

# Nanomechanics of negatively supercoiled diaminopurine-substituted DNA.

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## ABSTRACT

Single molecule experiments have demonstrated a progressive transition from a B- to an L-form helix as DNA is gently stretched and progressively unwound. Since the particular sequence of a DNA segment influences both base stacking and hydrogen bonding, the conformational dynamics of B-to-L transitions should be tunable. To test this idea, DNA with diaminopurine replacing adenine was synthesized to produce linear fragments with triply hydrogen-bonded A:T base pairs. Triple hydrogen bonding stiffened the DNA by 30% flexurally. In addition, DAP-substituted DNA formed plectonemes with larger gyres for both B- and L-form helices. Both unmodified and DAP-substituted DNA transitioned from a B- to an L-helix under physiological conditions of mild tension and unwinding. This transition avoids writhing by DNA stretched and unwound by enzymatic activity. The intramolecular nature and ease of this transition likely prevent cumbersome topological rearrangements in genomic DNA that would require topoisomerase activity to resolve. L-DNA displayed about tenfold lower persistence length indicating it is much more contractile and prone to sharp bends and kinks. However, left-handed DAP DNA was twice as stiff as unmodified L-DNA. Thus, significantly doubly and triply hydrogen bonded segments have very distinct mechanical dynamics at physiological levels of negative supercoiling and tension.

## INTRODUCTION

The conformation and flexibility of DNA depend on electrostatics, base pairing, and base stacking. These interactions produce sequence specific characteristics that influence topology and protein binding. Denaturation directly reflects the hybridization of DNA but does not reveal manifold conformational dynamics of the base-paired state, which confront enzymes, that package, untangle, and process the information contained in DNA. In addition, such enzymes must contend with a continuum of flexural and torsional DNA rigidity conferred by oxidative damage or natural modifications like methylation. Other structural changes can be produced chemically to modify the electrostatic, aromatic, or hydrogen bonding of base pairs (1). For example, 2, 6-diaminopurine (DAP) is an alternative nucleobase that can substitute adenine in A:T base pairs to add an exocyclic amine moiety in the minor groove (2). This amine group forms an additional hydrogen bond with the C2 carbon of thymine in a DAP:T base pair. Complete substitution of adenine- with DAP-deoxyribonucleotide triphosphate in polymerase chain reactions produces triple hydrogen bonding throughout the amplicon. While this modification seems quite uncommon in nature (3,4), it could be useful in DNA-based nanomachines or origami. Such constructs can also serve to probe the activity of enzymes that modify DNA topology or process genetic information (5), since the substitution of DAP for adenine changes base stacking and hydrogen bonding without altering the charge along the molecule.

Substitution of DAP for adenine is known to increase the melting temperature of the B-form (1,6,7). Sequence-specific effects persist and the Santa Lucia model for calculating the melting temperature (8) can be adjusted by uniformly scaling the dinucleotide enthalpies (9). Other data regarding the cyclization of DAP-substituted DNA (DAP-DNA) fragments or binding and obligate tight curling around histone octamers to form nucleosomes indicates that DAP-DNA is flexurally stiffer than unmodified DNA (10,11,12). This is substantiated by increased stiffness exhibited by DAP-DNA molecules a few hundred- to kilo-bases in length deposited on surfaces for atomic force microscopy (9,11,12) and the increased persistence lengths determined using optical or magnetic tweezers to stretch single molecules of DAP-DNA (9,12). Overtwisting DAP-DNA under slight tension also exhibits increased flexural stiffness associated with creating plectonemic coils (5). While the persistence length of DAP-DNA reflects increased rigidity, the overstretching transition of this DNA occurs at a lower tension (9). Thus, DAP-DNA is more rigid and also more susceptible to mechanically induced, structural, phase changes.

Perhaps the most striking phase change exhibited by DAP-DNA is the conversion from a right- to a left-handed helical form. Although DAP-substitution maintains the charge density, the molecule adopts an unusual X-DNA conformation in response to slight concentrations of magnesium in potassium phosphate buffer or 50% concentrations of ethanol (13). This has been shown to be a left-handed helical form with a zigzag phosphate backbone due to alternating syn- and anti-deoxyribose orientations (14). Circular dichroism studies of oligonucleotides of dGC repeats show that high salt concentrations, approximately 60% solutions of alcohols, or added nickel drive the formation of left-handed Z-DNA (14, 15). Critical to this conformational change is the dehydration of the minor groove. Substituting

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3 diaminopurine for adenine has a similar effect by inserting an exocyclic amine that displaces water from  
4 the minor groove, but in an appropriate solvent stabilizes the X-DNA form, which displays a highly  
5 negative CD signal at 280 nm (13). The fact that X-form DAP-DNA is recognized by Z-DNA antibodies  
6 and  $P^{31}$  NMR studies show phosphate signals similar to those recorded for poly(dAT) in an X-DNA form  
7 indicate a zigzag backbone.  
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11 It is noteworthy that low levels of unwinding and tension have also been shown to drive the transition  
12 from right- to left-handed helices in GC segments of DNA (16), which suggested that triple hydrogen  
13 bonding may favor the right- to-left transition. To investigate this possibility, magnetic tweezers were  
14 used to gently stretch and unwind both unmodified and DAP-substituted DNA. Magnetic tweezers are  
15 well established single molecule technique for the simultaneous application of torsion and tension to a  
16 single DNA molecule (17,18,19,20,21). By studying resulting variations of the extension of the molecule,  
17 we can explore the consequence of DAP substitution on nanomechanical characteristics of DNA  
18 filament in the different phases resulting in conditions of high negative torsion.  
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24 Overall, we report that the transition to a left-handed form occurred at low levels of unwinding and  
25 tension, and left-handed DAP DNA was significantly stiffer than unmodified L-DNA. The tenfold lower  
26 persistence length of this left-handed form and the higher susceptibility of triply hydrogen bonded  
27 regions to undergo this transition greatly expands the conformational complexity of DNA.  
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## 33 **MATERIAL AND METHODS**

### 34 **Preparation of DAP and WT DNA**

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38 WT and DAP DNA were prepared as previously described (9). Briefly, for the MT experiments T7  
39 DNA ligase was used to ligate a ~1.0 kbp, multiply digoxigenin-labeled (dig-tail) DNA fragment at one  
40 end of a 4642 bp (main) DNA fragment and a ~0.9 kbp, multiply biotin-labeled (bio-tail) DNA fragment  
41 to the other end.  
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### 44 **Magnetic tweezers setup and measurements**

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47 We used a custom made Magnetic Tweezers (MT) setup previously described in (22,23,24,25),  
48 consisting of an inverted optical microscope equipped with an oil-immersion objective (NIKON 100x,  
49 NA = 1.25) mounted on a piezoelectric focusing system (PIFoc, Physik Instrumente, Bresso, Italy). The  
50 objective, coupled with a 15 [cm] focal-length lens, led to a 75x magnification. The magnetic field was  
51 generated by two permanent neodymium magnets placed above the flow chamber and two piezoelectric  
52 motors controlled the position of the magnets along the optical axis (z-direction) and the rotation around  
53 the same axis in order to apply a stretching force, or a torque, to the torsionally constrained DNA.  
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### 58 **Microchamber preparation**

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3 The flow cell consisted in a square glass capillary tube (1x1 mm<sup>2</sup> section, 5-cm long, VitroCom,  
4 Mountain Lakes, NJ). For each measurement, 250 ml of DNA and magnetic beads suspension were  
5 injected, in the absence of a magnetic field, into the capillary previously functionalized as follows. First,  
6 a solution of 100 mg/ml polystyrene (average MW230000, Sigma Aldrich) in toluene was injected into  
7 the capillary. Then, the capillary was drained and dried with compressed air. In this way, the internal  
8 walls were uniformly coated with polystyrene (26). Next, 5 mg of sheep polyclonal anti-digoxigenin  
9 antibody (Roche) in 100 ml with 10 mM PBS were incubated in the capillary for 2 h at 37°C. Unbound  
10 anti-digoxigenin was eliminated by rinsing the capillary with PBS-Tween-20. The functionalized surface  
11 was passivated for 2 h at 37°C with a solution consisting of 10 mM PBS at pH 8 supplemented with  
12 0.1% Tween-20, 1 mg/ml fish sperm DNA (Roche) and 3 mM NaN<sub>3</sub> (27). Finally, the capillary was rinsed  
13 with PBS-Tween-20 and incubated for 1h to allow the DNA to bind to the lower capillary surface. For  
14 storage, several capillaries could be prepared simultaneously and kept at 20°C after air drying them.  
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### 24 **Data acquisition and analysis**

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26 Images were acquired by a CCD camera (Marlin Allied vision, USA) running at a frame rate of 60  
27 Hz and analyzed in real time by a custom made software developed in Java (Oracle, USA). Calculations  
28 on the diffraction rings profile of the magnetic beads allowed to measure the 3D position of the beads  
29 with a precision of 10 nm along the optical axes and 40 nm in x-y plane (22). The extracted data was  
30 recorded on the pc drive for successive analysis that was performed with an ad hoc software written in  
31 MatLab (Mathwork inc. USA).  
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## 38 **RESULTS**

### 39 **Phenomenology of the torsional behavior of WT and DAP DNA.**

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41 In this work, we studied the difference between the nanomechanical behavior of WT and DAP DNA  
42 in conditions of high negative torsion. Representative twist experiments are reported in Figure 1A where  
43 the DNA extension  $L_e$  is measured as a function of the imposed turns,  $n_t$ , or the corresponding  
44 supercoiling,  $\sigma = n_t / (N_b / 10.4)$ , where  $N_b = 4642$  is the number of base pairs of the DNA construct. Figure  
45 1A shows data acquired over a large range of imposed turns (from  $n_t = +100$  until  $n_t = -900$  turns) for WT  
46 and DAP DNA at three different forces  $F$  (0.2, 1.1, 2.3 [pN]). Additional data at different forces are  
47 reported in Figure S1. Whereas at low forces ( $F = 0.2$  [pN]), the behavior of WT and DAP DNA was  
48 basically indistinguishable, at intermediate ( $F = 1.1$  [pN]) and high ( $F = 2.3$  [pN]) forces, the two types of  
49 DNA show  $L_e$  vs  $n_t$  curves having qualitatively and quantitatively different characteristics. In particular,  
50 at intermediate and high forces the extension unwound DNA appeared systematically higher for DAP  
51 DNA than for WT DNA. Furthermore, for high forces ( $F = 2.3$  [pN]) three different linear regimes are  
52 clearly distinguishable.  
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3 These three regimes are highlighted in Figure 1B, a schematic representation of a typical  $L_e$  vs  $n_t$   
4 curve which subdivides it into three linear segments, where  $L_{eB}$ ,  $L_{eBL}$ , and  $L_{eL}$  define the DNA extension  
5 in the B, BL, and L forms (BL indicates a mixture of B and L forms). Each segment is characterized by  
6 a slope,  $dL_e/dn_t$ . For  $n_t > 0$  (red region), the DNA remains in the B-form and  $dL_{eB}/dn_t$  is negative. For  
7 intermediate negative values of  $n_t$  (green region), the DNA extension can either decrease or increase  
8 as a function of  $n_t$ , corresponding to a shallow positive or negative BL slope  $dL_{eBL}/dn_t$ , respectively. For  
9 higher negative values ( $\sigma < -1.8$ ; yellow region), the DNA is in the L-form and  $dL_{eL}/dn_t$  is positive and  
10 steep.  
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15 Following these observations, we first adopted a phenomenological approach to study the B, BL and  
16 L slopes for WT and DAP DNA, respectively, and for various applied force values.  
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19 *DNA B-form.* As shown in Figure 2, the absolute value of the slopes  $dL_{eB}/dn_t$ , calculated for the  
20 region of the experimental  $L_e$  vs  $n_t$  curves where  $n_t > 0$  (B slopes), is a decreasing function of the force.  
21 We note that the DAP slope absolute values, despite the large error bars, appear systematically larger  
22 than the WT DNA slope absolute values.  
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26 *DNA BL-form.* The slopes  $dL_{eBL}/dn_t$  (BL slopes), calculated in the intermediate BL region, are shown  
27 in Figure 3. The figure reveals that at fixed applied force the slopes of DAP DNA are significantly lower  
28 than those of WT DNA. Furthermore, the BL slopes show a change in sign for both DNA types, indicating  
29 an inversion force  $F^*$ , marked by stars in Figure 3, where the BL slope is zero. The inversion force is  
30 higher for WT DNA ( $F^*_{WT} = 2.7 \pm 0.3$  [pN]) than for DAP DNA ( $F^*_{DAP} = 1.6 \pm 0.3$  [pN]).  
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34 *DNA L-form.* Finally, Figure 4 shows the dependence of the L slopes  $dL_{eL}/dn_t$  (L slopes) on the  
35 applied force for both WT and DAP DNA. The values of L slopes are smaller with respect to the  
36 corresponding B slopes. Furthermore the L slopes for DAP appear significantly higher than those for  
37 WT DNA, highlighting, also in this case, the difference in the response to torsion between WT and DAP  
38 DNA.  
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#### 45 **Theoretical frame.**

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47 The variation of the DNA extension under an imposed torsion is due to the need of DNA to relax the  
48 torsion by forming plectonema and/or denaturation bubbles (28). For small applied forces ( $F < 0.5$  [pN]),  
49 the DNA is in B-form and the formation of plectonema is strongly favoured, consequently the variations  
50 in DNA extension as a function of torsion are significant. At negative and intermediate values of  $n_t$  or  $\sigma$   
51 ( $-800 < n_t < -50$ ;  $-1.8 < \sigma < -0.1$ ) and higher values of applied force ( $F > 1$  [pN]), i.e. in the DNA BL-form  
52 regime, the formation of denaturation bubbles induces slight variations in the DNA extension and the  
53 number of base pairs involved in such denaturation bubbles is proportional to  $n_t$  or  $\sigma$  (29). For large  $n_t$   
54 or  $\sigma$  absolute values ( $n_t < -800$ ;  $\sigma < -1.8$ ) the DNA is completely denatured, its conformation changes from  
55 B to L, and any further decrement of  $n_t$  or  $\sigma$  likely induces the formation of plectonema in segments of  
56 the L-form DNA.  
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In the following, we propose a simplified theoretical explanation for understanding the torsional behavior of the three DNA forms discussed thus far in either WT, or DAP, DNA.

*DNA B-form.* The variations of the DNA extension under positive imposed turns, or negative turns at low forces, is due to the formation of plectonema (30,31,32). A precise modeling of the geometrical structure of such plectonema has been already proposed, and it successfully predicts quantitatively the B slope of experimentally measured  $L_e$  vs  $n_t$  curves in WT DNA (33,34,35,36). Those authors extracted the geometrical conditions of the spiral-like plectonema which minimize the total energy by calculating the total energy as the sum of the potential, bending and electrostatic energies (34). The resulting prediction for the B slope according to that model are the dotted red (DAP) and blue (WT) curves in Figure 2, calculated at 150 [mM] NaCl for a WT DNA of persistence length  $L_{pB}=80$  [nm] and  $L_{pB}=40$  [nm], respectively. Such specific  $L_{pB}$  values are the best fit parameters for the proposed model (34) to the data and are qualitatively similar to the already reported values of persistence length of WT and DAP DNA (5,9) (see Discussion).

*DNA BL-form.* By further decreasing  $\sigma$ , especially in DNA molecules under increased tension, the DNA extension is well described by a combination of B and L-DNA (mixed BL phase) (33,37). Standard B-DNA is characterized by a persistence length,  $L_{pB}\approx 50$  [nm], and a distance between base pairs  $L_{0B}=0.34$  [nm]. The same parameters are generally difficult to estimate for L-DNA, but a value of  $L_{pL}\approx 3$  [nm] and a corresponding increment of the extension  $L_{0L}/L_{0B}= 1.41$  have been reported for WT DNA (33). Overall, L-DNA appears to be more flexible and elongated with respect to B-DNA. The percentage of L-DNA induced by torsion increases linearly with  $\sigma$ , until the critical value of  $n_{t,max}\approx -800$ , or  $\sigma_{max}\approx -1.8$ , where all B-DNA is converted to L-DNA. At the two extremes ( $n_t=0$  or  $\sigma=0$  and  $n_{t,max}\approx -800$  or  $\sigma_{max}\approx -1.8$ ), the DNA extension,  $L_{eB}$  and  $L_{eL}$ , respectively, is correctly modelled as a function of the external force  $F$ , by a classical Worm Like Chain (WLC) (38,39,40) with the corresponding  $L_{0B}$ ,  $L_{pB}$  and  $L_{0L}$ ,  $L_{pL}$  for pure B and L-DNA, respectively (41). The relation between  $F$  and both  $L_{eB}$  and  $L_{eL}$  is described by the function  $f$  given by the classical WLC model as follows:

$$F = f(L_{eB}, L_{pB}, L_{0B}) = \frac{k_B T}{L_{pB}} \left( -\frac{1}{4} + \frac{1}{4 \left( 1 - \frac{L_{eB}}{N_b L_{0B}} \right)^2} + \frac{L_{eB}}{N_b L_{0B}} \right) \quad (\text{eq. 1})$$

$$F = f(L_{eL}, L_{pL}, L_{0L}) = \frac{k_B T}{L_{pL}} \left( -\frac{1}{4} + \frac{1}{4 \left( 1 - \frac{L_{eL}}{N_b L_{0L}} \right)^2} + \frac{L_{eL}}{N_b L_{0L}} \right) \quad (\text{eq. 2})$$

From which, it is possible to calculate the DNA extension as  $L_{eB} = f^{-1}(F, L_{pB}, L_{0B})$

and  $L_{eL} = f^{-1}(F, L_{pL}, L_{0L})$ .

At intermediate  $\sigma$  values the DNA extension  $L_{eBL}$  of the BL mixed phase can be described as a linear combination of the WLC contribution from the two pure phases B and L:

$$L_{eBL} = L_{eB}(1 - \chi) + L_{eL}\chi = f^{-1}(F, L_{pB}, L_{0B})(1 - \chi) + f^{-1}(F, L_{pL}, L_{0L})\chi \quad (\text{eq. 3})$$

where  $\chi$  represents the fraction of L-DNA present in the sample and it can be evaluated as

$$\chi = \frac{n_t - n_b}{n_{t,max} - n_b} = \frac{\sigma - \sigma_b}{\sigma_{max} - \sigma_b} \quad (\text{eq. 4})$$

where  $n_b$  or  $\sigma_b$  are the number of turns, or supercoiling, corresponding to the buckling transition, the point at which the first supercoil forms, and  $n_{t,max}$  or  $\sigma_{max}$  are the number of turns, or supercoiling, necessary for a complete conversion to the L-form ( $\sigma_{max} = n_{t,max} / (N_b / 10.4)$ ).

The results of this heuristic model are represented in Figure 5A, where the WLC predictions for  $L_{eBL}$  vs  $F$  are shown in the case of the pure B (orange) and L (cyan) phases. The two limit cases of pure B and pure L phases correspond to the values of  $n_t = 0$  and  $n_t = n_{t,max}$ , respectively. As apparent from the Figure, the WLC model for the B-form is characterized by both a lower value of the asymptotic extension ( $L_{0B} < L_{0L}$ ) and a higher value of the persistence length with respect to the L-form ( $L_{pB} > L_{pL}$ ). The red star (Figure 5A) indicates the force at which extensions of B and L phases are equivalent, i.e., the inversion force.

For intermediate values of  $\sigma$ , this simple model predicts a behavior of  $L_{eBL}$  vs  $F$  described by the magenta curve, obtained for the exemplificative value of  $n_t = n_{t,max} / 2$  (Figure 5A). Furthermore, this model allows calculation of the value of the BL slope as  $dL_{eBL} / dn_t$  (Figures 5B and 5C) taking the  $n_t$  derivative of eq. 3. In all the plots of Figure 5, the theoretical predictions of the model are presented as dashed lines (outside the range of validity) or continuous lines (inside the range of validity), to emphasize that the proposed simplified model is valid only for high force values, where the percentage of denaturation bubbles is higher, and we can disregard the plectonema formation. Given the values of  $L_{0B}$  and  $L_{pB}$ , the free parameters of the model are the persistence length  $L_{pL}$  and the extension increment  $L_{0L}$  of the L-form. Consequently, in Figure 5B the predictions of the BL slopes are obtained by keeping constant the  $L_{0L}$  and considering a variable  $L_{pL}$ , while viceversa in Figure 5C. The main result of the simplified model is the prediction of the inversion force  $F^*$  (stars in Figure 5) where the BL slopes assume zero values for different values of  $L_{0L}$ , and  $L_{pL}$ . These inversion forces have been consistently calculated for several conditions (Figure S2) where we show  $F^*$  obtained by keeping constant the  $L_{0L} / L_{0B}$  ratio and by varying  $L_{pL}$  (Figure S2A), and vice versa (Figure S2B). From the predicted inversion force values it is possible to estimate the values of  $L_{0L} / L_{0B}$  and  $L_{pL}$  compatible with the experimentally measured values of  $F^*_{WT} = 2.7 \pm 0.3$  [pN] and  $F^*_{DAP} = 1.6 \pm 0.3$  [pN] which are indicated by the blue and red continuous horizontal lines in Figure S2 together with their confidence range indicated by horizontal dashed lines. Indeed, by keeping constant the values of  $L_{pB}$  and  $L_{0B}$ , it is possible to calculate a specific value of the inversion force by assuming the values  $L_{pL}$  and  $L_{0L}$ . Vice versa from a specific value of  $F^*$  it is possible to calculate the values of  $L_{pL}$  and  $L_{0L}$  compatible with such force values. Namely, the intersections of the horizontal lines of Figure S2 with the calculated lines at various  $L_{0L} / L_{0B}$  and  $L_{pL}$  indicates the  $L_{pL}$  and  $L_{0L} / L_{0B}$  values compatible with the measured values of  $F^*$ . The results are shown in the inset of Figure 3, where in the

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3 space of the two variables  $L_{0L}/L_{0B}$  and  $L_{pL}$ , we show the locus of values compatible with the range of  
4 measured values of the inversion force, represented as red and blue regions for DAP and WT,  
5 respectively.  
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8 *DNA L-form.* The L slopes are obtained at high negative values of the imposed supercoiling ( $n_t < -$   
9 800 or  $\sigma < -1.8$ ). In such a regime, all the DNA is in the L-form and any further torsion presumably induces  
10 the formation of plectonema of L-DNA. Since the physical phenomenon of plectonema formation is the  
11 same in both cases and the only difference is represented by the specific parameters of the phase  
12 under investigation, it would be possible to apply the same model (34) used to explain the B slopes  
13 also for L slopes. In Figure 4, we reported the experimentally measured values (symbols) and the values  
14 of the L slope predicted by the model (dashed lines) obtained for a specific parameter set:  $L_{pL, DAP} = 3$   
15 [nm] (red line) and  $L_{pL, WT} = 1.5$  [nm] (blue line) and reduced values of linear charge density and DNA  
16 radius (see Discussion).  
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## 25 DISCUSSION

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27 The experimental data presented here show significant analogies and differences between WT and  
28 DAP DNA. In particular, the measurements of the DNA extension under the effect of controlled  
29 mechanical stress, such as imposed torque and stretching, give access to the nanomechanical  
30 characterization of WT and DAP DNA in both the B and L forms. Indeed, using MTs it is possible to  
31 control the fraction of L-form. Such precise control is hardly achievable outside of the frame of a single  
32 molecule experimentation.  
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37 *DNA B-form.* Firstly we note that, in the regime of positive supercoiling and negative supercoiling at  
38 low forces ( $< 0.5$  pN), both the WT and DAP-substituted DNA adopt the classical right-handed double  
39 helix with unperturbed hydrogen bonds between the Watson and Crick bases (DNA in the B-form). In  
40 these regimes, the slope of the dependence of  $L_e$  vs. number of imposed turns,  $n_t$ , is due to the  
41 plectonema geometry, which depends on the B-DNA characteristics such as the persistence length and  
42 on the tension in the molecule (33,34,35,36,42). In the data presented here (Figure 2) we systematically  
43 observe steeper slopes for B-form DAP DNA with respect to WT DNA. The steeper slope of the of  $L_e$   
44 for DAP DNA vs. the number of imposed turns,  $n_t$  at low forces, indicates that the persistence length  
45 of DAP is definitely larger than that of WT DNA. The prediction of the model we used (34 or  
46 <https://home.uni-leipzig.de/mbp/index.php/software/>) quantifies precisely the B slope dependence on  
47 the persistence length, force, and ionic strength. Accordingly, the dashed lines in Figure 2 show the  
48 slopes calculated following the model in ref. (34), assuming a best fit value of persistence length  $L_p = 80$   
49  $\pm 15$  [nm] for DAP and  $40 \pm 10$  [nm] for WT B-DNA. The persistence length of DAP-substituted DNA  
50 obtained from MT measurements has been already reported in literature (5,9) and is compatible with  
51 our best fit values.  
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59 *DNA BL-form.* The negative supercoiling region at higher forces ( $> 1$  pN) exhibits significant  
60 differences between the BL and L slopes of WT and DAP-DNA samples. The difference between the



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3 inversion force  $F^*$  of DAP and WT DNA (Figure 3) is striking. This inversion force difference can be  
4 explained with the simplified model we obtained by combining in series a WLC model of B- and one of  
5 L-form DNA, each with its own characteristic  $L_p$  and  $L_0$  parameters.  
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8 Based on the assumption that the relative fraction  $\chi$  of L-DNA is linearly dependent on the number  
9 of imposed turns, our model can predict persistence length  $L_{pL}$  and contour length  $L_{0L}$  values, which are  
10 compatible with the measured inversion force  $F^*$ .  
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13 Our model is not sufficient to describe in more detail the L-form DNA mechanical parameters.  
14 Indeed, in principle, many pairs of values for  $L_{pL}$  and  $L_{0L}/L_{0B}$  within the areas shown in the inset of Figure  
15 3 satisfy the model. Nevertheless, it is possible to determine independently  $L_{0L}/L_{0B}$  from our data with a  
16 simple geometrical reasoning. The argument is based on the observation that all the DNA phases share  
17 the same length of the backbone and the differences in the contour lengths are related to the helix  
18 geometrical properties. In particular, the DNA contour length is determined by the helical pitch that can  
19 be easily calculated from the threshold ( $n_t = -800$ ) at which the BL region ends. A simple calculation (see  
20 Figure S3 for details) leads to a value  $L_{0L}/L_{0B} = 1.4$  equal for both WT and DAP, since both have the  
21 same threshold in  $n_t = -800$  for the BL to L transition. This value of 1.4 is also reported in the literature  
22 only for WT DNA (33,41). By using this value for  $L_{0L}/L_{0B}$ , we obtained the persistence length values  
23  $L_{pL,WT} \approx 3.5$  [nm] and  $L_{pL,DAP} \approx 7.0$  [nm] for the L-form.  
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31 The resulting increment of the L-DNA persistence length of DAP with respect to WT DNA suggests  
32 a more structured and rigid form of the L phase in DAP DNA. This increment of DAP stiffness confirms  
33 the qualitative observation that DAP is more easily extended with respect to the WT at constant force  
34 (see Figure 1A). Unlike B-DNA, the helically intertwined but non-hybridized L-DNA form is not well  
35 defined. The L-form is related to the torsional stress that drives the two strands of DNA counterclockwise  
36 while the electrostatic repulsion of the backbone strands gives rigidity to the chain. Some DNA bases  
37 might experience large enough negative twist to allow base pairing in the left hand configuration.  
38 However, in this case, Z-DNA would result, which is characterized by different mechanical parameters.  
39 In particular, for WT Z-DNA has a persistence length,  $L_{pZ,WT} = 200$  nm, much higher than that of L-DNA  
40 (33,43). In the BL region we cannot exclude that a small fraction of Z-DNA phase might be present  
41 along the chain. Differing tendencies to generate Z-DNA in the BL region could explain the observed  
42 difference of  $L_{pL}$  determined for DAP DNA.  
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50 *DNA L-form.* Finally, the measurements illustrated in Figure 4 shows a steeper dependence of  $L_{eL}$   
51 on  $n_t$  for DAP DNA in the region where all the DNA chain is completely converted to L-DNA. This  
52 observation is qualitatively confirmed by the model we adopted (34), since in this region the slope is  
53 presumably due to the formation of plectonemes and an increment can be justified with an increase in  
54 the persistence length of the L-form. The difference in persistence length qualitatively explains the  
55 difference between  $dL_{eB}/dn_t$  and  $dL_{eL}/dn_t$ , together with the difference between  $dL_{eL,WT}/dn_t$  and  
56  $dL_{eL,DAP}/dn_t$ . Unfortunately, the model used to describe the plectonemes formation for the B-DNA (34)  
57 can grasp only qualitatively the behaviour of L-DNA. Indeed, the model fits the data well (see dashed  
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lines of Figure 4) imposing  $L_{pL,WT} \approx 1.5$  [nm] and  $L_{pL,DAP} \approx 3$  [nm] which are different from what we derived above from the analysis of BL slopes ( $L_{pL,WT} \approx 3.5$  [nm]) and what reported in the literature ( $L_{pL,WT} \approx 3$  [nm]) (33). Furthermore, to obtain a reasonable fit with the model (34) is also necessary to impose a reduction of a factor of two of the filament linear charge density and of the DNA diameter, which would result in unrealistic parameters. We note, however, that the differences between the data and the model predictions (34) suggest that the model was extended beyond its range of validity. Assuming that the L slope simply depends on the formation of plectonema is likely naive, since as reported in (33) for the WT DNA, in the L regime, the torque does not seem constant, as should be expected in a plectonemic regime. A possible way to reconcile all these clues is to assume the presence of a new DNA phase after the complete denaturation process. However, whatever the conformation of DNA in this region might be, the large difference between the L slopes attests to a significant difference in the mechanical parameters of DAP with respect to WT DNA in this extremely negatively coiled regime.

In conclusion, L-DNA is the product of enzymes which unwind and/or denature the double helix and a target of regulatory proteins that selectively bind unwound or left-handed DNA. Therefore, given the low levels of unwinding and tension at which it forms, L-DNA likely plays a role in a wide range of DNA transactions. The higher bending rigidity of DAP DNA compared to normal DNA, imposed larger plectonemic gyres at high levels of unwinding. These features emphasize the dependence of the mechanics of L-form DNA on the extent of triple hydrogen bonding and/or on modifications that alter the minor groove and the associated stabilizing spine of water molecules (44) or the major groove in which methylation makes DNA more susceptible to B-Z transitions (45). More generally, the different mechanical properties of DAP expand the tool kit for the rational design of biomaterials.

#### **DATA AVAILABILITY**

Not applicable.

#### **SUPPLEMENTARY DATA**

Supplementary Data are available at NAR online.

#### **ACKNOWLEDGEMENT**

We thank C.A. Marrano for help with the measurements and R. Seidel for discussions about the theoretical model.

#### **FUNDING**

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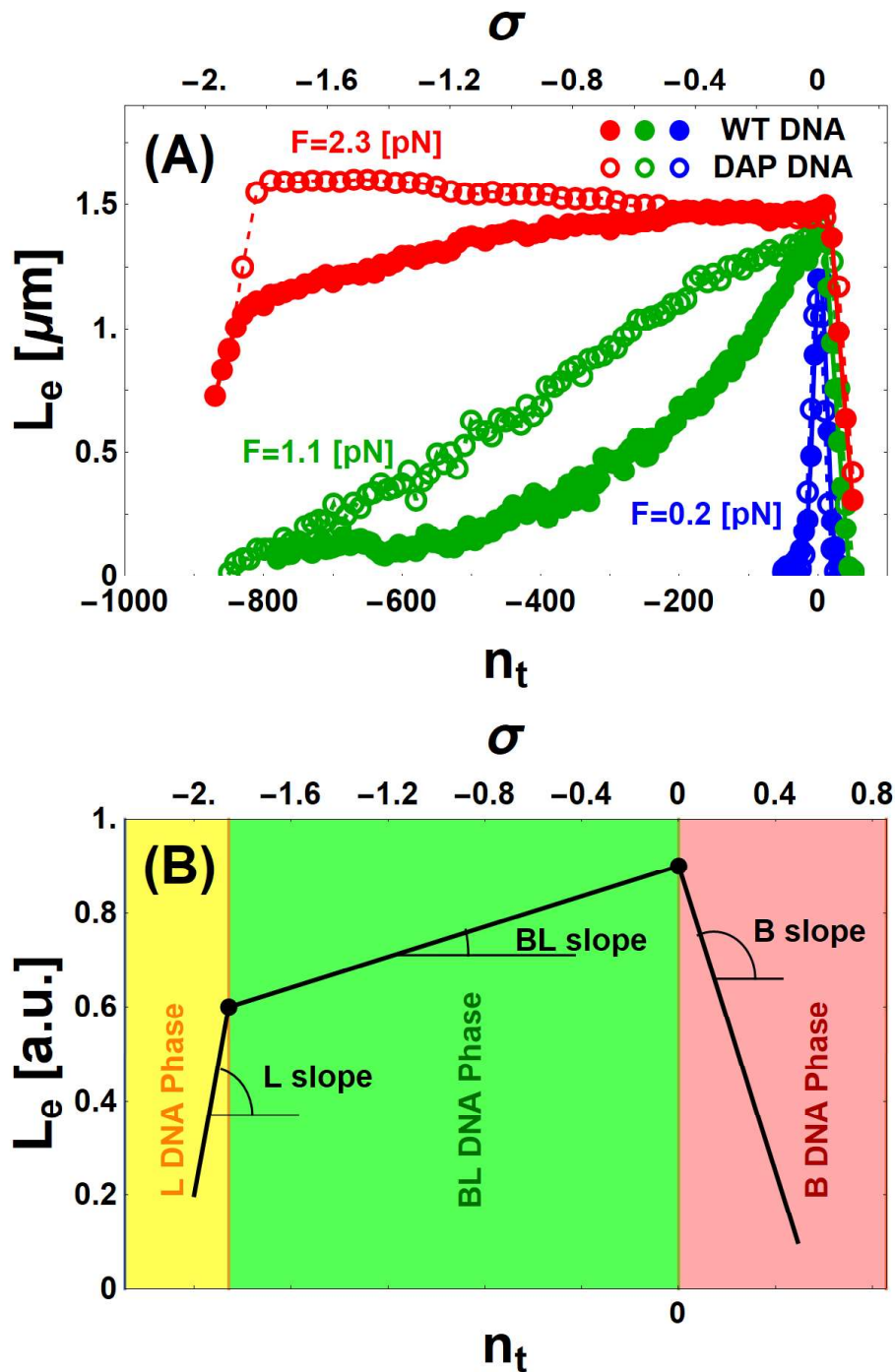
Funding for open access charge: shared between University of Milano-Bicocca and National Institutes of Health [R01GM084070].

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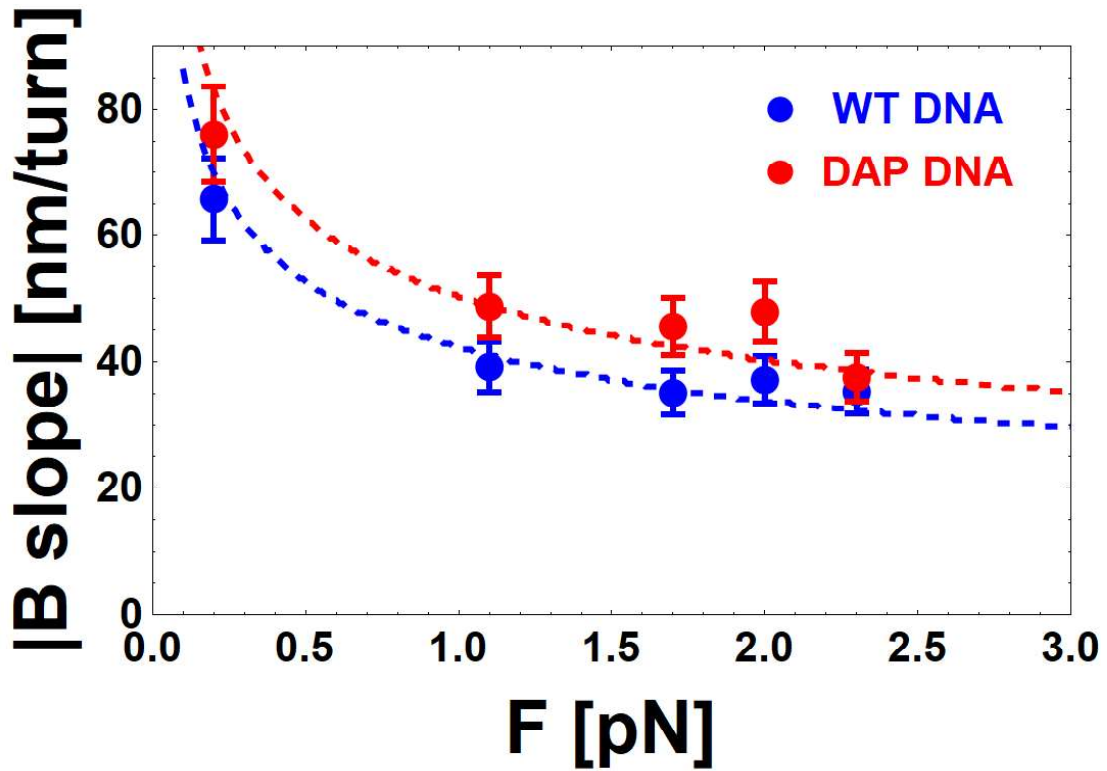
**CONFLICT OF INTEREST**

None declared.

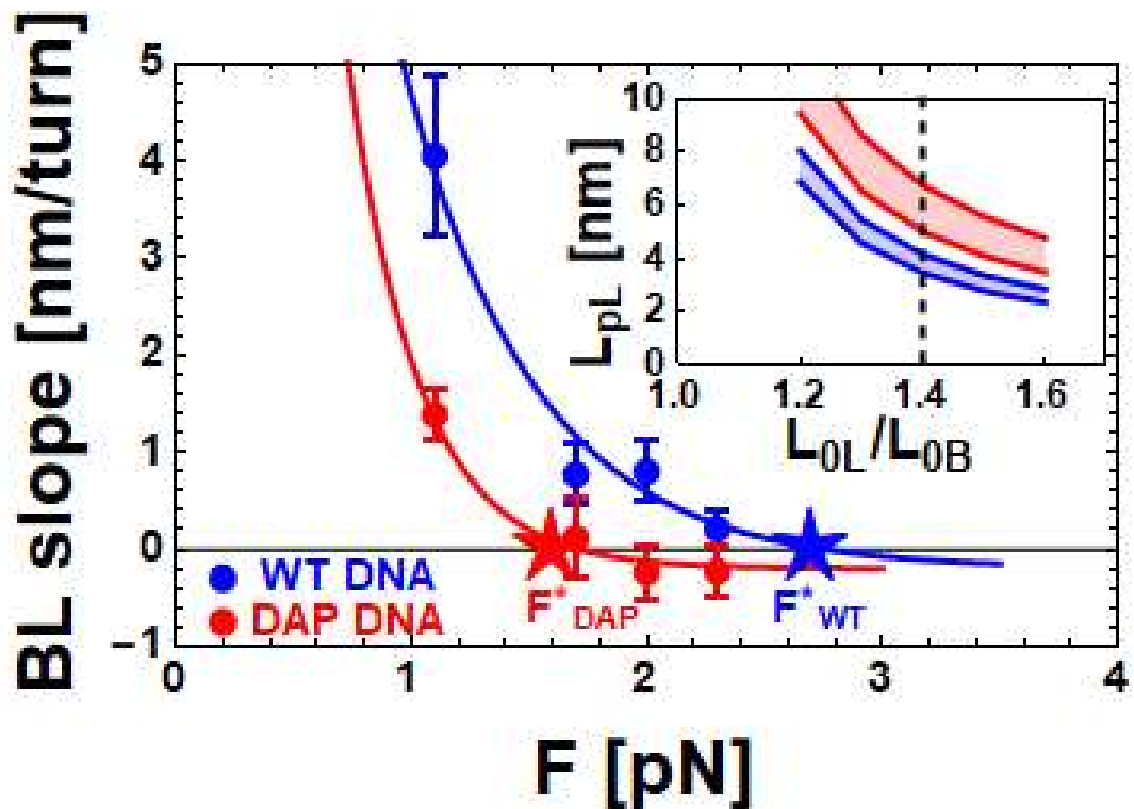
## TABLE AND FIGURES LEGENDS



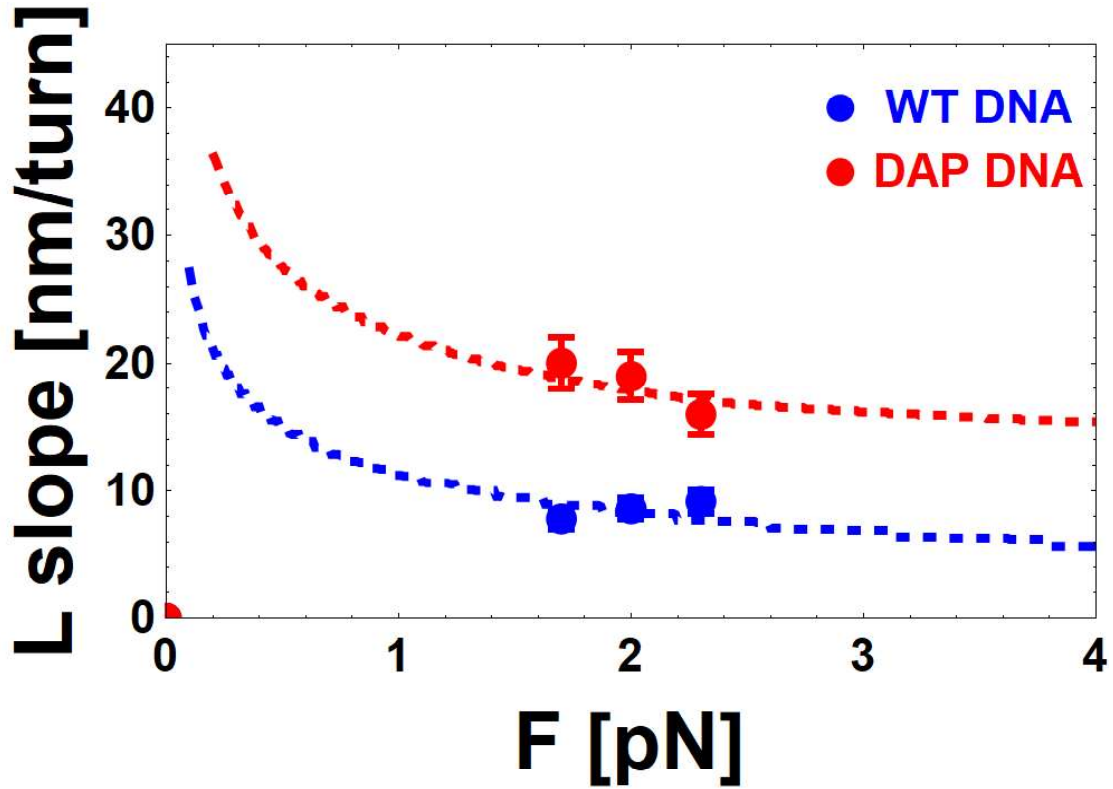
**Figure 1.** Torsional behavior of WT and DAP DNA. A: Representative MT data for WT (filled disks) and for DAP (open circles) DNA display different extensions,  $L_e$ , as a function of the number of imposed turns  $n_t$ , (lower axis) or supercoiling density  $\sigma$ , (upper axis) at fixed values of force  $F$ . B: Simplified sketch of theoretical DNA extension,  $L_e$ , obtained as a function of the number of imposed turns  $n_t$ , or supercoiling density  $\sigma$ , at fixed values of  $F$  shows three phases. The different, color-coded DNA phases are: L (yellow), mixed BL (green), and B (red). The specified angles illustrate the different slopes  $dL_e/dn_t$  characterizing the L, BL, and B phases.



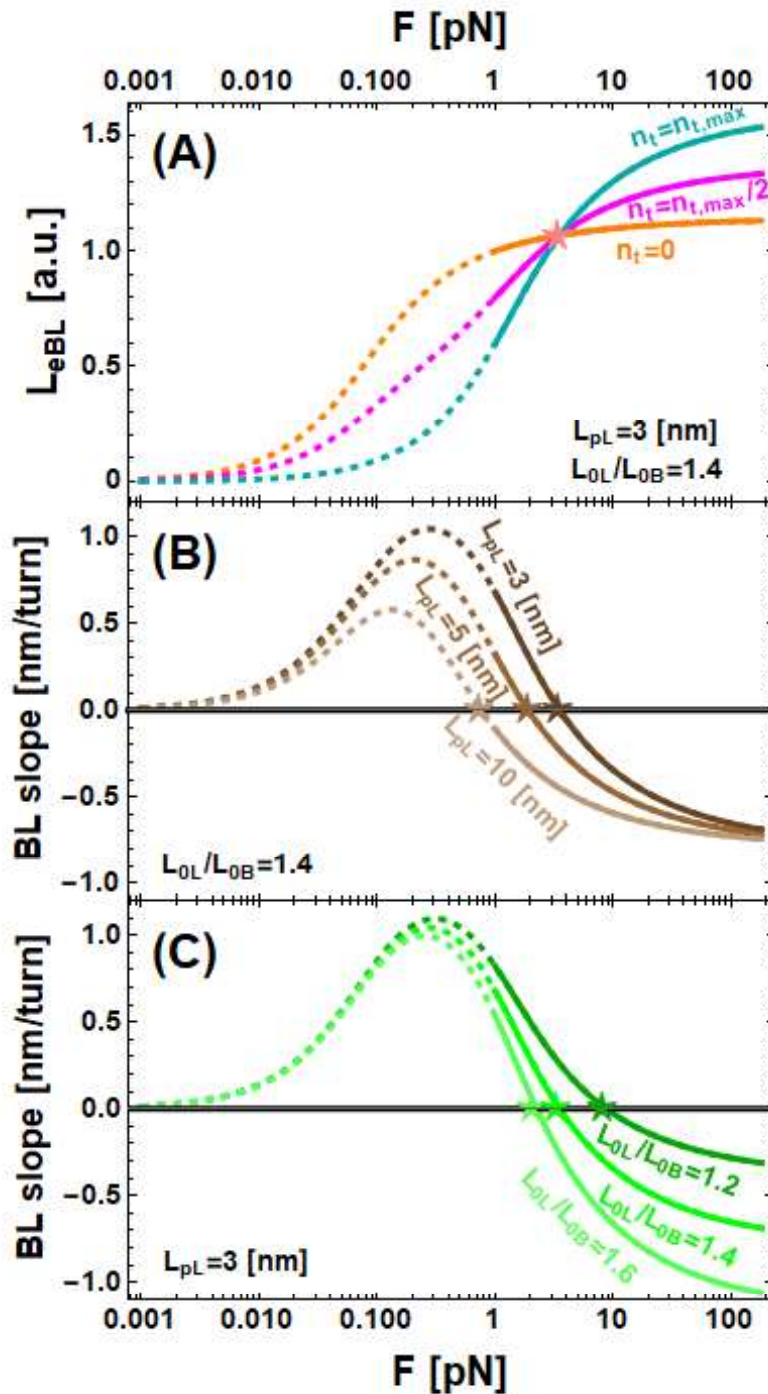
**Figure 2.** Force dependence of the slope of the extension versus twist data for B-form WT and DAP DNA. Experimental values (dots) and theoretical predictions (dashed curves) for the slope  $dL_{eB}/dn_t$  of B-WT (blue) and B-DAP (red) DNA. Slopes are given in absolute value. Predictions were obtained using the model (22) described in the text assuming:  $I_s=150$  [mM] and the best fit values  $L_{pB}=40$  [nm] (blue);  $L_{pB}=80$  [nm] (red).



**Figure 3.** Force dependence of the slope of the BL phase for WT and DAP DNA. Experimental data (dots) and theoretical predictions (solid lines) for the slope  $dL_{eBL}/dn_t$  of the BL-DNA phase are shown as a function of the applied force  $F$  for WT (blue) and DAP (red) DNA. The continuous lines are a guide for the eye.  $F^*_{DAP}=1.6$  [pN] and  $F^*_{WT}=2.7$  [pN] (stars), represent the measured inversion force for DAP and WT DNA, respectively. Inset: theoretical predictions of the regions of the values of  $L_{pL}$  and  $L_{0L}/L_{0B}$  compatible with the measured values of the inversion forces (red region for  $F^*_{DAP} \pm 0.3$  [pN] and blue region for  $F^*_{WT} \pm 0.3$  [pN], respectively). The vertical dashed line represents the assumed values of  $L_{0L}/L_{0B}$  for WT and DAP (see text for details).



**Figure 4.** Force dependence of the experimental data (symbols) and the theoretical predictions (dashed lines) for the slopes  $dL_{eL}/dn_t$  of the L-phase for WT (blue) and DAP (red) DNA. Parameters used in the model (22) are as described in the text: ionic strength  $I_s=150$  [mM],  $L_{pL}=3$  [nm] (red line), and  $L_{pL}=1.5$  [nm] (blue line) and reduced linear charge density and DNA radius (see text for details).



**Figure 5.** Simulations of DNA extension  $L_{eBL}$  and BL-phase slope,  $dL_{eBL}/dn_t$ , calculated according to the model described in the text. The theoretical predictions of the model are presented as dashed lines (outside the range of validity) or continuous lines (inside the range of validity). A: The calculated DNA extension  $L_{eBL}$  is plotted as a function of the applied force  $F$  with fixed persistence length of the L phase,  $L_{pL}=3$  [nm], and ratio between the L-form DNA and B-form DNA extension,  $L_{OL}/L_{OB}=1.41$ , for various  $n_t$  ( $n_t=0$  (orange);  $n_t=n_{t,max}/2$  (magenta);  $n_t=n_{t,max}$  (cyan)). B: The calculated BL-phase slope is plotted as a function of force with fixed  $L_{OL}/L_{OB}=1.41$  at  $L_{pL}=3$  [nm];  $L_{pL}=5$  [nm];  $L_{pL}=10$  [nm]. C: The calculated BL-phase slope is plotted as a function of force with fixed  $L_{pL}=3$  [nm] at  $L_{OL}/L_{OB}=1.2$ ;  $L_{OL}/L_{OB}=1.4$ ;  $L_{OL}/L_{OB}=1.6$ .



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4 **Nanomechanics of negatively supercoiled diaminopurine-**  
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10 Domenico Salerno<sup>1</sup>, Francesco Mantegazza<sup>1,\*</sup> Valeria Cassina<sup>1</sup>, Matteo Cristofalo<sup>1</sup>, Qing Shao<sup>2,†</sup>, Laura  
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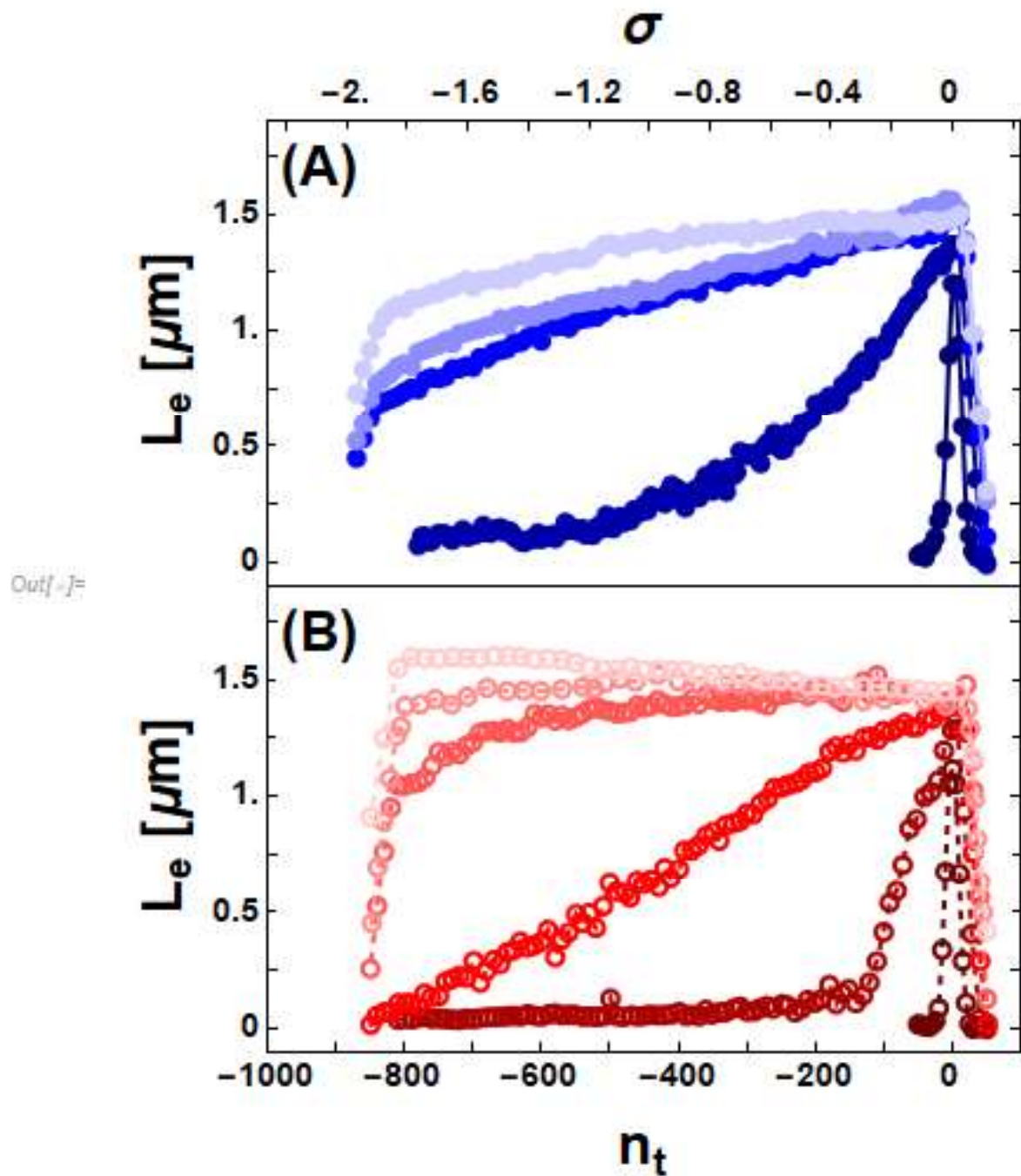
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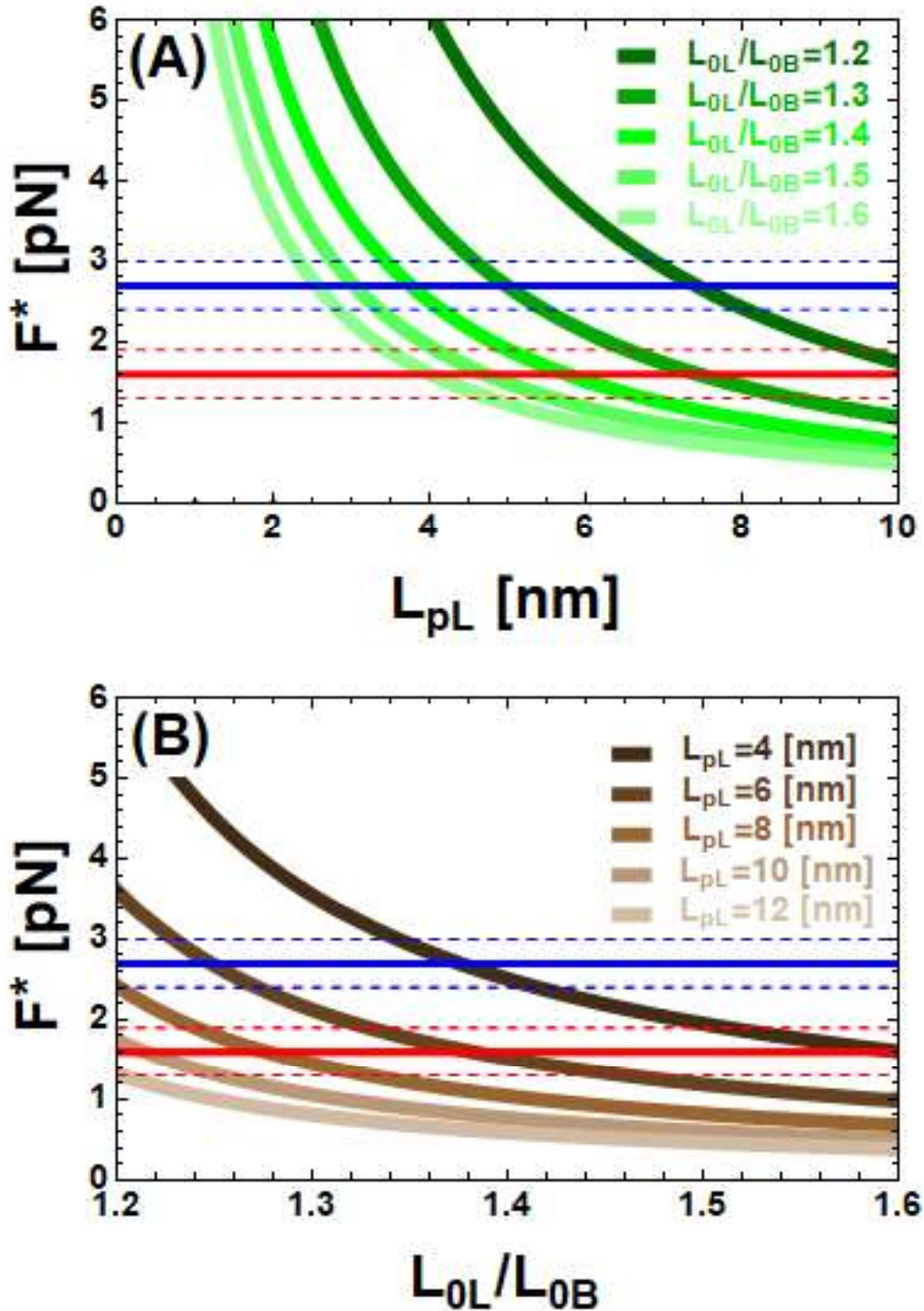
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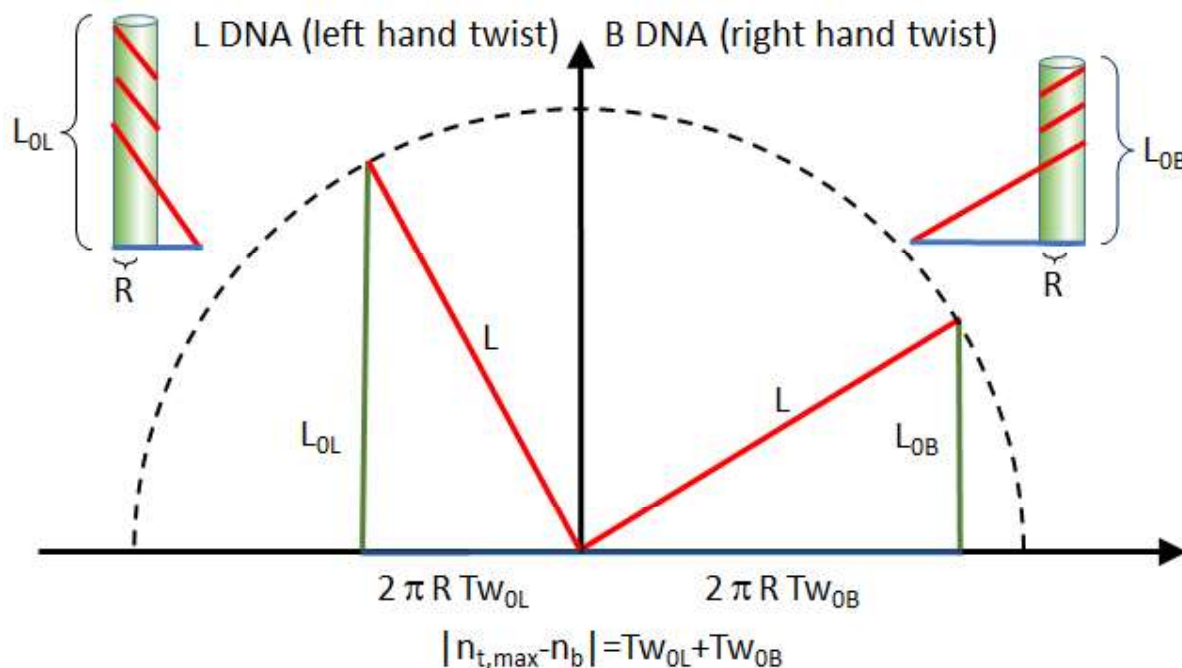
**SUPPLEMENTARY DATA**



**Figure S1.** Torsional behavior of WT and DAP DNA. The measured DNA extension,  $L_e$ , is plotted as a function of the number of imposed turns,  $n_t$ , or equivalently supercoiling density,  $\sigma$ , for WT (A, filled circles) and DAP DNA (B, open circles) at tensions of 0.2, 1.1, 1.7, 2.0, 2.3 [pN] (light to dark shades of blue or red).



**Figure S2.** Theoretical values of the inversion force  $F^*$  calculated as a function of the persistence length  $L_{pL}$  of the L phase and ratio  $L_{0L}/L_{0B}$  between the DNA L and B extension. The continuous horizontal lines represent the measured values of the inversion forces  $F^*$  for WT (blue line) and DAP (red line) DNA. The dashed lines show the uncertainty of the measured forces ( $F^*_{DAP} = 1.6 \pm 0.3$  [pN] and  $F^*_{WT} = 2.7 \pm 0.3$  [pN]).  $F^*$  was calculated as a function of  $L_{pL}$  for  $L_{0L}/L_{0B} = 1.2, 1.3, 1.4, 1.5, 1.6$ , (A) or as a function of  $L_{0L}/L_{0B}$  with  $L_{pL} = 4, 6, 8, 10, 12$  [nm] (B).



**Figure S3.** A geometrical model for calculating the contour length increment from B- to L-form DNA. Unwrapping a double-stranded filament, we obtain a triangle (top right for B DNA) and top left for L DNA). The hypotenuse (in red) corresponds to the length of the totally unwrapped filament  $L=0.6$  [nm] per bp (indicated here as the length of one of the two single strands) which is constant and fixes the locus of the points to which the triangle corner is allowed to belong (dashed circumference). The vertical cathetus (in green) of the triangle corresponds to the contour length of the double stranded DNA,  $L_{OL}$  or  $L_{OB}$ . The horizontal cathetus (in blue) is determined by the Pythagorean Theorem, and it corresponds to the circumference of the DNA ( $2\pi R$ , where  $R\approx 0.8$  [nm] is the DNA radius) multiplied the twist of the DNA in the different forms  $TW_{OB}=1$  turn/10.4 bps and  $TW_{OL}\approx 1$  turn/14.1 bps. Thus, with constant  $L$ , the twist of the DNA form determines its contour length. The relation  $|n_{t,max} - n_b| = TW_{OL} + TW_{OB}$  imposes the topological constraint on the twist number: the total twist applied to the DNA ( $n_{t,max}$ ) is absorbed as twist by the chain ( $n_b$ ) which converts the natural twist of the molecule from  $TW_{OB}$  to  $TW_{OL}$ . It follows that since WT and DAP DNA share the same threshold  $n_{t,max}$ , they also share the same  $TW_{OL}$  and, consequently, the same contour length of the L phase for DAP and WT DNA ( $L_{OL,WT}=L_{OL,DAP}$ ).