions together, we can predict frequency as:

$$f_0 = f_0 \left(u \left(A_{gl}(TA), V \left(P(RM), A_{gl}(TA, LCA, IA) \right) \right), x(CT, TA) \right)$$
$$518 \quad f_0 = \frac{V \left(P(RM), A_{gl}(TA, LCA, IA) \right)}{x(CT, TA)}$$

Lastly, we restricted the possible values for f_0 by implementing the whistle stability criteria
520 $d/x \leq St < 1$, where d is jet diameter and $St = f_0 \cdot d/u$ is the Strouhal number (25). Simulations were
implemented in Matlab and will be made available on Github. Because this is steady-state and
522 not dynamical model the time representation in Fig. 5 is arbitrary and chosen to fit experimental
data (44).

524

526 **References**

- 528 1. C. V. Portfors, D. J. Perkel, The role of ultrasonic vocalizations in mouse communication. *Curr. Opin. Neurobiol.* **28**, 115–120 (2014).
- 530 2. S. M. Brudzynski, Ethotransmission : communication of emotional states through ultrasonic vocalization in rats. *Curr. Opin. Neurobiol.* **23**, 310–317 (2013).
- 532 3. M. Fendt, M. Brosch, K. E. A. Wernecke, M. Willadsen, M. Wöhr, Predator odour but not TMT induces 22-kHz ultrasonic vocalizations in rats that lead to defensive behaviours in conspecifics upon replay. *Sci. Rep.* **8**, 11041 (2018).
- 534 4. I. Willuhn, *et al.*, Phasic Dopamine Release in the Nucleus Accumbens in Response to Pro-Social 50 kHz Ultrasonic Vocalizations in Rats. *J. Neurosci.* **34**, 10616–10623 (2014).
- 536 5. N. Simola, S. Granon, Ultrasonic vocalizations as a tool in studying emotional states in rodent models of social behavior and brain disease. *Neuropharmacology* **159** (2019).
- 538 6. S. M. Brudzynski, Ultrasonic calls of rats as indicator variables of negative or positive states: Acetylcholine-dopamine interaction and acoustic coding. *Behav. Brain Res.* **182**, 261–273 (2007).
- 540 7. G. D. Sewell, Ultrasonic Communication in Rodents. *Nature* **227**, 410 (1970).
- 542 8. G. D. Sales, “Ultrasonic calls of wild and wild-type rodents” in *Handbook of Mammalian Vocalization—An Integrative Neuroscience Approach*, S. M. Brudzynski, Ed. (Elsevier Academic Press, 2010), pp. 77–88.
- 544 9. D. E. Wilson, D. M. Reeder, *Mammal species of the world: a taxonomic and geographic reference. 3rd ed.* (Johns Hopkins University Press, Baltimore, Maryland, 2005).
- 546 10. A. C. Ericsson, M. J. Crim, C. L. Franklin, A brief history of animal modeling. *Mo. Med.* **110**, 201–205 (2013).
- 548 11. J. Fischer, K. Hammerschmidt, Ultrasonic vocalizations in mouse models for speech and socio-cognitive disorders: Insights into the evolution of vocal communication. *Genes, Brain Behav.* **10**, 17–27 (2011).
- 550 12. J. L. Silverman, M. Yang, C. Lord, J. N. Crawley, Behavioural phenotyping assays for mouse models of autism. *Nat. Rev. Neurosci.* **11**, 490–502 (2010).
- 552 13. M. Wöhr, F. I. Roulet, J. N. Crawley, Reduced scent marking and ultrasonic vocalizations in the BTBR T+tf/J mouse model of autism. *Genes, Brain Behav.* **10**, 35–43 (2011).
- 554 14. B. L. Zampieri, F. Fernandez, J. N. Pearson, M. R. Stasko, A. C. S. Costa, Ultrasonic vocalizations during male-female interaction in the mouse model of Down syndrome Ts65Dn. *Physiol. Behav.* **128**, 119–125 (2014).
- 556 15. M. L. Scattoni, *et al.*, Reduced ultrasonic vocalizations in vasopressin 1b knockout mice. *Behav. Brain Res.* **187**, 371–378 (2008).
- 558 16. K. L. Paumier, *et al.*, Intrastriatal injection of pre-formed mouse α -synuclein fibrils into rats triggers α -synuclein pathology and bilateral nigrostriatal degeneration. *Neurobiol. Dis.* **82**, 185–199 (2015).
- 560 17. E. Ey, *et al.*, The Autism ProSAP1 / Shank2 mouse model displays quantitative and structural abnormalities in ultrasonic vocalisations. *Behav. Brain Res.* **256**, 677–689 (2013).
- 562 18. M. L. Scattoni, L. Ricceri, J. N. Crawley, Unusual repertoire of vocalizations in adult BTBR T+tf/J mice during three types of social encounters. *Genes, Brain Behav.* **10**, 44–56 (2011).
- 564 19. L. H. Roberts, Evidence for the laryngeal source of ultrasonic and audible cries of rodents. *J. Zool.* **175**, 243–257 (1975).
- 566 20. L. H. Roberts, The rodent ultrasound production mechanism. *Ultrasonics* **13**, 83–88 (1975).
- 568 21. R. H. Henrywood, A. Agarwal, The aeroacoustics of a steam kettle. *Phys. Fluids* **25** (2013).
- 570 22. C.-M. Ho, S. N. Nossier, N. S. Nosseirp, Dynamics of an impinging jet. Part 1. The feedback phenomenon. *J. Fluid Mech.* **105**, 119–142 (1981).
- 572 23. D. Rockwell, E. Naudasher, Self Sustained Oscillations of Impinging Free Shear Layers. *Ann. Rev. Fluid Mech.* **11**, 67–94 (1979).
- 574 24. B. P. J. Morris, P. J. Morris, B. P. J. Morris, The spatial viscous instability of axisymmetric jets. *J. Fluid Mech.* **77**, 511–529 (1976).

- 582 25. E. Mahrt, A. Agarwal, D. Perkel, C. Portfors, C. P. H. Elemans, Mice produce ultrasonic
vocalizations by intra-laryngeal planar impinging jets. *Curr. Biol.* **26**, R880–R881 (2016).
- 584 26. G. D. Sewell, Ultrasonic signals from rodents. *Ultrasonics* **8**, 26–30 (1970).
- 586 27. T. Riede, H. L. Borgard, B. Pasch, Laryngeal airway reconstruction indicates that rodent
ultrasonic vocalizations are produced by an edge-tone mechanism. *R. Soc. open Sci.* **4**,
170976 (2017).
- 588 28. K. Inagi, E. Schultz, C. N. Ford, An anatomic study of the rat larynx: Establishing the rat
model for neuromuscular function. *Otolaryngol. - Head Neck Surg.* **118**, 74–81 (1998).
- 590 29. Y. S. Zhang, D. Y. Takahashi, D. A. Liao, A. A. Ghazanfar, C. P. H. Elemans, Vocal state
change through laryngeal development. *Nat. Commun.* **10**, 1–12 (2019).
- 592 30. T. Riede, Subglottal pressure, tracheal airflow, and intrinsic laryngeal muscle activity
during rat ultrasound vocalization. *J. Neurophysiol.* **106**, 2580–2592 (2011).
- 594 31. C. V Portfors, Types and functions of ultrasonic vocalizations in laboratory rats and mice.
J. Am. Assoc. Lab. Anim. Sci. **46**, 28–34 (2007).
- 596 32. L. H. Roberts, Correlation of respiration and ultrasound production in rodents and bats. *J.*
Zool. **168**, 439–449 (1972).
- 598 33. A. M. Johnson, M. R. Ciucci, J. A. Russell, M. J. Hammer, N. P. Connor, Ultrasonic output
from the excised rat larynx. *J. Acoust. Soc. Am.* **128**, EL75–EL79 (2010).
- 600 34. I. Sanders, D. Weisz, B. Y. Yang, K. Fung, A. Amirali, The mechanism of ultrasonic
vocalization in the rat. *Soc. Neurosci. Abstr.* **27(1)**, 241 (2001).
- 602 35. J. E. Rossiter, Wind-tunnel experiments on the flow over rectangular cavities at subsonic
and transonic speeds. *RAE Tech. Rep. No. 64037* (1964).
- 604 36. L. Shaw, R. Clark, D. Talmadge, F-111 generic weapons bay acoustic environment. *J.*
Aircr. **25**, 147–153 (1988).
- 606 37. R. Mittal, *et al.*, A versatile sharp interface immersed boundary method for incompressible
flows with complex boundaries. *J. Comput. Phys.* **227**, 4825–4852 (2008).
- 608 38. G. Smith, Structure of the normal rat larynx. *Lab. Anim.* **11**, 223–228 (1977).
- 610 39. L. B. Thomas, J. C. Stemple, R. D. Andreatta, F. H. Andrade, Establishing a New Animal
Model for the Study of Laryngeal Biology and Disease: An Anatomic Study of the Mouse
Larynx. *J. Speech Lang. Hear. Res.* **52**, 802–811 (2009).
- 612 40. R. M. May, How Many Species Are There on Earth? *Sci.* **241**, 1441–1449 (1988).
- 614 41. D. E. Okobi, A. Banerjee, A. M. M. Matheson, S. M. Phelps, M. A. Long, Motor cortical
control of vocal interaction in neotropical singing mice. *Sci.* **363**, 983–988 (2019).
- 616 42. A. Banerjee, S. M. Phelps, M. A. Long, Singing mice. *Curr. Biol.* **29**, R190–R191 (2019).
- 618 43. B. Pasch, I. T. Tokuda, T. Riede, Grasshopper mice employ distinct vocal production
mechanisms in different social contexts. *Proc. R. Soc. B Biol. Sci.* **284**, 2–11 (2017).
- 620 44. T. Riede, Stereotypic Laryngeal and Respiratory Motor Patterns Generate Different Call
Types in Rat Ultrasound Vocalization. *J. Exp. Zool. Part A Ecol. Genet. Physiol.* **319**, 213–
224 (2013).
- 622 45. T. Riede, Rat Ultrasonic Vocalization Shows Features of a Modular Behavior. *J. Neurosci.*
34, 6874–6878 (2014).
- 624 46. J. M. Tabler, *et al.*, Cilia-mediated hedgehog signaling controls form and function in the
mammalian larynx. *Elife* **6**, 1–26 (2017).
- 626 47. A. Gömmel, C. Butenweg, K. Bolender, A. Grunendahl, A muscle controlled finite-element
model of laryngeal abduction and adduction. *Comput. Methods Biomech. Biomed. Engin.*
10, 377–388 (2007).
- 628 48. J. T. Heaton, J. B. Kobler, D. M. Otten, R. E. Hillman, S. M. Zeitels, Development of a
Closed-Loop Stimulator for Laryngeal Reanimation: Part 2. Device Testing in the Canine
Model of Laryngeal Paralysis. *Ann. Otol. Rhinol. Laryngol.* **128**, 53S–70S (2019).
- 630 49. D. K. Chhetri, J. Neubauer, D. A. Berry, Neuromuscular control of fundamental frequency
and glottal posture at phonation onset. *J. Acoust. Soc. Am.* **131**, 1401–1412 (2012).
- 632 50. B. Geng, N. Pham, Q. Xue, X. Zheng, A three-dimensional vocal fold posturing model
based on muscle mechanics and magnetic resonance imaging of a canine larynx. *J.*
Acoust. Soc. Am. **147**, 2597–2608 (2020).
- 634 51. C. L. Ludlow, Central nervous system control of the laryngeal muscles in humans. *Respir.*
Physiol. Neurobiol. **147**, 205–222 (2005).

- 638 52. C. Hegoburu, *et al.*, The RUB cage: Respiration-ultrasonic vocalizations-behavior
640 acquisition setup for assessing emotional memory in rats. *Front. Behav. Neurosci.* **5**, 1–13
(2011).
- 642 53. T. E. Holy, Z. Guo, Ultrasonic songs of male mice. *PLoS Biol.* **3**, 1–10 (2005).
- 644 54. J. F. Y. Hoh, “Laryngeal muscles as highly specialized organs in airway protection,
646 respiration and phonation” in *Handbook of Mammalian Vocalization*, Handbook of
648 Behavioral Neuroscience., S. M. Brudzynski, Ed. (Elsevier, 2010), pp. 13–21.
- 650 55. J. F. Y. Hoh, Laryngeal muscle fibre types. *Acta Physiol. Scand.* **183**, 133–149 (2005).
- 652 56. Z. Chi, D. Margoliash, Temporal precision and temporal drift in brain and behavior of zebra
654 finch song. *Neuron* **32**, 899–910 (2001).
- 656 57. C. P. H. Elemans, A. F. Mead, L. C. Rome, F. Goller, Superfast vocal muscles control
658 song production in songbirds. *PLoS One* **3**, 6–11 (2008).
- 660 58. C. P. H. Elemans, A. F. Mead, L. Jakobsen, J. M. Ratcliffe, Superfast muscles set
662 maximum call rate in echolocating bats. *Sci.* **333**, 1885–1888 (2011).
- 664 59. H. J. Chiel, R. D. Beer, The brain has a body: adaptive behavior emerges from
666 interactions of nervous system, body and environment. *Trends Neurosci.* **20**, 553–557
(1997).
- 668 60. K. Nishikawa, *et al.*, Neuromechanics: an integrative approach for understanding motor
670 control. *Integr. Comp. Biol.* **47** **1**, 16–54 (2007).
- 672 61. R. Pfeifer, M. Lungarella, F. Iida, Self-Organization, Embodiment, and Biologically Inspired
674 Robotics. *Sci.* **318**, 1088–1093 (2007).
- 676 62. C. P. H. Elemans, *et al.*, Syringeal muscles fit the trill in ring doves (*Streptopelia risoria* L.).
678 *J. Exp. Biol.* **209**, 965–977 (2006).
- 680 63. N. Madhu, Note on measures for spectral flatness. *Electron. Lett.* **45**, 1195–1196(1) (2009).
- 682 64. J. P. Mortola, A. Noworaj, Two-sidearm tracheal cannula for respiratory airflow
684 measurements in small animals. *J. Appl. Physiol. Respir. Environ. Exerc. Physiol.* **55**, 250–
253 (1983).
- 686 65. Y. A. Çengel, J. M. Cimbala, *Fluid Mechanics*, 3rd Ed. (McGraw-Hill, 2014).

668 **Declarations**

670 **Ethics approval:** All experiments were conducted at the University of Southern Denmark and
672 were in accordance with the Danish Animal Experiments Inspectorate (Glostrup, Denmark).

674 **Consent for publication:** Not applicable.

676 **Availability of data and materials:** The datasets used and/or analysed during the current study
678 are available from the corresponding author on reasonable request.

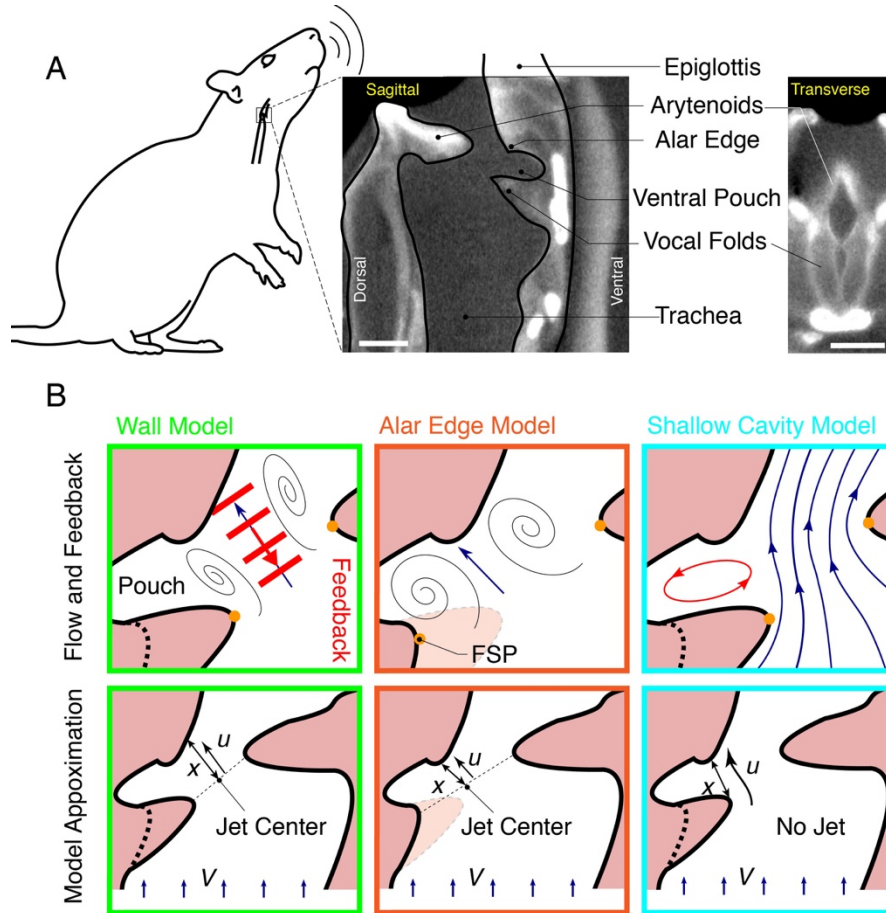
680 **Competing interests:** Authors declare that they have no competing interests.

682 **Funding:** This research was supported by a Janet Waldron Doctoral Research Fellowship, the
684 University of Maine (W.J.), and the Danish Research Council grant 7014-00270B
(C.P.H.E.).

686 **Author Contributions:** JH, AAG and CPHE designed research; JH, WJ, QX, XZ, MD, AAG and
CPHE performed research; QX, XZ, MD, CPHE contributed new reagents/analytic tools; JH, WJ,
QX, XZ, AAG and CPHE analyzed data; and JH, AAG and CPHE wrote the paper.

Acknowledgments: We thank T. Christensen for technical support; M. Wöhr for discussion and
comments on the manuscript.

Figures



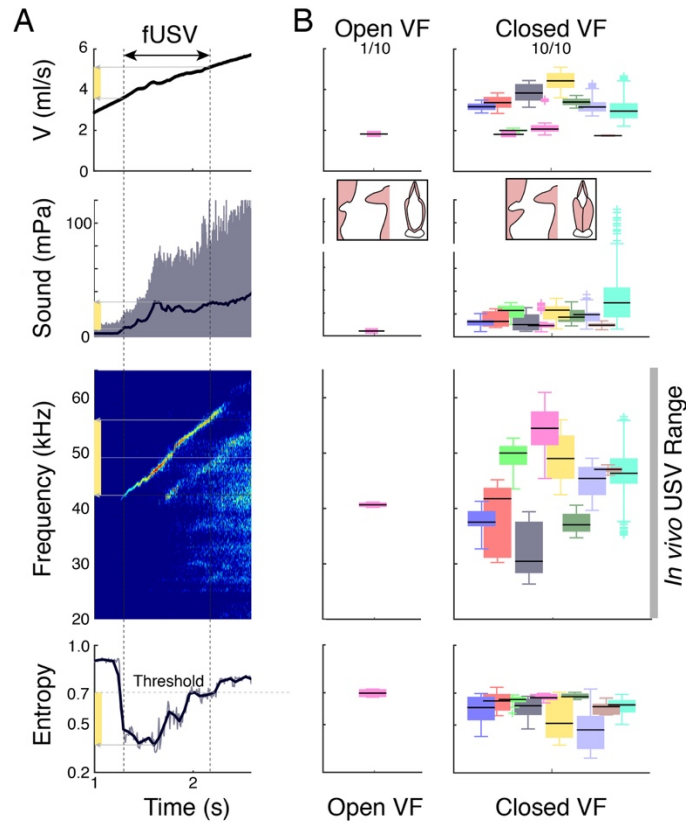
688

690

692 **Figure 1. Proposed aeroacoustic mechanisms of USV production in the rat and mouse**
 694 **larynx.** (A) Dice microCT scan of the rat larynx with cross-sections in medial sagittal plane (middle),
 696 and transversal plane parallel to the vocal folds (right). (B) Schematic of wall impingent (left), alar
 edge (middle), and shallow cavity (right) aerodynamic mechanisms of USV production in rats. The
 models are distinct in their local flow conditions (top row, black lines), feedback mechanism (red)

698

700



702

704

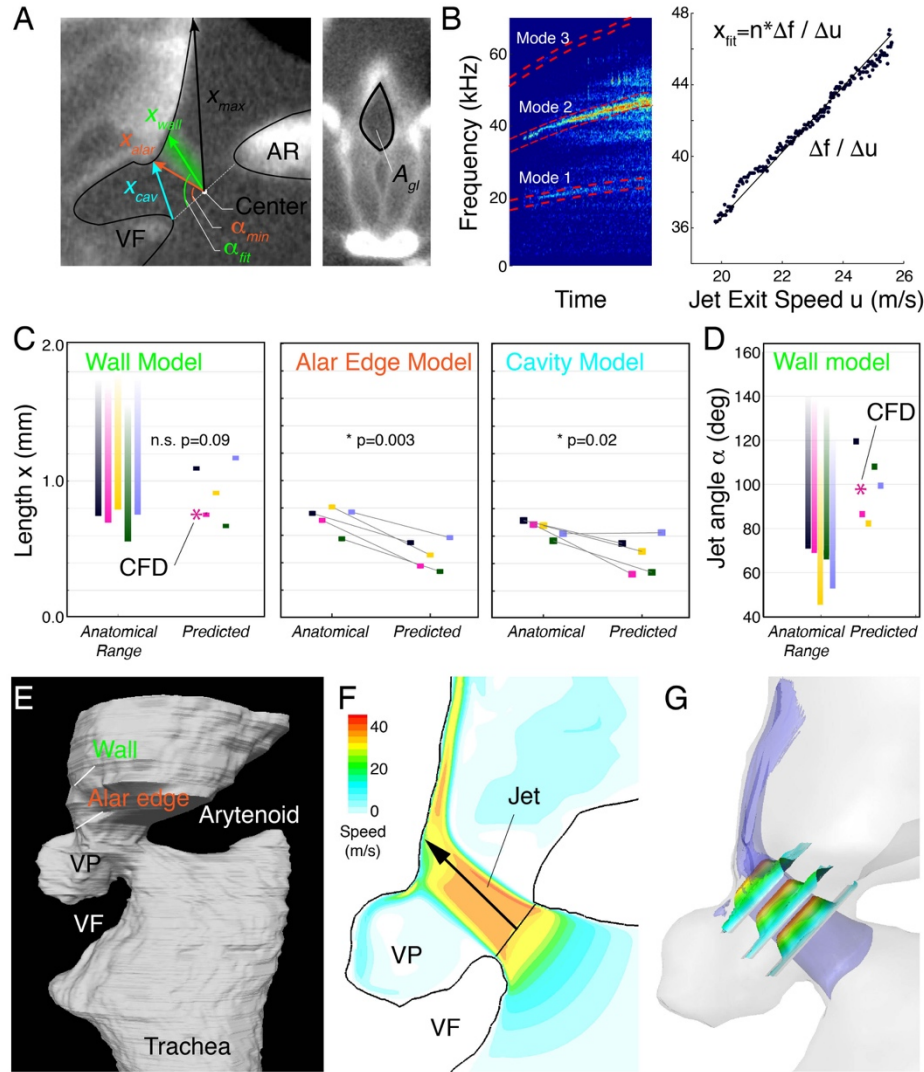
706

708

710

712

Figure 2. Rat fUSVs are produced with adducted vocal folds. (A) Above a threshold tracheal flow V , the isolated larynx produces fUSVs. From top to bottom: tracheal mass flow V , received sound pressure, sound spectrogram (NFFT=2048, overlap=50%, Hamming window), and scaled Shannon's entropy with the 0.7 threshold for USV detection. (B) With abducted vocal folds and open membranous glottis only 1 larynx produced fUSVs (left), while with adducted, opposed vocal folds all larynges ($N=10$) produced USVs (right) and within the *in vivo* frequency range of 18-96 kHz (31). Different colors represent different individuals.



714

716

718

720

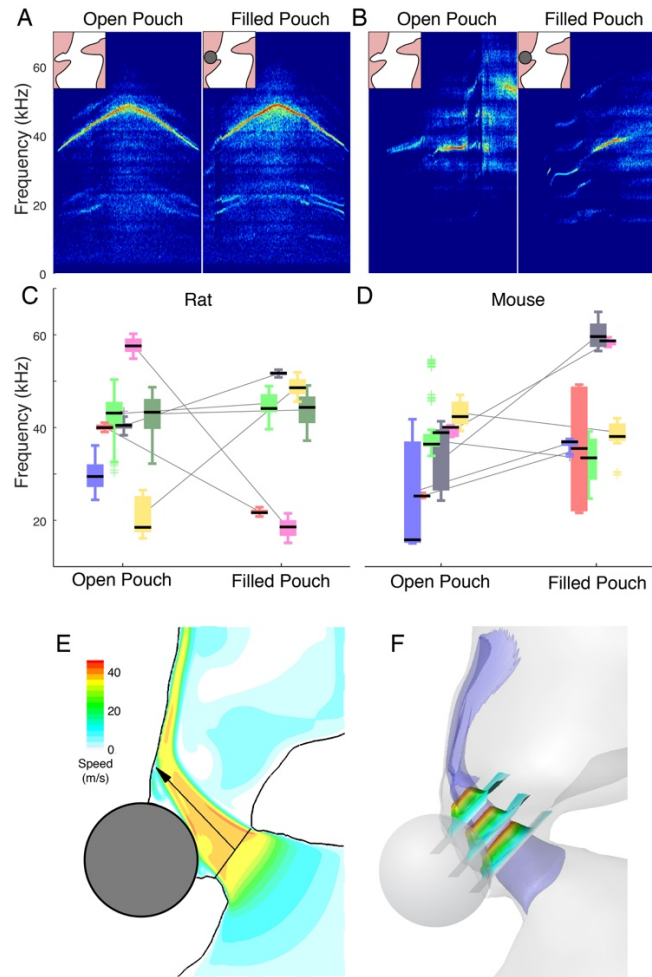
722

724

726

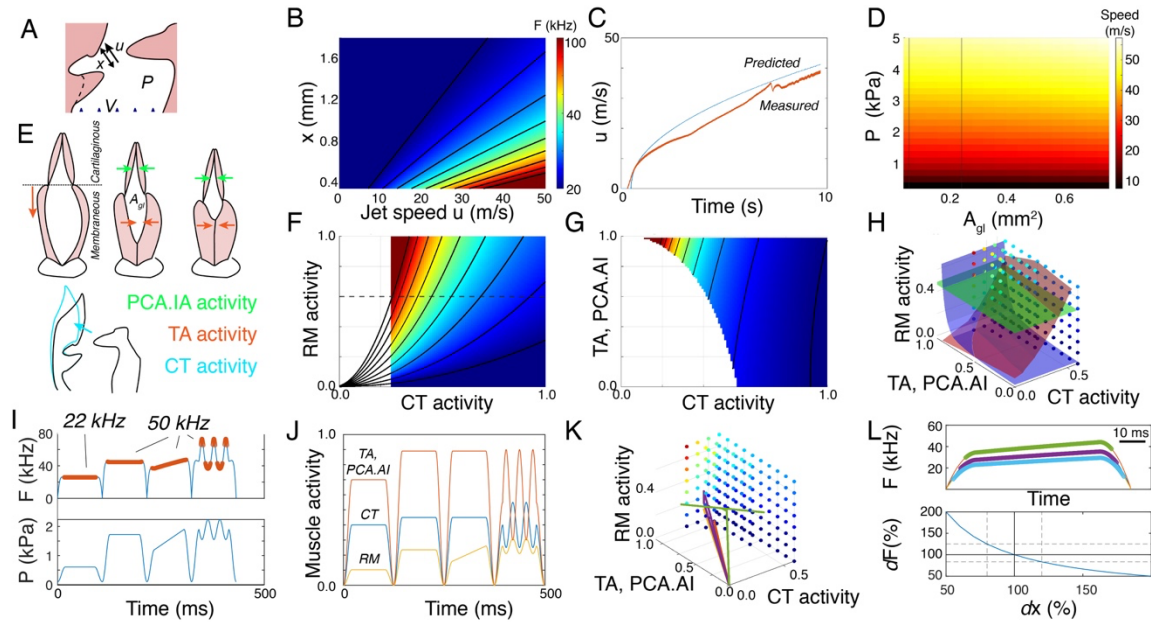
728

Figure 3. Glottal jet parameters support wall impingement model in rats. (A) The anatomical lengths of wall (x_{wall}) and alar edge (x_{alar}) jets, and ventral pouch cavity opening (x_{cav}) as measured in sagittal cross-sections of the glottis (left). Area of the cartilaginous glottis (A_{gl}) was measured in a transverse section parallel with the glottal opening (right). (B) Spectrogram (NFFT=2048, overlap=50%, Hamming window) of a fUSV shows multiple modes (red dashed boxes) essential to determine the dominant mode (See *Methods and Materials*). The slope between dominant frequency and jet speed equals the predicted jet/cavity length x (right). (C) Observed anatomical versus predicted values for x in wall, alar edge, and cavity model and (D) jet angle. These data show that wall-tone jet length and angle predictions fall within, while alar edge and cavity model predictions fall below the anatomical length range (C) or do not provide a solution for angle (D). (E) Flow was simulated in a fixed 3D mesh of the laryngeal airway. (F) 2D and (G) 3D flow show that a distinct jet is formed and impinges on the thyroid wall. Blue; iso-surface of jet speed equals 30 m/s. The three small planes present speed profiles and are contoured also by the speed value.



730

732 **Figure 4. The alar edge and ventral pouch are not required for USV production in rat and**
734 **mouse larynx.** Example spectrograms of normal fUSVs (left) and blocked alar edge and filled
736 ventral pouch (right) by small aluminum sphere in (A) rat and (B) mouse larynx. (C) Six out of 7 rat
738 larynges and (D) 6 out of 6 mice larynges produced fUSVs with filled ventral pouch. (E) Computational fluid dynamic simulations in the rat larynx slice and (F) 3D rendering show that also with a filled ventral pouch, a jet forms that impinges on the thyroid wall with negligible effect on the jet length and angle.



740

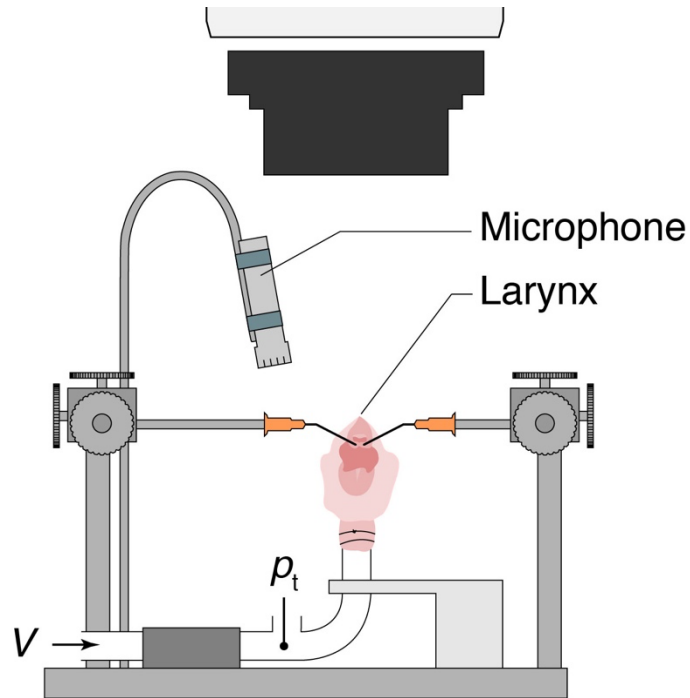
742 **Figure 5. Embodied motor control model extracts motor gestures for rat USVs call types.** (A)
 744 Exploring the parameter space of our aerodynamic wall impingement model results in (B)
 746 frequencies that overlap with in vivo range. Impingement length (x) and jet speed (u), tracheal flow
 748 V . (C) Predicted flow by orifice obstruction model (blue) corresponds well to measured flow (red)
 750 during subglottal pressure ramp through rat larynx in vitro. (D) Jet speed (u) as function of glottal
 752 area shows little dependency on glottal area in relevant range (vertical black lines represent 0.5
 754 and 2 times glottal areas from CT scans), whereas subglottal pressure strongly influences jet
 756 speed. (E) Effects of muscle shortening on laryngeal geometry (See *Methods and Materials*). Top,
 758 combinations of intrinsic laryngeal muscles affect the membranous and cartilaginous glottal area.
 760 Bottom, *m. cricothyroid* (CT) contraction leads to thyroid rotation (black to cyan outline), which
 762 increases impingement length. The rotatory action of CT is assumed to weakly counteracted by the
 764 smaller thyroarytenoid (TA) muscle. (F) Both respiratory muscle (RM) and CT activity affect USV
 frequency. The whistle is unstable in the white area. Black horizontal dashed line indicates the
 upper subglottal pressure limit during USVs in vivo ($p_{sub} = 3$ kPa). (G) TA action strongly
 influences the stability of the whistle and as such gates sounds, while it has little effect on f_0 .
 CT action affects both stability and f_0 . (H) Frequency f_0 is highly redundant in the three-
 dimensional motor space (red isosurface; $f_0 = 45$ kHz). Adding a given subglottal pressure
 (green isosurface; $p_t = 2.5$ kPa) and flow (blue isosurface; $V = 4.2$ ml/s) reduces this
 redundancy into a line or single point. (I) Driven by USV frequency (top; orange stable
 frequencies) and subglottal pressure (bottom), our model predicts muscle activity in (J) time
 and (K) as gestures in motor space for two common USV call types, 22 and 50 kHz, including
 the subtypes with frequency modulations and jumps. (L) Small changes in larynx geometry,
 such as impingement length x , alter the contours (top) and frequencies (bottom) of USVs
 whilst driven by identical motor gestures.

766

768

Additional Files

770



772 **Additional File 1.**

File format: Figure S1.jpg.

774 **Title:** Schematic of in vitro larynx sound production setup.

776 **Description:** The measurement position of tracheal pressure p_t and mass flow V are indicated. VF adduction is controlled with micro-manipulators.

778

780 **Additional File 2.**

File format: Movie M1.jpg.

782 **Title:** CFD simulation of airflow through rat larynx with adducted vocal folds (Fig 3FG).

784 **Description:** Flow was simulated in a fixed 3D mesh of the laryngeal airway. This movie shows that a distinct jet is formed and impinges on the thyroid wall.

786 **Additional File 3.**

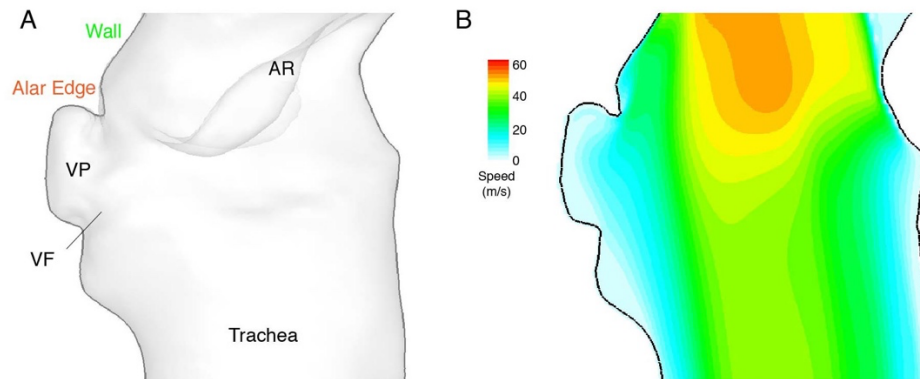
File format: Movie M2.jpg.

788 **Title: CFD simulation of airflow through rat larynx with filled ventral pouch (Fig 4EF).**

790 **Description:** This movie shows that also with a filled ventral pouch, a jet forms that impinges on the thyroid wall with negligible effect on the jet length and angle.

792

794



796 **Additional File 4.**

File format: Figure S2.jpg.

798 **Title:** CFD simulations does show no jet formation of a jet hitting alar edge or air circulation in ventral pouch in rat larynx with abducted vocal folds.

800 **Description:** (A) Geometry of previously published 3D airway reconstruction of the rat larynx (27). This air way geometry with abducted VFs was suggested to lead to USV production by either an air jet hitting the alar edge jet or by air circulation in the ventral pouch. (B) Modelled air speed through the air way (mass flow: 22 ml/s, tracheal pressure: 0.9 kPa) shows that no jet forms that hits either the alar edge or thyroid wall. Also, no air circulation takes place in the ventral pouch in this geometry. Therefore, the geometry of this air way does not support USV production following the alar edge model.

802

804

806