

1 **Short-Term Synaptic Dynamics Control the Activity Phase of**
2 **Neurons in an Oscillatory Network**

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4 **Abbreviated Title:** Synaptic dynamics influence activity phase of oscillatory neurons

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19 **Abstract**

20 In oscillatory systems, neuronal activity phase is often independent of network
21 frequency. Such phase maintenance requires adjustment of synaptic input with
22 network frequency, a relationship that we explored using the crab, *Cancer borealis*,
23 pyloric network. The burst phase of pyloric neurons is relatively constant despite a >2-
24 fold variation in network frequency. We used noise input to characterize how input
25 shape influences burst delay of a pyloric neuron, and then used dynamic clamp to
26 examine how burst phase depends on the period, amplitude, duration, and shape of
27 rhythmic synaptic input. Phase constancy across a range of periods required a
28 proportional increase of synaptic duration with period. However, phase maintenance
29 was also promoted by an increase of amplitude and peak phase of synaptic input with
30 period. Mathematical analysis shows how short-term synaptic plasticity can
31 coordinately change amplitude and peak phase to maximize the range of periods over
32 which phase constancy is achieved.

33 150/150

34 **Introduction**

35 Oscillatory neural activity is often organized into different phases across groups
36 of neurons, both in brain rhythms associated with cognitive tasks or behavioral states
37 (Hasselmo et al., 2002; Buzsaki and Wang, 2012; Buzsaki and Tingley, 2018), and in
38 central pattern generating (CPG) circuits that drive rhythmic motor behaviors (Marder
39 and Bucher, 2001; Marder et al., 2005; Grillner, 2006; Bucher et al., 2015; Katz, 2016;
40 Stein, 2018). The functional significance of different phases in the latter is readily
41 apparent, as they for example provide alternating flexion and extension of limb joints,
42 and coordination of movements between joints, limbs, and segments (Krantz and
43 Parks, 2012; Grillner and El Manira, 2015; Kiehn, 2016; Le Gal et al., 2017; Bidaye et
44 al., 2018). A hallmark of many such patterns is that the relative timing between neurons
45 is well maintained over a range of rhythm frequencies (Dicaprio et al., 1997; Hooper,
46 1997b, a; Wenning et al., 2004; Marder et al., 2005; Grillner, 2006; Mullins et al., 2011;
47 Le Gal et al., 2017). If the latency across different groups of neurons changes
48 proportionally to the rhythm period, phase (latency over period) is invariant, in some
49 cases providing optimal coordination at all speeds (Zhang et al., 2014).

50 The ability to maintain phase arises from central coordinating mechanisms
51 between circuit elements, as it is present in isolated nervous system preparations, but
52 the underlying cellular and circuit mechanisms are not well understood. Constant
53 phase lags between neighboring segments in the control of swimming can be
54 explained mathematically on the basis of asymmetrically weakly coupled oscillators,
55 but the role of intrinsic and synaptic dynamics within each segment is unknown (Cohen
56 et al., 1992; Skinner and Mulloney, 1998; Grillner, 2006; Mullins et al., 2011; Zhang et
57 al., 2014; Le Gal et al., 2017).

58 The pyloric circuit of the crustacean stomatogastric ganglion (STG) has inspired
59 a series of experimental and theoretical studies of cellular and synaptic mechanisms
60 underlying phase maintenance. The pyloric circuit generates a triphasic motor pattern
61 with stable phase relationships over a wide range of periods (Eisen and Marder, 1984;
62 Hooper, 1997b, a; Bucher et al., 2005; Goillard et al., 2009; Tang et al., 2012; Soofi et
63 al., 2014). Follower neurons burst in rebound from inhibition from pacemaker neurons
64 (Marder and Bucher, 2007; Daur et al., 2016), and post-inhibitory rebound delay scales
65 with the period of hyperpolarizing currents (Hooper, 1998). Voltage-gated
66 conductances slow enough for cumulative activation across cycles could promote such
67 phase maintenance (Hooper et al., 2009). Similarly, short-term depression of graded
68 inhibitory synapses is slow enough to accumulate over several pyloric cycles, meaning
69 that effective synaptic strength increases with increasing cycle period (Manor et al.,
70 1997; Nadim and Manor, 2000).

71 Theoretical studies have shown that short-term synaptic depression, by
72 increasing inhibition strength with cycle period, should promote phase maintenance
73 (Manor et al., 2003; Mouser et al., 2008), particularly in conjunction with inactivating (A-
74 type) potassium currents (Bose et al., 2004; Greenberg and Manor, 2005), which
75 control the rebound delay (Harris-Warrick et al., 1995b; Harris-Warrick et al., 1995a;
76 Kloppenburg et al., 1999). These predictions remain experimentally untested.

77 Additionally, postsynaptic responses also depend on the actual trajectory of
78 synaptic conductances, which are shaped by presynaptic voltage trajectories and
79 short-term synaptic plasticity (Manor et al., 1997; Mamiya et al., 2003; Zhao et al.,
80 2011; Tseng et al., 2014). If amplitude, duration, and trajectory of synaptic
81 conductance determine rebound delay, phase maintenance necessitates all three of
82 these parameters to change with cycle period in coordination. We used the dynamic

83 clamp technique to exhaustively explore the range of these parameters and
84 understand how the coordinated changes in synaptic dynamics determines the phase
85 of follower neurons in an oscillatory circuit. Our findings are consistent with a
86 mathematical framework that accounts for the frequency dependence of amplitude and
87 peak phase of the synaptic conductance.

88 **Results**

89 *Phase maintenance and latency maintenance*

90 The firing of neurons in oscillatory networks is shaped by a periodic synaptic
91 input. The relative firing latency of such neurons is often measured relative to a defined
92 reference time in each cycle of oscillation, and is used to determine the activity phase
93 of the neuron (see, e.g., Belluscio et al., 2012). For example, in a simple network
94 consisting of a bursting oscillatory neuron driving a follower neuron (Fig. 1A1), at a
95 descriptive level, the latency (L) of the follower neuron activity relative to the onset of
96 the oscillator's burst onset may depend on the oscillation cycle period (P). In response
97 to a change in period (say, to P_2), the follower neuron may keep constant latency ($L_2 =$
98 L), or constant phase, i.e., modify its latency proportionally to the change in period
99 ($L_2/P_2 = L/P$; Fig. 1A2). However, in many oscillatory systems, for example the pyloric
100 circuit (Hooper, 1997b, a), the relationship between L and P falls between these two
101 extremes. A sample recording of the bursting activity of the lateral pyloric (LP) neuron
102 in response to controlled changes in the pyloric cycle frequency demonstrates such a
103 relationship (Fig. 1B and 1C).

104 *The burst onset time of the isolated LP neuron depends on the temporal dynamics of 105 its input*

106 The LP neuron does not have intrinsic oscillatory properties, but oscillates due
107 to the synaptic input it receives from the pacemaker anterior burster (AB) and pyloric
108 dilator (PD) neurons, and the follower pyloric constrictor (PY) neurons (Fig. 2A). The
109 burst onset phase of the LP neuron ($\phi_{LP} = \Delta t / P$; Fig. 2A) is shaped by the interaction
110 between synaptic inputs and the neuron's intrinsic dynamics that influence post-
111 inhibitory rebound. We measured an overall burst onset phase of the LP neuron of ϕ_{LP}
112 $= 0.34 \pm 0.03$ (N=9).

113 As a first-order quantification, we measured how inputs to the LP neuron
114 interact with its intrinsic properties to determine the timing between its bursts, in the

115 absence of network oscillations. To this end, we synaptically isolated the LP neuron
116 (Fig. 2B) and drove it with a noise current input (see Methods). In response to the
117 noise input, the LP neuron produced an irregular pattern of spike times, which included
118 a variety of bursting patterns with different spike numbers (Fig. 2C). We were
119 interested in the characteristics of inputs producing different burst onset latencies.
120 However, unlike a periodic input, noise input does not provide a well-defined reference
121 point to measure the burst onset latency. We therefore categorized bursts with respect
122 to the preceding inter-burst intervals (IBIs) during which no other action potentials
123 occurred. We classified these IBIs in bins (300, 500, 700 and 900 ms) and tagged
124 bursts based on the IBI values (Fig. 2C). We characterized the driving input leading to
125 bursts with specific IBIs by burst-triggered averaging the input current (I_{BTA} ; an example
126 shown in Fig. 2D). Our analysis produced a single I_{BTA} for each of the four IBIs in each
127 preparation (N=23). I_{BTA} 's of each preparation were first normalized in amplitude by the
128 I_{BTA} amplitude at IBI = 300 ms (Fig. 2E; average shown in Fig. 2F) to examine how
129 peak amplitude (I_{peak}) varied with IBI. These data were then normalized in time (Fig.
130 2G) to examine the effect of IBI on peak phase (Δ_{peak}) and the rise ($slope_{up}$) and fall
131 ($slope_{down}$) slopes of the input current across preparations. We found that IBI had a
132 significant effect on I_{peak} , Δ_{peak} , $slope_{up}$ and $slope_{down}$ (all one-way RM-ANOVA on
133 ranks; data included in Figure 2-source data). In particular, larger IBIs corresponded to
134 larger I_{peak} values (Fig. 2F-2H; $p < 0.001$, $\chi^2 = 65.87$) with smaller (more advanced) Δ_{peak}
135 (Fig. 2I; $p < 0.001$, $\chi^2 = 41.35$). The change in Δ_{peak} was due to a decrease in $slope_{up}$
136 ($p < 0.001$, $\chi^2 = 65.25$), whereas $slope_{down}$ did not vary as much (Figs. 2J-2K; $p = 0.002$, χ^2
137 = 14.77).

138 *The burst onset phase of the LP neuron oscillation depends on its synaptic input*

139 Injection of noise current revealed that the timing of the LP response is
140 exquisitely sensitive to the duration and amplitude of inputs. In the intact system, the
141 primary determinant of input duration and amplitude is the network period (P), as
142 increasing P increases both presynaptic pacemaker burst duration (Hooper, 1997b, a)
143 and synaptic strength (Manor et al., 1997; Nadim and Manor, 2000). To explore the
144 effect of the duration and strength of the synaptic input, we used dynamic clamp to
145 drive the LP neuron with a realistic synaptic conductance waveform.

146 We constructed this realistic waveform by measuring the synaptic current input
147 to the LP neuron during ongoing pyloric oscillations (Fig. 3A). These measurements

148 showed the two components of inhibitory synaptic input, those from the pacemaker AB
149 and PD neurons (left arrow) and those from the follower PY neurons (right arrow). In
150 each cycle, the synaptic current always had a single peak, but the amplitude and
151 phase of this peak showed variability across preparations (Fig. 3B, average in blue).

152 The realistic conductance input was injected periodically with strength g_{max} (Fig.
153 3C). For any fixed g_{max} , φ_{LP} decreased as a function of P (Fig. 3D), i.e., the relative
154 onset of the LP burst was advanced in slower rhythms. In contrast to the effect of P , for
155 any given P , φ_{LP} increased sublinearly as a function of g_{max} (Fig. 3E). Fig. 3F combines
156 the simultaneous influence of both parameters on φ_{LP} . The results shown in Fig. 3D
157 indicate that the LP neuron intrinsic properties alone do not produce phase constancy.
158 However, level sets of φ_{LP} (highlighted for three values in Fig. 3F), indicate that phase
159 could be maintained over a range of P values, if g_{max} increases as a function of P . This
160 finding was predicted by our previous modeling work, in which we suggested that
161 short-term synaptic depression promotes phase constancy by increasing synaptic
162 strength as a function of P (Manor et al., 2003; Bose et al., 2004). We will further
163 discuss the role of synaptic depression below.

164 To clarify the results of Fig. 3, it is worth examining the extent of phase
165 maintenance for fixed g_{max} . An example of this is shown in Fig. 4A (turquoise plots). A
166 comparison of these data with the theoretical cases in which either delay or phase is
167 constant suggests that the LP neuron produces relatively good phase maintenance, at
168 least much better in comparison with constant delay. However, this conclusion is
169 misleading because, in these experiments, the duty cycle of the synaptic input was
170 kept constant. Therefore, most of the phase maintenance is due the fact that the
171 synaptic input keeps perfect phase. In fact, if the reference point measures phase
172 relative to the end –rather than onset– of the PD burst (Fig. 4B), phase maintenance of
173 the LP neuron is barely better than in the constant delay case (Fig. 4A, purple plots). It
174 is therefore clear that phase maintenance by the LP neuron would require the
175 properties of the synaptic input to change as a function of P , a hallmark of short-term
176 synaptic plasticity (Fortune and Rose, 2001; Grande and Spain, 2005). As mentioned
177 above, short-term plasticity such as depression could produce changes in g_{max} as a
178 function of P . Independently of g_{max} , the peak time of the synaptic current is another
179 parameter that could change with P and influence the timing of the postsynaptic burst.

180 We therefore proceeded to systematically explore the influence of P , g_{max} and the
181 synaptic peak time on φ_{LP} .

182 *A systematic exploration of synaptic input parameters on the phase of the LP neuron*

183 For a detailed exploration of the influence of the synaptic input on φ_{LP} , we
184 approximated the trajectory of the (unitary) synaptic conductance in one cycle by a
185 simple triangle (Fig. 5A), which could be defined by three parameters: duration (T_{act}),
186 peak time (t_{peak}) and amplitude (g_{max}) (Fig. 5B). This simplified triangular synaptic
187 conductance waveform could then be repeated with any period (P) to mimic the
188 realistic synaptic input to the LP neuron. For a given synaptic duration T_{act} , the peak
189 phase of the synapse can be defined as $\Delta_{peak} = t_{peak} / T_{act}$. The parameter Δ_{peak} is
190 known to vary as a function of P (Tseng et al., 2014) and, in a previous study, we found
191 that Δ_{peak} may influence the activity of the postsynaptic neuron, independent of P and
192 g_{max} (Mamiya and Nadim, 2004). We therefore systematically explored the influence of
193 three parameters of the synaptic input (P , g_{max} and Δ_{peak}) on φ_{LP} .

194 As with the realistic synaptic waveforms (Fig. 3), we used the dynamic clamp
195 technique to apply the triangular conductance waveform periodically to the synaptically
196 isolated LP neuron. Across different runs within the same experiment, the parameters
197 P , g_{max} and Δ_{peak} were changed on a grid (see Methods). In addition, all combinations
198 of these three parameter values were run in two conditions in the same experiment, 1:
199 with constant duration, i.e., constant T_{act} across different P values (C-Dur of 300 ms),
200 and 2: with constant duty cycle, i.e., T_{act} changing proportionally to P (C-DC of 0.3; Fig.
201 5C). Using these protocols, we measured the effects of synaptic parameters on φ_{LP}
202 (Fig. 5D).

203 The LP neuron produced burst responses that followed the synaptic input in a
204 1:1 manner across all values of P that were used (Fig. 6A1). When g_{max} and Δ_{peak} were
205 kept constant, φ_{LP} decreased as a function of P (Fig. 6A2). This decrease was always
206 larger for the C-Dur case than the C-DC case. For both C-DC and C-Dur, this trend
207 was seen across all values of Δ_{peak} and g_{max} (Fig 6A3). The effect of P on φ_{LP} was
208 highly significant for both C-DC (Three-Way ANOVA, $p < 0.001$, $F = 100.677$) and C-Dur
209 (Three-Way ANOVA, $p < 0.001$, $F = 466.424$), indicating that the period and duration of
210 the inhibitory input to the LP neuron had a significant effect on its phase.

211 Changing g_{max} produced a large effect on the level of hyperpolarization in the
212 LP neuron, but this usually translated to only a small or modest effect on the time to the
213 first spike following inhibition (Fig. 6B1). Overall, increasing g_{max} at constant values of P
214 and Δ_{peak} produced a significant but only small to moderate increase in φ_{LP} (Three-Way
215 ANOVA, $p < 0.001$, $F = 10.798$). Although increasing g_{max} produced the same qualitative
216 effect for both the C-DC and C-Dur (e.g., Fig. 6B2), φ_{LP} in the C-DC case was
217 restricted to a smaller range (Fig. 5F top vs. bottom panels). Overall, this increase was
218 robust for most values of P and Δ_{peak} (Fig. 6B3).

219 Increasing Δ_{peak} for a constant value of P and g_{max} (Fig. 6C1), produced a small
220 but significant increase in φ_{LP} (Three-Way ANOVA, $p < 0.001$, $F = 17.172$). This effect
221 was robust for most values of P and g_{max} , for both C-DC and C-Dur (Fig. 6C2 and 6C3).

222 These results showed that all three parameters that define the shape of the
223 IPSC influence φ_{LP} . Clearly, the strongest effect is the decrease in φ_{LP} as a function of
224 P . However, φ_{LP} modestly increases as a function of the other two parameters, g_{max}
225 and Δ_{peak} . This raised the question how g_{max} and Δ_{peak} would have to change in
226 coordination as a function of P to counteract the effect of P on φ_{LP} and achieve phase
227 constancy.

228 *Coordinated changes of g_{max} and Δ_{peak} produce the largest effect on phase*

229 To explore how g_{max} and Δ_{peak} might interact to influence φ_{LP} , we examined the
230 sensitivity of φ_{LP} to these two parameters, individually and in combination, for all values
231 of P in our data (see Methods). Sensitivity of φ_{LP} to these two parameters varied across
232 P values, with larger sensitivity at lower values of P (data not shown; Two-Way RM-
233 ANOVA, $p < 0.001$, $F = 16.054$). For simplicity, we averaged the sensitivity values across
234 different P values to obtain an overall measure of the influence of g_{max} and Δ_{peak} . These
235 results showed that, for the C-DC case, φ_{LP} had a positive sensitivity to g_{max} and a
236 smaller positive sensitivity to Δ_{peak} (Fig. 7A). The sensitivity was largest if the two
237 parameters were varied together ($g_{max} + \Delta_{peak}$) and smallest if they were varied in
238 opposite directions ($g_{max} - \Delta_{peak}$; Two-Way RM-ANOVA, $p < 0.001$, $F = 3.330$). Similarly,
239 these sensitivity values were also significantly different for the C-Dur case (Fig. 7B;
240 Two-Way RM-ANOVA, $p < 0.001$, $F = 2.892$), with largest sensitivity for $g_{max} + \Delta_{peak}$ and
241 smallest for $g_{max} - \Delta_{peak}$.

242 *Level sets of φ_{LP} in the P - g_{max} - φ_{peak} space for C-DC and C-Dur cases*

243 To search for phase constancy across different P values in our dataset, we
244 expressed φ_{LP} as a function of the three IPSC parameters, P , g_{max} and Δ_{peak} :
245 $\varphi_{LP} = \Phi(P, g_{max}, \Delta_{peak})$. **Figure 8** shows heat map plots of the function Φ , plotted for the
246 range of values of P and Δ_{peak} and four values of g_{max} . In these plots, phase constancy
247 can be seen as the set of values in each graph that are isochromatic, indicating the
248 level sets of the function Φ . These level sets are mathematically defined as
249 hypersurfaces on which the function has a constant value: $\Phi(P, g_{max}, \Delta_{peak}) = \varphi_c$. For the
250 C-DC case, in each g_{max} section of the plot, the level sets (e.g. $\varphi_c = 0.34$ denoted in
251 white) spanned a moderate range of P values as Δ_{peak} increased (**Fig. 8A1**). The span
252 of P values across all four panels indicates the range of cycle periods for which phase
253 constancy could be achieved by varying g_{max} and Δ_{peak} . This range of P values
254 (spanned by the white curves) was considerably smaller for the C-Dur case (**Fig. 8A2**).

255 For any constant phase value φ_c , these level sets can be expressed as

$$256 \quad P = P_{\varphi_c}(g_{max}, \Delta_{peak}),$$

257 which describes a surface in the 3D space, yielding the P value for which phase can be
258 maintained at φ_c , for the given values of g_{max} and Δ_{peak} . The level set indicated by the
259 white curves in panel A for the C-DC case is plotted as a heat map in **Fig. 8B1** and can
260 be compared with the same plot for the C-Dur case in **Fig. 8B2**. The range of colors in
261 each plot (marked next to each panel) indicates the range of P values for which phase
262 can be kept at $\varphi_c = 0.34$. To reveal how this range depends on the desired phase, we
263 measured this range for all values of φ_c between 0.2 and 0.8 (**Figs. 8C1 and 8C2**). We
264 found that the LP neuron could not achieve phases below 0.3 in the C-DC case (**Fig.**
265 **8C1**), which is simply because the neuron never fired during the inhibitory synaptic
266 current (which had a duty cycle of 0.3). Furthermore, the range of P values for which
267 the LP phase could be maintained by varying g_{max} and Δ_{peak} was much larger for C-DC
268 inputs compared to C-Dur Inputs, for all φ_c values between 0.31 and 0.54.

269 *A model of synaptic dynamics could predict activity onset phase of LP neuron*

270 To gain a better understanding of our experimental results, we considered
271 **Equations (7) and (8)**—the mathematical description of φ_{LP} as a function of P , g_{max} and
272 Δ_{peak} , for the C-Dur and C-DC cases, respectively, that we derived in the Methods
273 section—repeated here for convenience:

$$274 \quad \varphi_{LP} = F(P, g_{max}, \Delta_{peak}) = \frac{\tau_s}{P} \ln \frac{g_{max}}{g^*} + \frac{\Delta_{peak} T_{act}}{P}, \quad (7)$$

$$275 \quad \varphi_{LP} = F(P, g_{max}, \Delta_{peak}) = \frac{\tau_s}{P} \ln \frac{g_{max}}{g^*} + \Delta_{peak} DC. \quad (8)$$

276 In the C-Dur case, described by Equation (7), the input period has the most
 277 significant affect and φ_{LP} decays like $1/P$. In contrast, in the C-DC case, described by
 278 Equation (8), φ_{LP} is bounded from below by $\Delta_{peak} \cdot DC$ and thus behaves very differently
 279 than in the C-Dur case. In particular, as P increases, φ_{LP} approaches $\Delta_{peak} \cdot DC$ for the
 280 C-DC case, whereas it approaches 0 in the C-Dur case.

281 Keeping g_{max} (respectively, Δ_{peak}) constant in these equations allows us to
 282 obtain a relationship between P and Δ_{peak} (respectively, g_{max}), for which φ_{LP} is kept
 283 constant at φ_c . Consider Equations (7) and (8) for fixed values of both φ_{LP} ($= \varphi_c$) and
 284 g_{max} . Then these equations reduce to simple functional relationships where Δ_{peak} can be
 285 expressed as a function of P . In the C-DC case, for example, evaluating Δ_{peak} from
 286 Equation (8) produces

$$287 \quad g_{max} = g^* \cdot \exp\left(\frac{P}{\tau_s} (\varphi_c - \Delta_{peak} DC)\right) \quad (1)$$

288 Equation (12) describes how g_{max} must vary with P for the system to maintain a
 289 constant phase φ_c for any given Δ_{peak} .

290 Alternatively, Δ_{peak} can be expressed as a function of P . In the C-DC case,
 291 evaluating Δ_{peak} from Equation (8) produces

$$292 \quad \Delta_{peak} = \frac{\varphi_c}{DC} - \frac{\tau_s}{DC \cdot P} \ln \frac{g_{max}}{g^*}, \quad (2)$$

293 This equation describes how Δ_{peak} must vary with P for the system to maintain a
 294 constant phase φ_c for any given g_{max} . A comparison of these two cases can be seen in
 295 Fig. 9A, where either g_{max} (green) or Δ_{peak} (blue) is varied, while keeping the other
 296 parameter constant, to keep φ_{LP} constant at $\varphi_c=0.34$ across different P values. (The
 297 red curve is the depressing case, described below.) As the figure shows, phase
 298 constancy can be achieved by varying either parameter, but each parameter produces
 299 a different range of P across which phase is maintained.

300 In fact, Equation (13) can be used to calculate the range of P values over which
 301 changing Δ_{peak} (from 0 to 1) can maintain a constant phase φ_c . Solving $0 < \Delta_{peak} < 1$
 302 using Equation (13) yields

$$303 \quad \frac{\tau_s}{\varphi_c} \ln \frac{g_{max}}{g^*} < P_{DC} < \frac{\tau_s}{\varphi_c - DC} \ln \frac{g_{max}}{g^*} \quad (3)$$

304 Performing the same procedure in the C-Dur case, we find

$$305 \quad \frac{\tau_s}{\varphi_c} \ln \frac{g_{max}}{g^*} < P_{Dur} < \frac{T_{act}}{\varphi_c} + \frac{\tau_s}{\varphi_c} \ln \frac{g_{max}}{g^*}. \quad (4)$$

306 If ΔP denotes the range of P values that respectively satisfy Equation (14) or (15), it is
 307 straightforward to show that $\Delta P_{DC} > \Delta P_{Dur}$ (compare black and blue curves in Fig. 9B
 308 and 9C). To see this, note that the lower bound for each interval is the same, thus we
 309 need to show that the upper bound for the C-DC case is larger than that for the C-Dur
 310 case. That is,

$$311 \quad \varphi_c < DC \left(1 + \frac{\tau_s}{T_{act}} \ln \frac{g_{max}}{g^*} \right). \quad (5)$$

312 Equation (16) is true for τ_s and g_{max} large enough.

313 Two additional points are notable from Fig. 9B. First, the lower bound on φ_{LP} for
 314 which phase constancy can occur (i.e., φ_c) is smaller in the C-Dur than C-DC case.
 315 This is because we have assumed that in the C-DC case the LP neuron cannot fire
 316 during inhibition (i.e., until after $\Delta_{peak} DC$). Second, for φ_c larger than ~ 0.5 , ΔP is larger
 317 for the C-Dur case. This occurs because, when φ_c is sufficiently large, Equation (16)
 318 can no longer be satisfied. These findings are consistent with our experimental results
 319 described above, indicating that although phase constancy can be achieved when
 320 either g_{max} or Δ_{peak} increases with P , a concomitant increase of both—which could occur
 321 for example with a depressing synapse—greatly expands the range of P values for
 322 which a constant phase is maintained.

323 We now consider how short-term depression of the synapse—a property known
 324 to exist in the pacemaker to LP synapse (Zhao et al., 2011)—influences phase
 325 constancy by changing g_{max} and Δ_{peak} . We will restrict this section to the C-DC case. A
 326 similar derivation can be made for the C-Dur case. As mentioned in the Methods, the

327 effect of synaptic depression on synaptic strength can be obtained by Equation (11)
328 (repeated from the Methods):

$$329 \quad g_{max} = \bar{g}_{max} \cdot s_{max}(P) \quad (11)$$

330 where s_{max} is the maximum value of s_d at the onset of the pacemaker burst:

$$331 \quad s_{max}(P) = \frac{1 - e^{-P(1-DC)/\tau_\alpha}}{1 - e^{-P(1-DC)/\tau_\alpha} e^{-DC \cdot P/\tau_\beta}}.$$

332 Note that s_{max} is a monotonically increasing function with values between 0 and 1. Its
333 value approaches 1 as P increases, indicating that the synapse becomes stronger. In
334 this equation, \bar{g}_{max} is constant and is chosen so that the non-depressing and
335 depressing conductances match at $P = 1$ s. As seen in Fig. 9A, when synaptic
336 depression dictates how g_{max} varies with P as in Equation (11), and Δ_{peak} varies with P
337 and g_{max} according to Equation (13), the simultaneous changes in g_{max} and Δ_{peak} (red)
338 greatly increase the range of P values over which ϕ_{LP} is constant.

339 Returning to Fig. 9B, note that the C-DC case with depression spans a larger
340 range of P values than the non-depressing case. Similarly, in Fig. 9C, the range of P
341 values for which phase can be maintained is larger than the non-depressing case
342 across ϕ_{LP} values, except where ϕ_{LP} is so large that the depressing synapse operates
343 outside its dynamic range.

344 Discussion

345 *The importance of phase in oscillatory networks*

346 A common feature of oscillatory networks is that the activities of different
347 neuron types are restricted to specific phases of the oscillation cycle. For example,
348 different hippocampal and cortical neurons are active in at least three distinct phases of
349 the gamma rhythm (Hajos et al., 2004; Hasenstaub et al., 2005), and distinct
350 hippocampal neuron types fire at different phases of the theta rhythm and sharp wave-
351 associated ripple episodes (Somogyi and Klausberger, 2005).

352 Experimental studies quantify the latency of neural activity with respect to a
353 reference time in the cycle, but in most cases, these latencies are normalized and
354 reported as phase. Distinct neuron types can maintain a coherent activity phase,

355 despite wide variations in the network frequency (30-100 Hz for gamma rhythms, 4-7
356 Hz for theta rhythms, and 120-200 Hz for sharp wave-associated ripple episodes).
357 Phase-specific activity of different neuron types is proposed to be important in rhythm
358 generation (Wang, 2010), and indicates the necessity of precise timing for producing
359 proper circuit output and behavior (Kopell et al., 2011). For example, phase locking of
360 spike patterns to oscillations is important for auditory processing, single cell and
361 network computations and Hebbian learning rules (Kayser et al., 2009; McLelland and
362 Paulsen, 2009; Panzeri et al., 2010). For brain oscillations, phase relationships may
363 provide clues about the underlying circuit connectivity and dynamics, but a behavioral
364 correlate of varying frequencies is not obvious. In contrast, the activity phase of distinct
365 neuron types in rhythmic motor circuits is a tangible readout of the timing of motor
366 neurons and muscle contractions, thus defining phases of movement (Grillner and El
367 Manira, 2015; Kiehn, 2016; Le Gal et al., 2017; Bidaye et al., 2018). Because
368 meaningful behavior depends crucially on proper activity phases, whether neurons
369 maintain their activity phase in face of changes in frequency simply translates to
370 whether the movement pattern changes as it speeds up or slows down.

371 *Determinants of phase*

372 In oscillatory networks, the activity phases of different neuron types depend to
373 different degrees on the precise timing and strength of their synaptic inputs (Oren et
374 al., 2006). Our results from noise current injections showed that the timing of the LP
375 neuron is strongly dependent on the timing of inputs it receives. Dynamic clamp
376 injection of realistic or triangular conductance waveforms with different periods (P)
377 indicated that φ_{LP} was largely determined by the duration of the synaptic input. φ_{LP}
378 changed substantially with P when inputs had constant duration, but much less when
379 inputs had a constant duty cycle, i.e., when duration scaled with P . However, our
380 experiments also showed that inputs of constant duty cycles alone are insufficient for
381 phase constancy. φ_{LP} decreased with P even with a constant duty cycle of inputs, but
382 increased with either synaptic strength (g_{max}) or peak phase of the synaptic input
383 (Δ_{peak}). The increase in φ_{LP} had similar sensitivity to g_{max} and Δ_{peak} , and therefore a
384 larger sensitivity to a simultaneous increase in both. Consequently, it was possible to
385 keep φ_{LP} constant over a wide range of cycle periods by increasing both parameters
386 with P .

387 The fact that an increase in g_{max} with P promotes phase constancy is
388 biologically relevant, as short-term depression in pyloric synapses means that synaptic
389 strength indeed increases with P (Manor et al., 1997). Previous modeling studies show
390 that short-term synaptic depression of inhibitory synapses promotes phase constancy
391 (Nadim et al., 2003; Bose et al., 2004), largely because of longer recovery times from
392 depression at larger values of P .

393 The case is less clear for the finding that an increase of Δ_{peak} with P promotes
394 phase maintenance, as we have previously shown that Δ_{peak} in LP actually decreases
395 with P (Manor et al., 1997; Tseng et al., 2014). On the face of it, this suggests that an
396 increase in Δ_{peak} is not a strategy employed in the intact circuit. However, the caveat is
397 that such results may critically depend on the cause of the change in P , either
398 technically and biologically. While in our current study we varied Δ_{peak} with direct
399 conductance injection into LP, previous results were obtained by changing the
400 waveform and period of the presynaptic pacemaker neurons. When P is changed in an
401 individual preparation by injecting current into or voltage-clamping the pacemakers,
402 phase of follower neurons is not particularly well maintained. An example of this is
403 shown in Fig. 1, where ϕ_{LP} values fall between constant phase and constant duration
404 and, additionally, all pyloric neurons show behavior that falls between constant phase
405 and constant latencies (Hooper, 1997b, a). This may reflect that individuals are not
406 keeping phase particularly well when the only cause of changing P is the presynaptic
407 input. This is supported by the observation that even during normal ongoing pyloric
408 activity, phases change with cycle-to-cycle variability of P in individual preparations
409 (Bucher et al., 2005). However, it does not preclude the possibility that Δ_{peak} plays an
410 important role in stable phase relationships when P differs because of temperature,
411 neuromodulatory conditions, or inter-individual variability (discussed below).

412 It is noteworthy that a change in the synaptic strength or peak phase with P is
413 not peculiar to graded synapses. The fact that short-term synaptic plasticity can act as
414 a frequency-dependent gain control mechanism is well known for many spike-mediated
415 synaptic connections. In bursting neurons, the presence of a combination of short-term
416 depression and facilitation in the same spike-mediated synaptic interaction could also
417 result in changes in the peak phase of the summated synaptic current as a function of
418 burst frequency and duration, and the intra-burst spike rate (Markram et al., 1998).

419 *Phase relationships in changing temperatures*

420 An interesting case is provided by the observation that phases are remarkably
421 constant when pyloric rhythm frequency is changed with temperature. Tang et al.
422 (2012) report a 4-fold decrease in P of the pyloric rhythm between 7 and 23° C. In this
423 study, none of the pyloric phases changed significantly, and it is worth noting that
424 under conditions of changing temperatures, the relationships between P , g_{max} , and
425 Δ_{peak} appeared to be fundamentally different from when P is changed at a constant
426 temperature. Presynaptic voltage trajectories scaled with changing P , and Δ_{peak} of
427 postsynaptic currents was independent of P , in contrast to the decrease described at
428 constant temperature (Manor et al., 1997; Tseng et al., 2014). Amplitudes of synaptic
429 potentials did not change with temperature, despite an increase in synaptic current
430 amplitudes with increasing temperature (and associated decrease in P). This is in
431 contrast to the positive relationship between g_{max} and P that results from synaptic
432 depression at a constant temperature (Manor et al., 1997). Therefore, it appears that
433 the likely substantial effects of temperature on synaptic dynamics and ion channel
434 gating are subject to a set of compensatory adaptations different from when P is
435 changed at constant temperature.

436 *Slow compensatory regulation of phase*

437 Phase maintenance in the face of changing P in an individual animal requires
438 the appropriate short-term dynamics of synaptic and intrinsic neuronal properties. The
439 fact that characteristic (and therefore similar) phase relationships can also be observed
440 under the same experimental conditions across individual preparations is a different
441 conundrum, particularly when P can vary substantially, as is true for brain oscillations
442 (Hajos et al., 2004; Hasenstaub et al., 2005; Somogyi and Klausberger, 2005). Phases
443 show different degrees of variability across individuals in a variety of systems, e.g.,
444 leech heartbeat (Wenning et al., 2018), larval crawling in *Drosophila* (Pulver et al.,
445 2015), and fictive swimming in zebrafish (Masino and Fetcho, 2005), but in all of these
446 cases phases are not correlated with P . In the pyloric rhythm, phases are also variable
447 to a degree across individuals, but not correlated with the mean P , which varies >2-fold
448 (Bucher et al., 2005; Goillard et al., 2009). This phase constancy occurs despite
449 considerable inter-individual variability in ionic currents, and is considered the ultimate
450 target of slow compensatory regulation, i.e., homeostatic plasticity (Marder and
451 Goillard, 2006; Ma and LaMotte, 2007; Marder et al., 2014). Slow compensation can
452 also be observed directly when rhythmic activity is disrupted by decentralization, and

453 subsequently recovers to similar phase relationships over the course of many hours
454 (Luther et al., 2003). It is interesting to speculate if our findings about how synaptic
455 parameters must change to keep phase constant would hold across individuals with
456 different mean P . The prediction would be coordinated positive correlations of both g_{max}
457 and Δ_{peak} with P .

458 *Phase relationships under different neuromodulatory conditions.*

459 The flipside of the question how neurons maintain phase is the question how
460 their phase can be changed. In motor systems in particular, changes in phase
461 relationships are functionally important to produce qualitatively different versions of
462 circuit output, for example to produce different gaits in locomotion (Vidal-Gadea et al.,
463 2011; Grillner and El Manira, 2015; Kiehn, 2016). The activity of neural circuits is
464 flexible, and much of this flexibility is provided by modulatory transmitters and
465 hormones which alter synaptic and intrinsic neuronal properties (Brezina, 2010; Harris-
466 Warrick, 2011; Jordan and Slawinska, 2011; Bargmann, 2012; Marder, 2012; Bucher
467 and Marder, 2013; Nadim and Bucher, 2014). The pyloric circuit is sensitive to a
468 plethora of small molecule transmitters and neuropeptides which affect cycle frequency
469 and phase relationships (Marder and Bucher, 2007; Stein, 2009; Daur et al., 2016).
470 With respect to our findings, any given neuromodulator could act presynaptically to
471 alter P , duration, or duty cycle on the one hand, and g_{max} and Δ_{peak} on the other. In
472 addition, the neuromodulator could affect the postsynaptic neuron's properties and
473 alter its sensitivity to any of these parameters. Therefore, our findings could not just
474 further our understanding of how phase can be maintained across different rhythm
475 frequencies, but also provide a framework for testing if and how changes in synaptic
476 dynamics may contribute to altering phase relationships under different
477 neuromodulatory conditions.

478 **Materials and Methods**

479 Adult male crabs (*Cancer borealis*) were acquired from local distributors and
480 maintained in aquaria filled with chilled (10-13°C) artificial sea water until use. Crabs
481 were prepared for dissection by placing them on ice for 30 minutes. The dissection was
482 performed using standard protocols as previously described (Tohidi and Nadim, 2009;
483 Tseng and Nadim, 2010). The STNS, including the four ganglia (esophageal ganglion,
484 two commissural ganglia, and the STG) and their connecting nerves, and the motor

485 nerves arising from the STG, were dissected from the stomach and pinned into a
486 Sylgard (Dow-Corning) lined Petri dish filled with chilled saline. The STG was
487 desheathed, exposing the somata of the neurons for intracellular impalement.
488 Preparations were superfused with chilled (10-13°C) physiological *Cancer* saline
489 containing: 11 mM KCl, 440 mM NaCl, 13 mM CaCl₂ · 2H₂O, 26 mM MgCl₂ · 6H₂O,
490 11.2 mM Trizma base, 5.1 mM maleic acid with a pH of 7.4.

491 Extracellular recordings were obtained from identified motor nerves using
492 stainless steel electrodes, amplified using a differential AC amplifier (A-M Systems,
493 model 1700). One lead was placed inside a petroleum jelly well created to electrically
494 isolate a small section of the nerve, the other right outside of it. For intracellular
495 recordings, glass microelectrodes were prepared using the Flaming-Brown
496 micropipette puller (P97; Sutter Instruments) and filled with 0.6 M K₂SO₄ and 20 mM
497 KCl. Microelectrodes used for membrane potential recordings had resistances of 25-
498 30MΩ; those used for current injections had resistances of 15-22 MΩ. Intracellular
499 recordings were performed using Axoclamp 2B and 900A amplifiers (Molecular
500 Devices). Data were acquired using pClamp 10 software (Molecular Devices) and the
501 Netsuite software (Gotham Scientific), sampled at 4-5 kHz and saved on a PC using a
502 Digidata 1332A (Molecular Devices) or a PCI-6070-E data acquisition board (National
503 Instruments).

504 Individual pyloric neurons were impaled and identified via their membrane
505 potential waveforms, correspondence of spike patterns with extracellular nerve
506 recordings, and interactions with other neurons within the network (Weimann et al.,
507 1991).

508 *Constructing realistic IPSC waveforms*

509 Inhibitory postsynaptic currents (IPSCs) were recorded from LP neurons during
510 the ongoing rhythm using two-electrode voltage clamp and holding the LP neuron
511 at -50mV, far from the IPSC reversal potential of ~ -80 mV (Fig. 3A). When the LP
512 soma is voltage clamped at -50 mV, the axon (which is electrotonically distant from the
513 soma) can still produce action potentials following the synaptic inhibition from the
514 pacemaker neurons. The onset of the LP neuron action potentials (recorded in the
515 current trace) was used to calculate the mean IPSC for each experiment averaging the
516 IPSCs over 10-20 cycles. The IPSC waveforms were then extracted by normalizing
517 both the amplitude and the duration of the mean IPSC.

518 *Driving the LP neuron with noise current*

519 In these experiments, the preparation was superfused in *Cancer* saline plus
520 10^{-5} M picrotoxin (PTX; Sigma Aldrich) for 30 minutes to block the synaptic currents to
521 the LP neuron. The removal of synaptic inhibition onto LP neurons changed the
522 activity of these neurons from bursting to tonic firing. Then, noise current, generated by
523 the Ornstein-Uhlenbeck process (Lindner), was injected into the isolated LP neurons
524 for 60 minutes using the Scope software (available at
525 <http://stg.rutgers.edu/Resources.html>, developed in the Nadim laboratory). The baseline
526 of the noise current was adjusted by adding DC current so that it can provide enough
527 inhibition to produce silent periods alternating with bursts of action potentials.

528 *Driving the LP neuron with realistic or triangular IPSC waveforms in dynamic clamp*

529 The dynamic clamp current was injected using the Netclamp software (Netsuite,
530 Gotham Scientific). The synaptically isolated (10^{-5} M PTX) LP neuron was driven with
531 an artificial synaptic current in dynamic clamp. The synaptic current was given as

532
$$I_{syn} = g_{syn} (V_{LP} - E_{syn})$$

533 where the synaptic conductance g_{syn} was a pre-determined waveform, repeated
534 periodically with period P , and E_{syn} was the synaptic reversal potential set to -80 mV
535 (Zhao et al., 2011).

536 Two sets of dynamic clamp experiments were performed on different animals.
537 In one set of experiments, g_{syn} was set to be a triangular waveform. We measured the
538 effects of four different parameters in these triangle conductance injections (Fig. 1):
539 peak phase (Δ_{peak}), duration (T_{act}), period (P = time between onsets of dynamic clamp
540 synaptic injections), and maximal conductance (g_{max} , the peak value of g_{syn}). This
541 allowed us to explore which combinations of the different parameters influences the LP
542 phase. Five values for P were used: 500, 750, 1000, 1500, and 2000 ms, which cover
543 the typical range of pyloric cycle periods. Three values of g_{max} were used: 0.1, 0.2 and
544 0.4 μ S, consistent with previous measurements of synaptic conductance (Zhao et al.,
545 2011; Tseng et al., 2014). The value of Δ_{peak} was varied to be 0, 0.25, 0.5, 0.75 or 1. In
546 the same experiment, all runs were done in two conditions: with T_{act} constant across
547 different P values (C-Dur case with $T_{act} = 300$ ms) or with T_{act} changing proportionally
548 to P (C-DC case with duty cycle $DC = T_{act} / P = 0.3$).

549 In the other set of experiments, g_{syn} was a realistic IPSC waveform, based on a
550 pre-recorded IPSC in the LP neuron. In these experiments, P was varied to be 500,
551 750, 1000, 1250, 1500, or 2000 ms by scaling the realistic waveform in the time
552 direction. In these experiments, g_{max} was set to be 0.1, 0.2, 0.4, 0.6, or 0.8 μ S. The LP
553 neuron burst onset delay (Δt) was measured relative to the onset of the pacemaker
554 component of the synaptic input (identified by the kink in the synaptic conductance
555 waveform) in each cycle. The burst phase was calculated as $\phi_{LP} = \Delta t / P$. Phase
556 constancy means that Δt changed proportionally to P . To measure the LP neuron
557 phase with respect to the end of the pacemaker input, this reference used was the
558 point on the synaptic conductance waveform marked by drawing a horizontal line from
559 the kink that identified the onset of the pacemaker input.

560 *Determining relationship between cycle period (P), synaptic strength (g_{max}) and LP* 561 *phase (ϕ_{LP}) using the realistic IPSC waveform*

562 We determined how well the mathematical model derived for constant input
563 duty cycles (see Equation 8 below), matched the experimental data obtained with
564 realistic IPSC waveforms. To this end, we fit the model to ϕ_{LP} values measured for all
565 values of g_{max} and P , using the standard fitting routine 'fit' in MATLAB (Mathworks).

566 *Sensitivity of ϕ_{LP} to g_{max} and Δ_{peak} across all P values*

567 To explore how g_{max} and Δ_{peak} might interact to influence ϕ_{LP} , we first examined
568 the sensitivity of ϕ_{LP} to these two parameters, individually and in combination, for all
569 values of P in our data. For each cycle period, we computed the mean phase across all
570 of our experiments (N=9) and all values of g_{max} (0.1, 0.2 and 0.4 μ S) and we
571 interpolated the ϕ_{LP} for g_{max} (0.2 and 0.4) to obtain ϕ_{LP} for 0.3 and Δ_{peak} (0, 0.25, 0.5,
572 0.75 or 1). This produced a 4 by 5 matrix of all of the values. For each data point in the
573 matrix, we moved along eight different directions (increase both g_{max} and Δ_{peak} , increase
574 g_{max} , increase Δ_{peak} , decrease both g_{max} and Δ_{peak} , decrease g_{max} , decrease Δ_{peak} ,
575 increase g_{max} and decrease Δ_{peak} , and decrease g_{max} and increase Δ_{peak}) and calculated
576 the change in phase per unit (normalized) in g_{max} , Δ_{peak} , or both.

577 *A model of synaptic dynamics*

578 In the derivation of the model, the firing time of the LP neuron was assumed to
579 be completely determined by its synaptic input. This synaptic conductance (g_{syn}) was

580 assumed to rise and fall with distinct time constants. The following holds over one cycle
581 period and therefore time is reset with period P ($t \pmod{P}$):

$$582 \quad \frac{dg_{syn}}{dt} = \begin{cases} (g_{max} - g_{syn})\tau_r & t \pmod{P} < t_{peak} \\ -g_{syn} / \tau_s & t \pmod{P} \geq t_{peak} \end{cases} \quad (6)$$

583 where the time t_{peak} , corresponding to Δ_{peak} , is $t_{peak} = \Delta_{peak} T_{act}$. We assumed that LP
584 neuron remained inactive when g_{syn} was above a fixed threshold (g^*) less than g_{max} .
585 Because the synaptic input is periodic with period P , we solved for the minimum and
586 maximum values of g_{syn} in each cycle. The minimum (g_{lo}) occurred just before the onset
587 ($t = 0$) of AB/PD activity, whereas the maximum occurred at the peak synaptic phase
588 Δ_{peak} for the C-Dur case. In the C-DC case, $T_{act} = DC \cdot P$, where DC is the duty cycle
589 (fixed at 0.3 in our experiments).

590 To calculate g^* , we set the value $t = 0$ so that $g_{syn}(0) = g_{lo}$ (and, by periodicity,
591 $g_{syn}(P) = g_{lo}$), and solved the first part of Equation (1) where g_{syn} increases until $t = t_{peak}$.
592 This yielded

$$593 \quad g_{peak} = g_{syn}(t_{peak}) = g_{max} + (g_{lo} - g_{max})e^{-t_{peak}/\tau_r} \quad (7)$$

594 We then used the second part of Equation (1) to track the decay of g_{syn} for $t_{peak} < t < P$:

$$595 \quad g_{syn}(t) = g_{peak}e^{-(t-t_{peak})/\tau_s} \quad (8)$$

596 Using Equation (3), we calculated the time Δt at which the synaptic conductance
597 $g_{syn}(\Delta t) = g^*$ as follows:

$$598 \quad g^* = g_{peak}e^{-(\Delta t - t_{peak})/\tau_s} \quad (9)$$

599 Solving Equation (4) for Δt yielded

$$600 \quad \Delta t = \tau_s \ln \frac{g(t_{peak})}{g^*} + t_{peak}.$$

601 Dividing this equation by P yielded ϕ_{LP} :

$$602 \quad \phi_{LP} = F(P, g_{max}, \Delta_{peak}) = \frac{\tau_s}{P} \ln \frac{g_{peak}}{g^*} + \frac{t_{peak}}{P}, \quad (10)$$

603 where g_{peak} is given by Equation (2). This expression provides a description of the
 604 dependence of φ_{LP} as a function of P , g_{max} and Δ_{peak} . To explore the role of the
 605 parameters in this relationship, we made a simplifying assumption that the synaptic
 606 conductance $g_{syn}(t)$ rapidly reached its peak (i.e., τ_r was small), stayed at this value and
 607 started to decay at $t = t_{peak}$. In this case $g(t) = g_{max}$ on the interval $(0, t_{peak})$ and the value
 608 of g_{lo} is irrelevant. With this assumption, Equation (5) reduced to

$$609 \quad \varphi_{LP} = \frac{\tau_s}{P} \ln \frac{g_{max}}{g^*} + \frac{t_{peak}}{P}. \quad (11)$$

610 Substituting $t_{peak} = \Delta_{peak} \cdot T_{act}$ in Equation (6), gave

$$611 \quad \varphi_{LP} = F(P, g_{max}, \Delta_{peak}) = \frac{\tau_s}{P} \ln \frac{g_{max}}{g^*} + \frac{\Delta_{peak} T_{act}}{P}, \quad (12)$$

612 which we used to describe the LP phase in the C-Dur case. To describe the C-DC
 613 case, after substituting $t_{peak} = \Delta_{peak} \cdot DC \cdot P$, we obtained

$$614 \quad \varphi_{LP} = F(P, g_{max}, \Delta_{peak}) = \frac{\tau_s}{P} \ln \frac{g_{max}}{g^*} + \Delta_{peak} DC. \quad (13)$$

615 Note that these equations also describe the relationship between φ_{LP} with T_{act} (or DC , in
 616 the C-DC case).

617 Equations (7) and (8) can be used to approximate a range of parameters over
 618 which φ_{LP} is maintained at a constant value. To do so, we assumed a specific
 619 parameter set, say $(\hat{P}, \hat{g}_{max}, \hat{\Delta}_{peak})$, satisfies

$$620 \quad F(\hat{P}, \hat{g}_{max}, \hat{\Delta}_{peak}) = \hat{\varphi}_{LP},$$

621 for some fixed phase value, $\hat{\varphi}_{LP}$. We could now ask whether there are nearby
 622 parameters for which phase remains constant, i.e., F remains equal to $\hat{\varphi}_{LP}$. The Implicit
 623 Function Theorem (Krantz and Parks, 2012) guarantees that this is the case, provided
 624 certain derivatives evaluated at $(\hat{P}, \hat{g}_{max}, \hat{\Delta}_{peak})$ are non-zero, which turns out to be true
 625 over a large range of parameters. Since the partial derivative with respect to Δ_{peak} of
 626 $F(P, g_{max}, \Delta_{peak})$ at this point is a non-zero constant equal to T_{act}/P (or DC) in the C-Dur
 627 (or C-DC) case, there is a function $\Delta_{peak} = h(P, g_{max})$ such that

628
$$F(P, g_{max}, h(P, g_{max})) = \hat{\phi}_{LP} \quad (14)$$

629 for values of P and g_{max} near (\hat{P}, \hat{g}_{max}) . In other words, the Implicit Function Theorem
 630 guarantees that small changes in P and g_{max} can be compensated for by an appropriate
 631 choice of Δ_{peak} in order to maintain a constant LP phase. A similar analysis can be
 632 done by solving for g_{max} in terms of P and Δ_{peak} or by solving for P in terms of g_{max} and
 633 Δ_{peak} .

634 *Adding synaptic depression to the model of synaptic dynamics*

635 In a previous modeling study, we explored how the phase of a follower neuron
 636 was affected when the inhibitory synapse from an oscillatory neuron to this follower
 637 had short-term synaptic depression (Manor et al., 2003). In that study the role of the
 638 parameter Δ_{peak} was not considered. It is, however, straightforward to add synaptic
 639 depression to **Equations (7) and (8)**.

640 An *ad hoc* model of synaptic depression can be made using a single variable s_d
 641 which will be a periodic function that denotes the extent of depression and takes on
 642 values between 0 and 1 (Bose et al., 2004). s_d decays during the AB/PD burst (from
 643 time 0 to T_{act} , indicating depression) and then recovers during the inter-burst interval
 644 (from T_{act} to P , indicating recovery). Thus, s_d can be described by an equation of the
 645 form:

646
$$\frac{ds_d}{dt} = \begin{cases} -s_d / \tau_\beta & t \pmod{P} \leq T_{act} \\ (1 - s_d) / \tau_\alpha & T_{act} < t \pmod{P} < P \end{cases}$$

647 Using periodicity, it is straightforward to show that the maximum value of s_d , which
 648 occurs at the start of the AB/PD burst, is given by:

649
$$s_{max}(P) = \frac{1 - e^{-P(1-DC)/\tau_\alpha}}{1 - e^{-P(1-DC)/\tau_\alpha} e^{-DC \cdot P/\tau_\beta}}. \quad (15)$$

650 Note that s_{max} is a monotonically increasing function with values between 0 and 1. Its
 651 value approaches 1 as P increases, indicating that the synapse becomes stronger. For
 652 a complete derivation and description, see (Bose et al., 2004). The effect of synaptic
 653 depression on synaptic strength can be obtained by setting

654
$$g_{max} = \bar{g}_{max} \cdot s_{max}(P) \quad (16)$$

655 where s_{max} is given by [Equation \(10\)](#).

656 *Software, analysis and statistics*

657 Data were analyzed using MATLAB scripts to calculate the time of burst onset
658 and the phase. Statistical analysis was performed using Sigmaplot 12.0 (Systat).
659 Significance was evaluated with an α value of 0.05, error bars and error values
660 reported denote standard error of the mean (SEM).

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665 **Competing interests**

666 The authors declare no competing financial interests.

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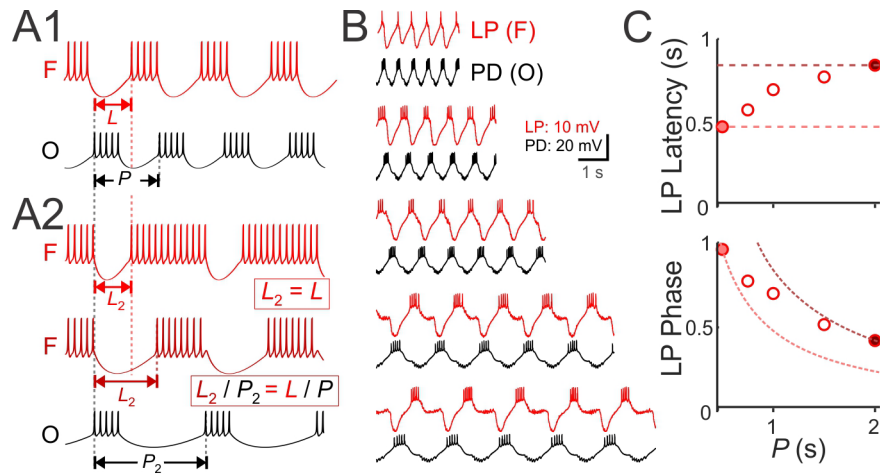
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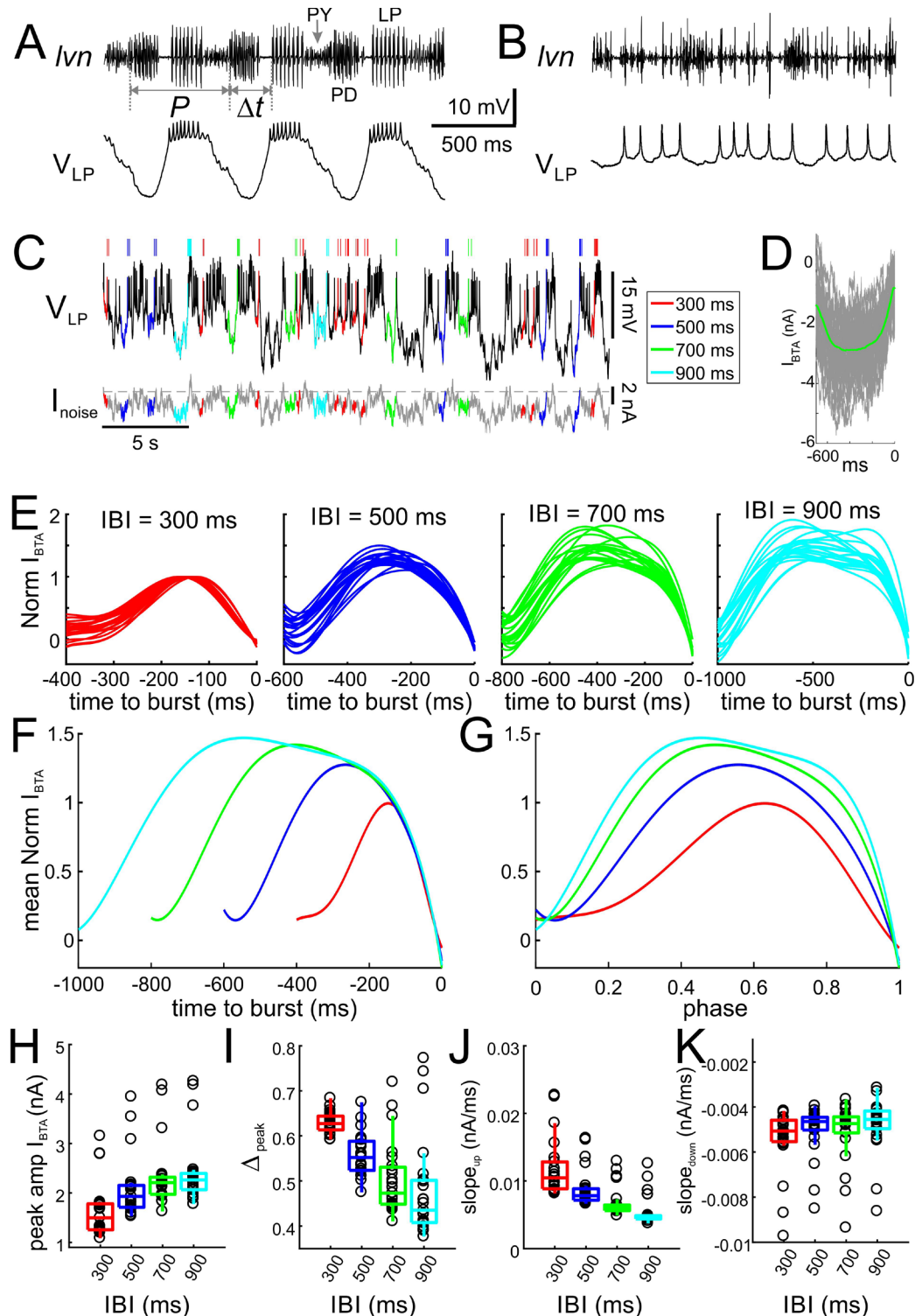
862 **Figures**



863

864 **Figure 1: Latency constancy and phase constancy as a function of period**

865 **A1.** Schematic diagram showing that a follower neuron (F) strongly inhibited by a
 866 bursting oscillatory neuron (O) with period P can produce rebound bursts with the
 867 same period at a latency L . **A2.** If the period of O changes to a new value (P_2), the new
 868 F burst latency (L_2) typically falls between two extremes: it could stay constant (top
 869 trace) or change proportionally to P_2 , so that the burst phase (L / P) remains constant
 870 (middle trace). **B.** Example traces of the pyloric pacemaker PD neuron and the follower
 871 LP neuron represent the O and F relationship in panel A. Here, the PD neuron is
 872 voltage clamped and a pre-recorded waveform of the same neuron is used to force the
 873 neuron to follow different cycle periods. **C.** A measurement of the LP neuron burst
 874 onset phase with respect to the onset of the PD neuron burst shows that it falls
 875 between the two limits of constant latency and constant phase. Dotted curves
 876 represent constant latency matched to the latencies at the two extreme P values.



877

878

Figure 2: Inputs to the LP neuron influence burst time, spike number and interval

879

A. Simultaneous intracellular recording of the LP neuron and extracellular recording of

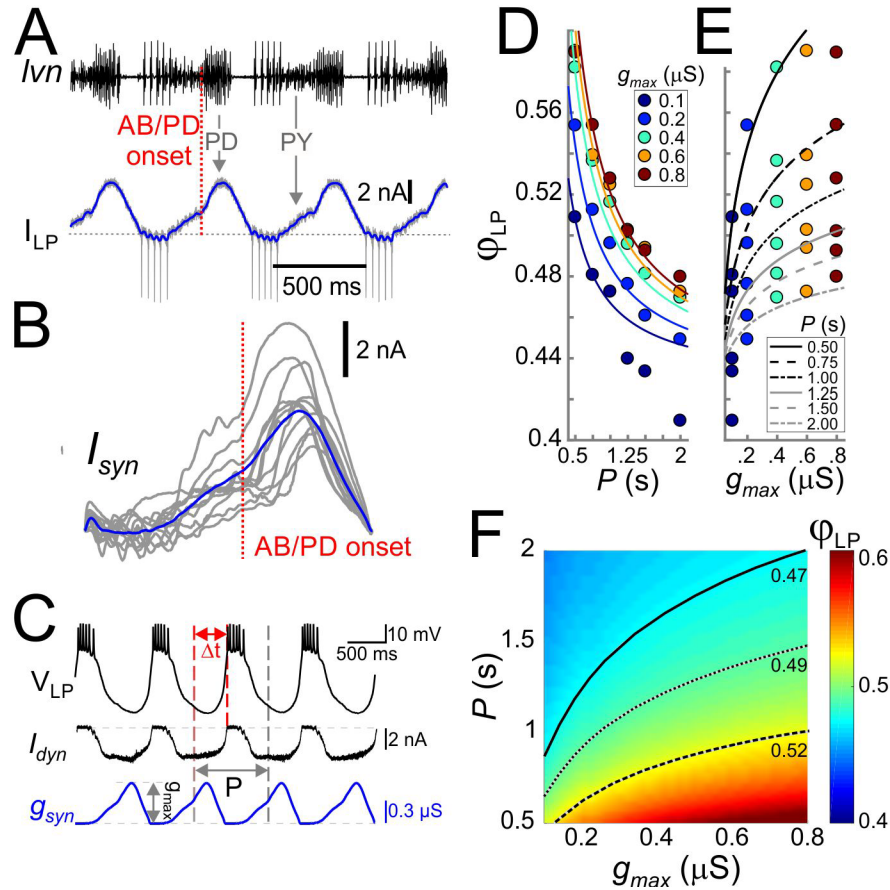
880

the lateral ventricular nerve (*I_{vn}*), containing the axons of the LP, PD and PY neurons

881

(arrows). Period (P) and the burst onset time (Δt) of the LP neuron are defined in

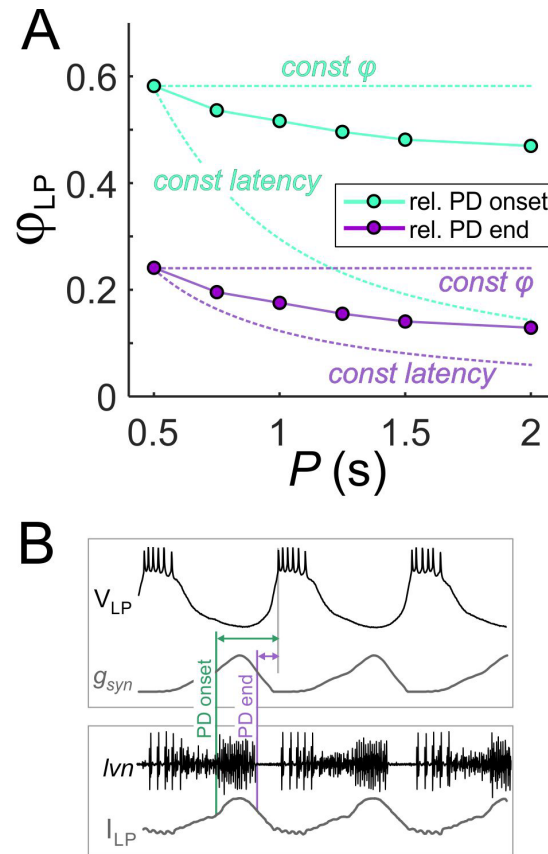
882 reference to the pacemaker group (PD) burst. **B.** Blocking the synaptic inputs (10 μ M
883 picrotoxin) to the LP neuron disrupts its bursting oscillations. **C.** The synaptically
884 isolated LP neuron was driven with a noise current input (I_{noise}) for 60 minutes. In
885 response, the LP neuron produced an irregular pattern of bursting. Specific inter-burst
886 intervals ($IBIs$) were tagged and used for burst-triggered averaging. **D.** Example of
887 burst-trigger-averaged input current (I_{BTA} , green). Individual traces are shown in grey.
888 **E.** For each IBI (300, 500, 700, 900 ms), I_{BTA} was calculated and normalized to the
889 amplitude of I_{BTA} for $IBI=300$ ms. Different traces in each panel show the I_{BTA} of
890 different preparations. **F.** The mean (across preparations) of the normalized I_{BTA} s
891 shown in panel E. **G.** Traces in panel F normalized by IBI . **H-K.** Four parameters define
892 the shape of the I_{BTA} : peak amplitude I_{amp} (**H**), peak phase Δ_{peak} (**I**), slope_{up} (**J**) and
893 slope_{down} (**K**) across preparations. IBI had a significant effect on amplitude I_{amp}
894 ($p<0.001$), peak phase Δ_{peak} ($p<0.001$), slope_{up} ($p<0.001$) and slope_{down} ($p=0.002$).



895

896 *Figure 3: Cycle period and synaptic strength affect the phase of LP burst onset in*
 897 *opposite directions.*

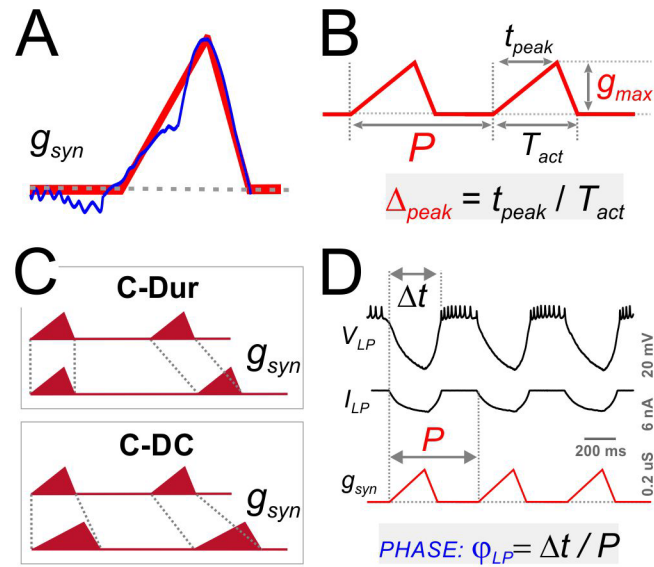
898 **A.** The synaptic input to the LP neuron was measured by voltage clamping it at a
 899 holding potential of -50mV during ongoing oscillations. The onset of the pacemaker
 900 (AB/PD) activity is seen as a kink in the synaptic current (I_{LP} , blue). **B.** Synaptic input
 901 averaged across cycles from 9 different LP neurons. The blue trace shows the
 902 average. **C.** An example of the LP neuron driven by the realistic synaptic waveform in
 903 dynamic clamp. The burst onset time (Dt) was measured relative to the AB/PD onset
 904 and used to measure the LP phase (ϕ_{LP}). g_{max} denotes the conductance amplitude. **D.**
 905 Mean ϕ_{LP} ($N=9$ preparations) shown as a function of P and fit with the function given by
 906 Equation (8) (fit values $\tau_s = 26.0$ ms, $g^* = 0.021$ μS and $\Delta_{peak} \cdot DC = 0.43$). **E.** Mean ϕ_{LP}
 907 plotted against g_{max} also shown with the fit to Equation (8). **F.** Heat map, obtained from
 908 fitting Equation (8) to the data in panels D and E, shows ϕ_{LP} as a function of both g_{max}
 909 and P . Black curves show the level sets of phase constancy for three values of ϕ_{LP}
 910 (0.47, 0.49, and 0.52).



911

912 *Figure 4: The constant duty cycle of synaptic conductance is a major factor in phase*
 913 *maintenance.*

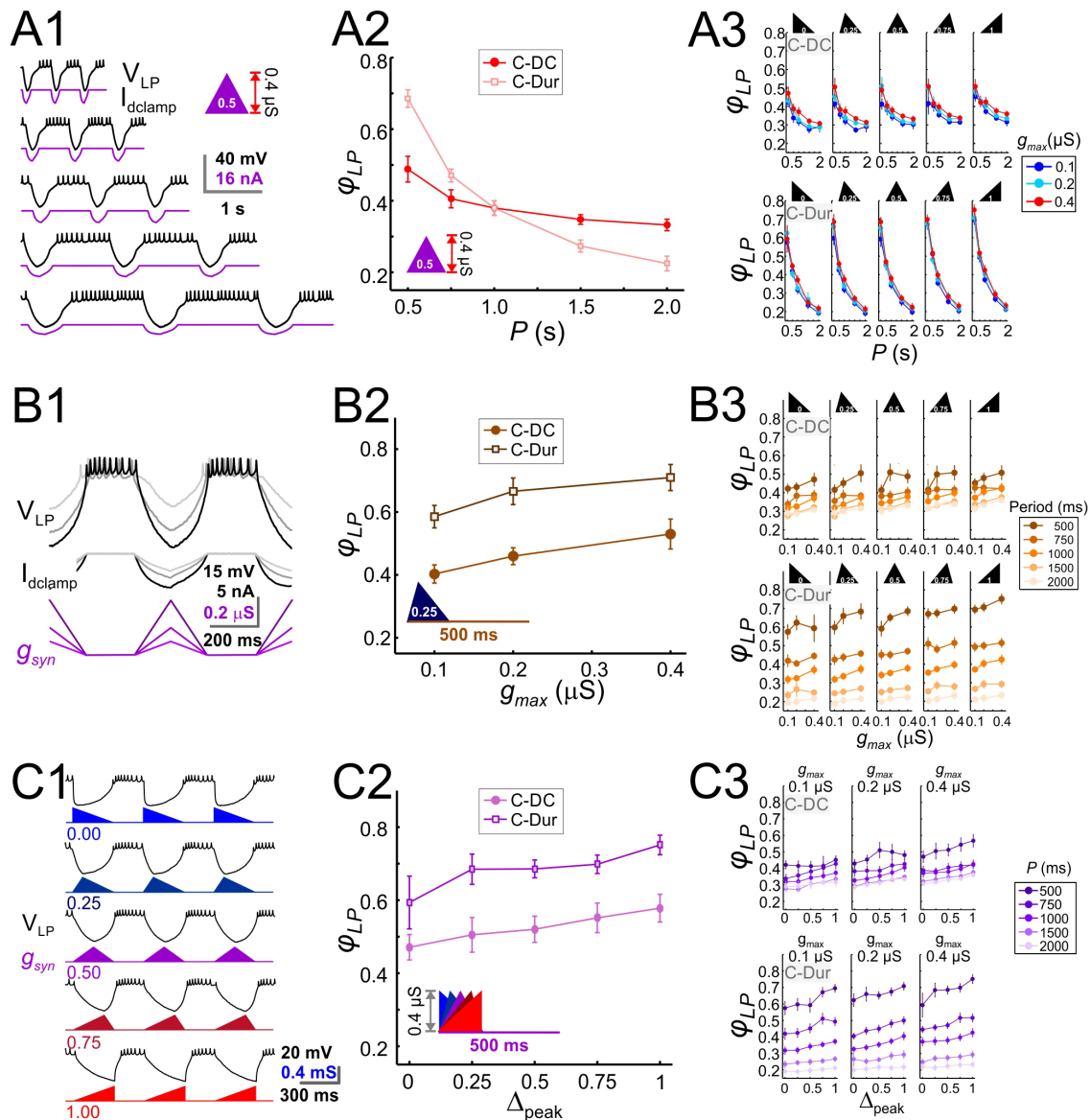
914 **A.** The change in ϕ_{LP} values with P are compared with the constant phase (solid curve)
 915 and constant duration (dashed curve) extremes. Lime traces show the usual values of
 916 ϕ_{LP} , calculated from the LP burst onset latency with respect to the onset of the PD
 917 burst. Lavender traces show ϕ_{LP} calculated from the LP burst onset latency with
 918 respect to the end of the PD burst. Data shown are the same as in Fig. 3D for $g_{max}=0.4$
 919 μS . **B.** Schematic diagram shows the latency of LP burst onset measured with respect
 920 to the (estimated) onset and end of the PD burst in the dynamic clamp experiments
 921 (see Methods). Bottom panel shows the synaptic current waveform measured in the
 922 voltage-clamped LP neuron during ongoing pyloric activity. Top panel shows the
 923 dynamic clamp injection of the synaptic conductance waveform into a synaptically-
 924 isolated LP neuron. The current waveform of the bottom panel is aligned to the
 925 conductance waveform of the top panel for the comparison used in determining the PD
 926 burst onset and end in the top panel.



927

928 *Figure 5: Four parameters describing synaptic shape were varied in the experimental*
 929 *paradigm.*

930 **A.** A triangle shaped conductance was used to mimic the synaptic input to the LP
 931 neuron. **B.** The triangular waveform can be described by period (P), duration (T_{act}),
 932 peak time (t_{peak}) and amplitude (g_{max}). **C.** In dynamic clamp runs, the synapse duration
 933 T_{act} was kept constant at 300 ms (C-Dur) or maintained at a constant duty cycle (T_{act}
 934 $/P$) of 0.3 (C-DC) across all values of P . **D.** Intracellular voltage recording of the
 935 synaptically isolated LP neuron during a dynamic clamp stimulation run using the
 936 triangle conductance. The burst onset time (Δt , calculated in reference to the synaptic
 937 conductance onset) was used to calculate the activity phase ($\phi_{LP} = \Delta t/P$).



938

939

Figure 6: The LP burst onset phase decreases as a function of P , but increases as a function of g_{max} and Δ_{peak} .

940

941 Periodic injection of an inhibitory triangular waveform conductance into the isolated LP

942 neuron produced bursting activity from which ϕ_{LP} was calculated. The parameters g_{max} ,

943 Δ_{peak} and P were varied across runs for both C-Dur and C-DC cases. **A.** ϕ_{LP} decreases

944 as a function of P . **A1.** Intracellular recording of an LP neuron showing a C-DC

945 conductance input across five periods. **A2.** ϕ_{LP} for the example shown in A1 plotted as

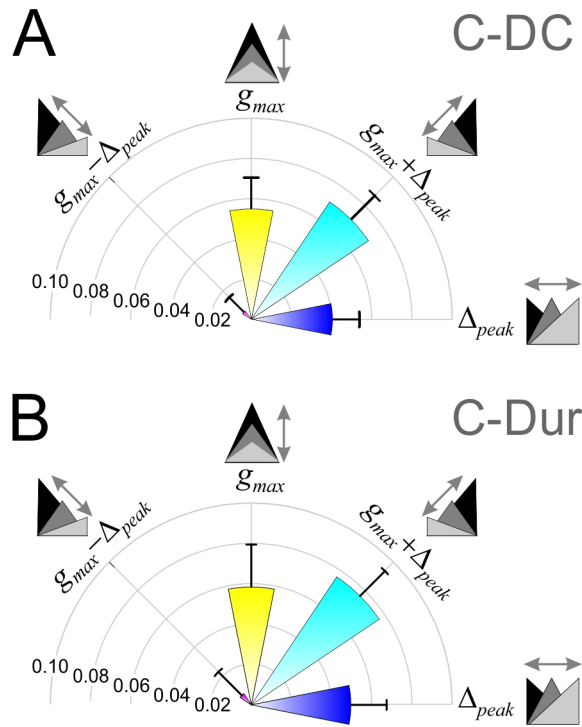
946 a function of P (for $g_{max} = 0.4 \mu\text{S}$, $\Delta_{peak} = 0.5$) for both C-Dur and C-DC cases. ϕ_{LP}

947 decreases rapidly with P and the drop is larger for the C-Dur case. **A3.** ϕ_{LP} decreased

948 with P in both the C-DC case (Three-Way RM ANOVA, $p < 0.001$, $F = 100.7$) and the C-

949 Dur case (Three-Way RM ANOVA, $p < 0.001$, $F = 466.4$) and all values of Δ_{peak} . The

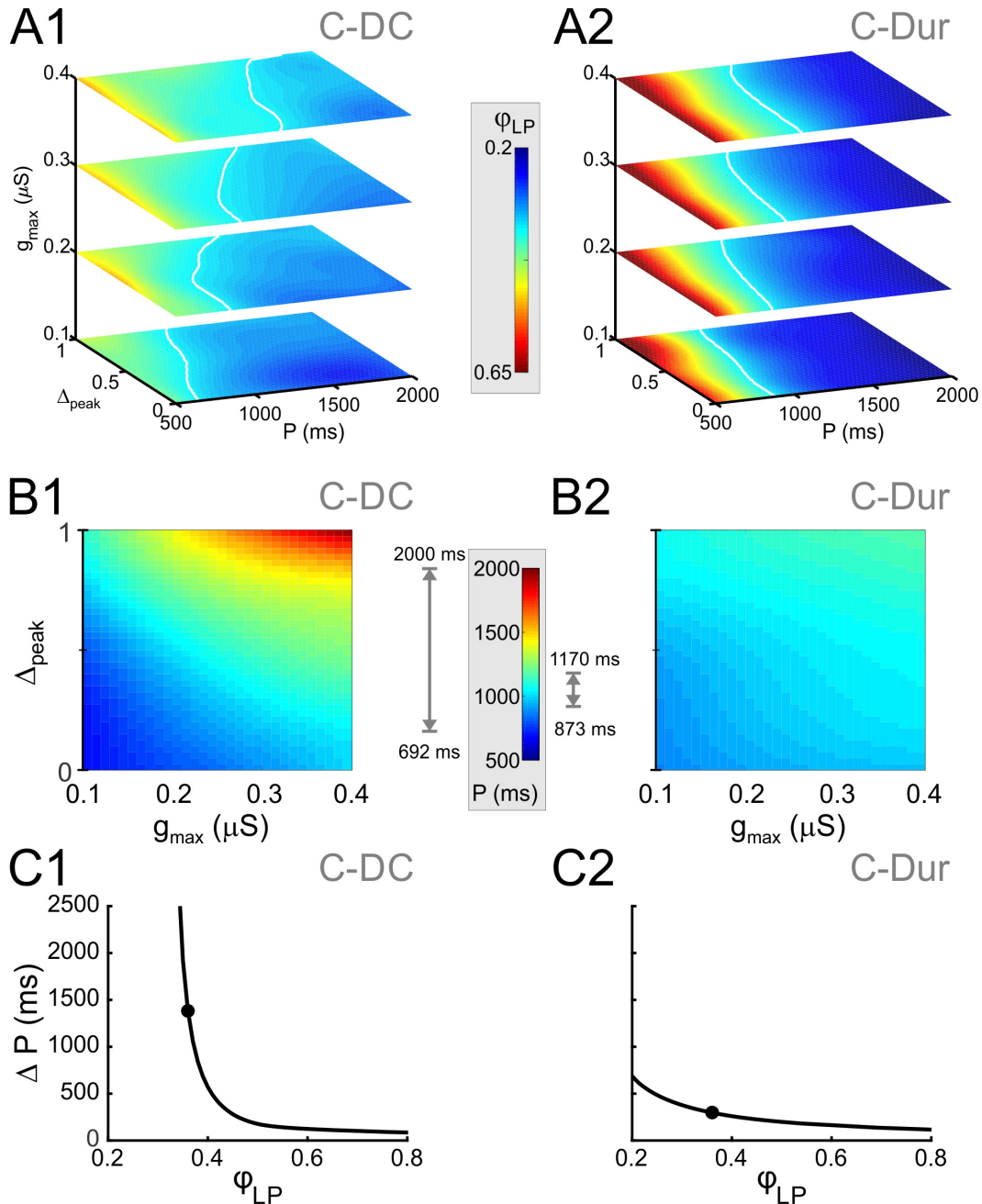
950 range of φ_{LP} drop was greater for the C-Dur case compared to the C-DC case. **B.** φ_{LP}
 951 increases as a function of g_{max} . **B1.** Intracellular recording of an LP neuron showing the
 952 conductance input across three values of g_{max} . **B2.** φ_{LP} for the example shown in B1
 953 plotted as a function of P (for $P = 500$ ms, $\Delta_{peak} = 0.25$) shows a small increase for both
 954 C-Dur and C-DC cases. **B3.** φ_{LP} increased with g_{max} in almost all trials for both C-DC
 955 and C-Dur cases and all values of Δ_{peak} . **C.** φ_{LP} increases as a function of Δ_{peak} . **C1.**
 956 Intracellular recording of the LP neuron showing the conductance input for five values
 957 of Δ_{peak} . **C2.** φ_{LP} for the example neuron in C1 plotted as a function of Δ_{peak} (for $P = 500$
 958 ms, $g_{max} = 0.4$ μ S) for both C-DC and C-Dur cases. **C3.** φ_{LP} increased with Δ_{peak} for
 959 both C-DC and C-Dur cases and all values of g_{max} .



960

961 *Figure 7: Sensitivity analysis shows that φ_{LP} increases more effectively if g_{max} and Δ_{peak}*
 962 *increase together.*

963 **A.** The sensitivity φ_{LP} to local changes in g_{max} and Δ_{peak} was averaged across all values
 964 of P for the C-DC case. The sensitivity was largest if both parameters were increased
 965 together ($g_{max} + \Delta_{peak}$) and smallest if they were varied in opposite directions ($g_{max} -$
 966 Δ_{peak} ; One-Way RM-ANOVA, $p < 0.001$, $F = 3.330$). **B.** The same sensitivity analysis in
 967 the C-Dur case shows similar results (One-Way RM-ANOVA, $p < 0.001$, $F = 2.892$).



968

969

Figure 8: Simultaneous increase of both Δ_{peak} and g_{max} across their range of values can produce phase maintenance across a large P range in the C-DC case and a much smaller P range in the C-Dur case

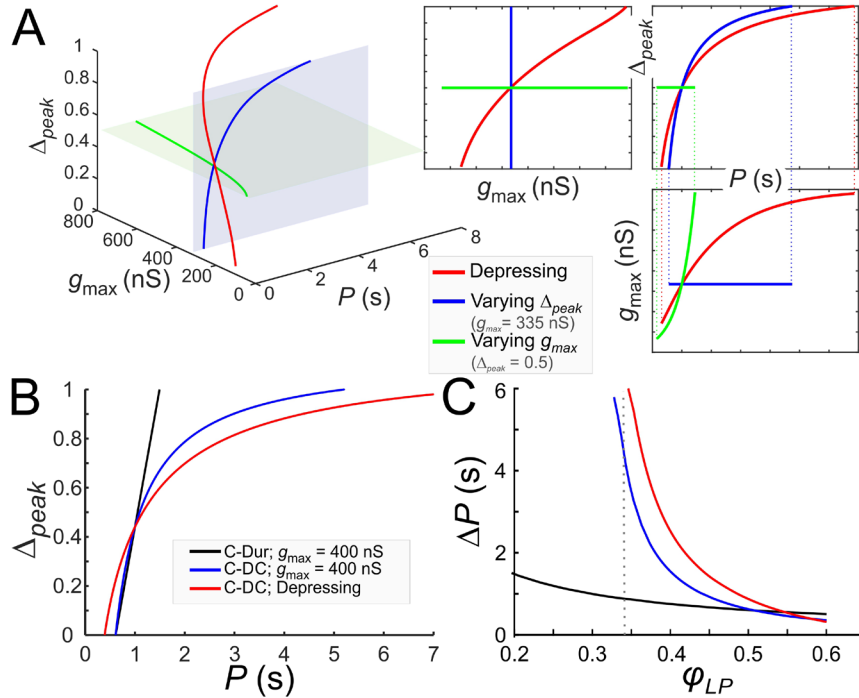
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A. Heat map plots of the function Φ , plotted for the range of values of P and Δ_{peak} and 4 values of g_{max} for the C-DC (**A1**) and C-Dur (**A2**) cases. The white curves show the level set of $\varphi_{LP} = 0.34$, shown as an example of phase constancy. The color maps are interpolated from sampled data (see Methods; $N=9$) **B.** Heat map for the level sets $\varphi_{LP} = 0.34$ for the C-DC (**B1**) and C-Dur (**B2**) cases. Range of colors in each panel indicate

976

977 the range of P values for which φ_{LP} could remain constant at 0.34 for each case, as
978 indicated by the grey arrows on the side of the heatmap color legend. **C.** The range
979 (ΔP) of P values for which φ_{LP} could remain constant at any value between 0.2 and 0.8
980 for the C-DC (**C1**) and C-Dur cases (**C2**). Filled circles show the values shown in panel
981 B. The LP neuron cannot achieve φ_{LP} values below 0.3 in the C-DC case. For φ_{LP}
982 values between 0.3 and ~ 0.65 , the range was larger in C-DC case.



983

984 **Figure 9: Model prediction of the range of phase constancy.**

985 **A.** For the C-DC case, a constant phase of $\phi_{LP} = 0.34$ can be maintained across a
986 range of cycle periods P when g_{max} is constant (at 335 nS; blue plane) and Δ_{peak} varies
987 from 0 to 1 according to Equation (13) (blue), or when Δ_{peak} is fixed (at 0.5; green
988 plane) and g_{max} varies from 200 to 800 nS according to Equation (12). Alternatively,
989 g_{max} and Δ_{peak} can covary to maintain phase, as in a depressing synapse, where g_{max}
990 varies with P according to Equation (11), and Δ_{peak} is calculated for each P and g_{max}
991 value according to Equation (13). As seen in the 2D coordinate-plane projections of the
992 3D graph (right three graphs), the range of P values for which phase constancy is
993 achieved is largest when g_{max} and Δ_{peak} covary (dotted lines show limits of P for phase
994 constancy). The depressing synapse conductance is chosen to be 335 nS at $P = 1$ s.
995 **B, C.** A comparison between the C-DC and C-Dur cases shows that in the latter case a
996 constant phase of ϕ_{LP} can be maintained across a larger range of P values when Δ_{peak}
997 increases with P (and g_{max} is fixed at 400 nS) according to Equation (13). The
998 relationship of Δ_{peak} and P is shown in **B** for $\phi_{LP} = 0.34$. **C** shows the range of P values
999 (ΔP) of cycle periods for which phase remains constant at ϕ_{LP} . If g_{max} also varies with
1000 P , as in a depressing synapse (red; Equation (11)), the range of P values for which
1001 phase is constant is further increased. (Dotted line: $\phi_{LP} = 0.34$.)

1002 **Source Data Files**

1003 **Figure 2-1. File: *Figure2_sourcedata.xlsx***

1004 This Excel file contains 4 sheets, including all measured attributes of the burst-

1005 triggered average current (I_{BTA}) for different IBIs (N=23) as shown in **Fig. 2H-2K**.