## 1 Biophysical mechanisms in the mammalian respiratory oscillator

## 2 re-examined with a new data-driven computational model

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#### 13 Abstract

14 An autorhythmic population of excitatory neurons in the brainstem pre-Bötzinger complex is 15 a critical component of the mammalian respiratory oscillator. Two intrinsic neuronal 16 biophysical mechanisms—a persistent sodium current  $(I_{NaP})$  and a calcium-activated nonselective cationic current  $(I_{CAN})$ —were proposed to individually or in combination generate 17 18 cellular- and circuit-level oscillations, but their roles are debated without resolution. We re-19 examined these roles with a new computational model of an excitatory population with 20 randomly distributed  $I_{NaP}$  and  $I_{CAN}$  conductances and synaptic connections. This model 21 robustly reproduces experimental data showing contrary to previous hypotheses, rhythm 22 generation is independent of  $I_{CAN}$  activation, which instead determines population activity 23 amplitude. The novel insight is that this occurs when  $I_{CAN}$  is primarily activated by neuronal 24 calcium fluxes driven by synaptic mechanisms. Rhythm depends critically on  $I_{NaP}$  in a 25 subpopulation forming the rhythmogenic kernel. The model explains how the rhythm and 26 amplitude of respiratory oscillations involve distinct biophysical mechanisms.

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#### 29 Introduction

Defining cellular and circuit mechanisms generating the vital rhythm of breathing in mammals
remains a fundamental unsolved problem of wide-spread interest in neurophysiology (Richter
& Smith, 2014; Del Negro et al., 2018), with potentially far-reaching implications for

33 understanding mechanisms of oscillatory circuit activity and rhythmic motor pattern generation 34 in neural systems (Marder & Calabrese, 1996; Buzsaki, 2006; Grillner, 2006; Kiehn, 2006). 35 The brainstem pre-Bötzinger complex (pre-BötC) region (Smith et al., 1991) located in the 36 ventrolateral medulla oblongata is established to contain circuits essential for respiratory 37 rhythm generation (Smith et al., 2013; Del Negro et al., 2018), but the operational cellular 38 biophysical and circuit synaptic mechanisms are continuously debated. Pre-BötC excitatory 39 neurons and circuits have autorhythmic properties and drive motor circuits that can be isolated 40 and remain rhythmically active in living rodent brainstem slices in vitro. Numerous 41 experimental and theoretical analyses have focused on the rhythmogenic mechanisms 42 operating in these in vitro conditions to provide insight into biophysical and circuit processes 43 involved, with potential relevance for rhythm generation during breathing in vivo (Feldman & 44 Del Negro, 2006; Richter & Smith, 2014). The ongoing rhythmic activity in vitro has been 45 suggested to arise from either a subset(s) of intrinsically bursting neurons which, through 46 excitatory synaptic interactions, recruit and synchronize neurons within the network 47 (pacemaker-network model) (Rekling & Feldman, 1998; Toporikova & Butera, 2011; Ramirez 48 et al., 2004), or as an emergent network property through recurrent excitation (Jasinski et al., 49 2013) and/or synaptic depression (group pacemaker model) (Rubin et al., 2009).

50 From these previous analyses, involvement of two possible cellular-level biophysical 51 mechanisms have been proposed. One based on a slowly inactivating persistent sodium current 52  $(I_{NaP})$  (Butera et al., 1999a), and the other on a calcium-activated non-selective cation current 53  $(I_{CAN})$  coupled to intracellular calcium  $([Ca]_i)$  dynamics (for reviews see (Rybak et al., 2014; 54 Del Negro et al., 2010), or a combination of both mechanisms (Jasinski et al., 2013). Despite 55 the extensive experimental and theoretical investigations of these sodium- and calcium-based mechanisms, the actual roles of  $I_{NaP}$ ,  $I_{CAN}$  and the critical source(s) of  $[Ca]_i$  transients in the 56 57 pre-BötC are still unresolved. Furthermore, in pre-BötC circuits the process of rhythm 58 generation must be associated with an amplitude of circuit activity sufficient to drive 59 downstream circuits to produce adequate inspiratory motor output. Biophysical mechanisms 60 involved in generating the amplitude of pre-BötC circuit activity have also not been 61 established.

62  $I_{NaP}$  is proposed to mediate an essential oscillatory burst-generating mechanism since 63 pharmacologically inhibiting  $I_{NaP}$  abolishes intrinsic neuronal rhythmic bursting as well as 64 pre-BötC circuit inspiratory activity and rhythmic inspiratory motor output in vitro. Theoretical 65 models of cellular and circuit activity based on  $I_{NaP}$  - dependent bursting mechanisms closely 66 reproduce experimental observations such as voltage-dependent frequency control, spike-67 frequency adaptation during bursts, and pattern formation of inspiratory motor output (Butera 68 et al., 1999b; Pierrefiche et al., 2004; Smith et al., 2007). This indicates the plausibility of  $I_{NaP}$ -

69 dependent rhythm generation.

70 In the pre-BötC, I<sub>CAN</sub> was originally postulated to underlie intrinsic pacemaker-like 71 oscillatory bursting at the cellular level and contribute to circuit-level rhythm generation, since 72 intrinsic bursting in a subset of neurons in vitro was found to be terminated by the  $I_{CAN}$  inhibitor flufenamic acid (FFA) (Pena et al., 2004). Furthermore, inhibition of I<sub>CAN</sub> in the pre-BötC 73 74 reduces the amplitude of the rhythmic depolarization (inspiratory drive potential) driving 75 neuronal bursting and can eliminate inspiratory motor activity in vitro (Pace et al., 2007). I<sub>CAN</sub> 76 became the centerpiece of the group pacemaker model for rhythm generation, in which this 77 conductance was proposed to be activated by inositol trisphosphate (IP3) receptor/ER-78 mediated intracellular calcium fluxes initiated via glutamatergic metabotropic receptormediated signaling in the pre-BötC excitatory circuits. The molecular correlate of  $I_{CAN}$  was 79 80 postulated to be the transient receptor potential channel M4 (TRPM4) (Mironov 2008; Pace et 81 al. 2007)- one of the two known Ca<sup>2+</sup>-activated TRP channels (Guinamard et al., 2010; Ullrich 82 et al., 2005). TRPM4 has now been identified by immunolabeling and RNA expression 83 profiling in pre-BötC inspiratory neurons in vitro (Koizumi et al., 2018).

Investigations into the sources of intracellular  $Ca^{2+}$  activating  $I_{CAN}$ /TRPM4 suggested 84 that (1) somatic calcium transients from voltage-gated sources do not contribute to the 85 86 inspiratory drive potential (Morgado-Valle et al., 2008), (2) IP3/ER-mediated intracellular Ca<sup>2+</sup> 87 release does not contribute to inspiratory rhythm generation in vitro), and (3) in the dendrites calcium transients may be triggered by excitatory synaptic inputs and travel in a wave 88 89 propagated to the soma (Mironov 2008). Theoretical studies have demonstrated the plausibility 90 of  $[Ca]_i$ - $I_{CAN}$  dependent bursting (Rubin et al., 2009; Toporikova & Butera, 2011), however these models omit  $I_{NaP}$  and/or depend on additional unproven mechanisms to generate 91 92 intracellular calcium oscillations to provide burst termination, such as IP3 dependent calcium-93 induced calcium release (Toporikova & Butera, 2011), partial depolarization block of action 94 potentials (Rubin, et al., 2009), and the Na<sup>+</sup>/K<sup>+</sup> pump (Jasinski et al., 2013). Surprisingly, 95 pharmacological inhibition of  $I_{CAN}$ /TRPM4 has recently been shown to produce large 96 reductions in the amplitude of pre-BötC inspiratory neuron population activity without 97 significant perturbations of inspiratory rhythm (Koizumi et al., 2018). These new observations 98 re-define the role of  $I_{CAN}$ , and require theoretical re-examination of pre-BötC neuronal

99 conductance mechanisms and network dynamics, particularly how rhythm generation 100 mechanisms can be independent of  $I_{CAN}$ -dependent mechanisms generating the amplitude of 101 network activity.

102 In this theoretical study, we examine the role of  $I_{CAN}$  in the pre-BötC by considering 103 two plausible mechanisms of intracellular calcium fluxes: (1) from voltage-gated, and (2) from 104 synaptically activated sources. We deduce that  $I_{CAN}$  is primarily activated by calcium transients 105 that are coupled to rhythmic excitatory synaptic inputs originating from  $I_{NaP}$  dependent bursting inspiratory neurons. Additionally, we show that  $I_{CAN}$  contributes to the inspiratory 106 107 drive potential by mirroring the excitatory synaptic current. Our new model explains the recent 108 experimental observations obtained from in vitro neonatal rodent slices isolating the pre-BötC, 109 showing large reductions in circuit activity amplitude by inhibiting  $I_{CAN}$ /TRPM4 without 110 perturbations of inspiratory rhythm generation in pre-BötC excitatory circuits in vitro. The 111 model supports the novel concept that  $I_{CAN}$  activation in a subpopulation of pre-BötC excitatory 112 neurons are critically involved in amplifying synaptic drive from a subset of neurons whose 113 rhythmic bursting is critically dependent on  $I_{NaP}$  and forms the kernel for rhythm generation 114 in vitro. The model shows how the functions of generating the rhythm and amplitude of 115 inspiratory oscillations in pre-BötC excitatory circuits are determined by distinct biophysical 116 mechanisms.

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119 **Results** 

# 120 $\overline{g}_{CAN}$ Variation has Opposite Effects on Amplitude and Frequency of Network Bursting121in the $Ca_V$ and $Ca_{Svn}$ Models

122 Recent experimental work (Koizumi et al., 2018) has demonstrated that pharmacological 123 inhibition of  $I_{CAN}$ /TRPM4 in the pre-BötC in slices from in vitro neonatal mouse/rat slice 124 preparations, strongly reduces the amplitude of (or completely eliminates) the inspiratory 125 hypoglossal (XII) motor output, as well as the amplitude of pre-BötC excitatory circuit activity 126 that is highly correlated with the decline of XII activity, while having little effect on inspiratory 127 burst frequency. Here, we systematically examine the relationship between I<sub>CAN</sub> conductance 128  $(\bar{g}_{CAN})$  on amplitude and frequency of circuit activity for voltage-gated  $(Ca_V)$  and synaptically 129 activated sources ( $Ca_{Syn}$ ) of intracellular calcium. We found that that reduction of  $\bar{g}_{CAN}$  drives 130 opposing effects on circuit activity amplitude and frequency that are dependent on the source 131 of intracellular calcium transients (Fig. 1). In the  $Ca_{V}$  network, where calcium influx is

132 generated exclusively from voltage-gated calcium channels, increasing  $\bar{g}_{CAN}$  has no effect on 133 amplitude but increases the frequency of network oscillations (Fig. 1A, C, D). Conversely, in 134 the  $Ca_{Syn}$  network where calcium influx is generated exclusively by excitatory synaptic input, 135 increasing  $\bar{g}_{CAN}$  strongly increases the amplitude and slightly decreases the oscillation 136 frequency (Fig. 1B, C, D).

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139 Figure 1. Manipulations of  $\bar{g}_{CAN}$  in the  $Ca_V$  and  $Ca_{Syn}$  networks produce opposite effects on network activity 140 amplitude (spikes/s) and frequency. (A & B) Histograms of neuronal firing and voltage traces for pacemaker and 141 follower neurons in the  $Ca_V$ , and  $Ca_{Syn}$  networks with linearly increasing  $\bar{g}_{CAN}$ . (C) Plot of  $\bar{g}_{CAN}$  vs network 142 activity amplitude for the  $Ca_V$  and  $Ca_{SVR}$  networks in A and B. (D) Plot of  $\bar{g}_{CAN}$  vs network frequency for the  $Ca_V$ 143 and  $Ca_{Syn}$  networks in A and B.  $Ca_V$  network parameters:  $\bar{g}_{Ca} = 1.0$  (nS),  $P_{Ca} = 0.0$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.05$ 144 0.2 (nS).  $Ca_{Syn}$  network parameters:  $\bar{g}_{Ca} = 0$  (nS),  $P_{Ca} = 0.01$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.2$  (nS).

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#### 146 Effects of Subthreshold Activation of I<sub>CAN</sub> on Network Frequency

147 In  $I_{NaP}$ -dependent bursting neurons in the pre-BötC, bursting frequency depends on their 148 excitability (i.e., baseline membrane potential) which can be controlled in different ways, e.g. 149 by directly injecting a depolarizing current (Smith et al., 1991) or varying the conductance 150 and/or reversal potentials of some ionic channels (Butera et al., 1999a). Due to their relatively 151 short duty cycle, the bursting frequency in these neurons is largely determined by the interburst 152 interval, defined as the time between the end of one burst and the start of the next. During the 153 burst,  $I_{NaP}$  slowly inactivates resulting in burst termination and abrupt hyperpolarization of the 154 membrane. The interburst interval is then determined by the amount of time required for  $I_{NaP}$ 155 to recover from inactivation and return the membrane potential back to the threshold for burst 156 initiation. This process is governed by the kinetics of  $I_{NaP}$  inactivation gating variable  $h_{NaP}$ . 157 Higher neuronal excitability reduces the value of  $h_{NaP}$  required to initiate bursting. 158 Consequently, the time required to reach this value is decreased, which results in a shorter 159 interburst interval and increased frequency.

160 To understand how changing  $\bar{g}_{CAN}$  affects network bursting frequency we quantified 161 the values of  $h_{NaP}$  averaged over all pacemaker neurons immediately preceding each network 162 burst and, also, the average  $I_{CAN}$  values between the bursts in the  $Ca_V$  and  $Ca_{syn}$  networks (Fig. 2). In the  $Ca_V$  network,  $I_{Ca}$  as modeled remains residually activated between the bursts thus 163 164 creating the background calcium concentration which partially activates  $I_{CAN}$ . Therefore, 165 between the bursts  $I_{CAN}$  functions as a depolarizing leak current. Consistently, we found that in the  $Ca_V$  network increasing  $\bar{g}_{CAN}$  increases  $I_{CAN}$  (Fig. 2A) progressively depolarizing the 166 167 network, which reduces the  $h_{NaP}$  threshold for burst initiation (Fig. 2B) and, thus, increases 168 network frequency (Fig. 1D).

- 169 In the  $Ca_{syn}$  model the intracellular calcium depletes entirely during the interburst 170 interval. Consequently, increasing  $\bar{g}_{CAN}$  has no effect on  $I_{CAN}$  (Fig. 2A) and frequency is 171 essentially unaffected (Fig. 1D).
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Figure 2 Calcium source and  $\bar{g}_{CAN}$ -dependent effects on cellular properties regulating network frequency for the simulations presented in Fig. 1. (A) Average magnitude of  $I_{CAN}$  in pacemaker neurons during the interburst interval for the  $Ca_V$  (red) and  $Ca_{Syn}$  (blue) networks. (B) Average inactivation of the burst generating current  $I_{NaP}$  in pacemaker neurons immediately preceding each network burst as a function of  $\bar{g}_{CAN}$  for the voltage-gated and synaptic calcium networks.  $Ca_V$  Network Parameters:  $\bar{g}_{Ca} = 1.0$  (*nS*),  $P_{Ca} = 0.0$ ,  $P_{Syn} = 0.05$  and  $W_{max} =$ 0.2 (*nS*).  $Ca_{Syn}$  network parameters:  $\bar{g}_{Ca} = 0$  (*nS*),  $P_{Ca} = 0.01$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.2$  (*nS*).

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#### 181 Changes in Network Activity Amplitude Are Driven by Recruitment of Neurons

182 Network activity amplitude is defined as the total number of spikes produced by the 183 network per a time bin. Consequently, changes in amplitude can only occur by increasing the 184 number of neurons participating in bursts (recruitment) and/or increasing the firing rate of the 185 recruited neurons. To analyze changes in amplitude, we quantified the number of recruited 186 neurons (Fig. 3A) and the average spike frequency in recruited neurons (Fig. 3B) as a function of  $\bar{g}_{CAN}$  for both network models. In the  $Ca_V$  network, increasing  $\bar{g}_{CAN}$  increases the number 187 188 of recruited neurons (Fig. 3A), but decreases the average spiking frequency in recruited neurons 189 (Fig. 3B) which, together result in no change in amplitude (Fig. 1C). In the  $Ca_{syn}$  network, 190 increasing  $\bar{g}_{CAN}$  strongly increases the number of recruited neurons (Fig. 3A) and increases the 191 spike frequency of recruited neurons (Fig. 3B) resulting in a large increase in network activity 192 amplitude (Fig. 3C).



194 Figure 3 Calcium source and  $\bar{g}_{CAN}$ -dependent effects on cellular properties regulating network activity amplitude 195 for the simulations presented in Fig. 1. (A) Number of recruited neurons in the modeled population of 100 neurons 196 as a function of  $\bar{g}_{CAN}$  for voltage-gated and synaptic calcium sources. The number of recruited neurons is defined 197 as the peak number of spiking neurons per bin during a network burst. (B) Average spiking frequency of recruited 198 neurons as a function of  $\bar{g}_{CAN}$  for the voltage-gated and synaptic calcium mechanism. Average spiking frequency 199 is defined the number of spikes per bin divided by the number of recruited neurons. The parameters used in these 200 simulations are:  $Ca_V$ :  $\bar{g}_{Ca} = 1.0$  (*nS*),  $P_{Ca} = 0.0$ ,  $P_{Syn} = 0.05$  and W = 0.2 (*nS*).  $Ca_{Syn}$ :  $\bar{g}_{Ca} = 0$  (*nS*),  $P_{Ca} = 0$ 201 0.01,  $P_{Syn} = 0.05$  and W = 0.2 (*nS*).

## 203 Manipulating $\overline{g}_{CAN}$ in the $Ca_{Syn}$ Model is Qualitatively Equivalent to Changing the 204 Strength of Synaptic Interactions

205 Since changes in  $\bar{g}_{CAN}$  in the  $Ca_{Svn}$  model primarily affect network activity amplitude 206 through recruitment of follower neurons, and the network activity amplitude strongly depends 207 on the strength of synaptic interactions, we next examined the relationship between  $\bar{g}_{CAN}$ , 208 synaptic strength and network activity amplitude and frequency (Fig. 4). Synaptic strength is 209 defined as the number of neurons multiplied by the connection probability multiplied by the average weight of synaptic connections  $(N \cdot P_{Syn} \cdot \frac{1}{2}W_{max})$ . We found that the effects of varying 210  $\bar{g}_{CAN}$  or the synaptic strength on network activity amplitude and frequency are qualitatively 211 212 equivalent in the  $Ca_{Svn}$  network which is indicated by symmetry of the heat plots (across the X=Y line) in Fig. 4A, B. We further investigated and compared the effect of reducing  $\bar{g}_{CAN}$  or 213 214 the synaptic strength on network activity amplitude and frequency as well as the effects on the 215 recruitment of follower neurons (Fig. 4C-F). To make this comparison, picked a starting point 216 in the 2D parameter space between  $\bar{g}_{CAN}$  and synaptic strength where the network is bursting. 217 Then in separate simulations we linearly reduced either  $\bar{g}_{CAN}$  or the synaptic strength to zero. 218 We show that reducing either  $\bar{g}_{CAN}$  or the synaptic strength have very similar effects on 219 network activity amplitude and frequency (Fig. 4C, D). Furthermore, the effect on follower 220 neurons in both cases is nearly identical (Fig. 4E, F). Reducing either  $\bar{g}_{CAN}$  or the synaptic 221 strength decreases the excitatory input to follower neurons during network oscillations which 222 is a major component of the inspiratory drive potential. Therefore, in the  $Ca_{Syn}$  network, 223 manipulations of  $\bar{g}_{CAN}$  will affect the strength of the inspiratory drive potential in follower 224 neurons in a way that is equivalent to changing the synaptic strength of the network. In contrast, 225 manipulations of  $\bar{g}_{CAN}$  in the  $Ca_V$  network will only slightly affect the inspiratory drive 226 potential in follower neurons due to changes in the average firing rate of active neurons (see 227 Fig. 3B).



230 Figure 4 Manipulations of synaptic strength  $(N \cdot P_{Syn} \cdot \frac{1}{2}W_{max})$  and  $\bar{g}_{CAN}$  have equivalent effects on network 231 activity amplitude, frequency and recruitment of follower neurons. (A & B) Relationship between  $\bar{g}_{CAN}$ , synaptic 232 strength and the amplitude and frequency in the  $Ca_{Syn}$  network. Notice the symmetry about the X=Y line in panels 233 A and B, which, indicates that changes in  $\bar{g}_{CAN}$  and or synaptic strength are qualitatively equivalent. Synaptic 234 strength was changed by varying  $W_{max}$ . (C) Relationship between network activity amplitude and the reduction 235 of  $\bar{g}_{CAN}$  (BLUE) or synaptic strength (GREEN). (D) Relationship between network frequency and the reduction 236 of  $\bar{g}_{CAN}$  (BLUE) or synaptic strength (GREEN). (E & F) Decreasing  $\bar{g}_{CAN}$  or synaptic strength de-recruits 237 follower neurons by reducing the inspiratory drive potential, indicated by the amplitude of subthreshold 238 depolarization, right traces. The solid blue and green lines in panels A and B represent the location in the 2D 239 parameter space of the corresponding blue and green curves in C and D. The action potentials in the right traces 240 of E and F are truncated to show the change in neuronal inspiratory drive potential. The parameters used for these 241 simulations are  $Ca_{Syn}$ :  $\bar{g}_{Ca} = [0,0]$ ,  $P_{Ca} = 0.01$ ,  $P_{Syn} = 1.0$  and  $W_{max} = var$ 

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#### 243 Robustness of Amplitude and Frequency Effects

244 We also examined if the effects are conserved in both the  $Ca_V$  and  $Ca_{Syn}$  networks over 245 a range of network parameters. To test this, we investigated the dependence of network activity 246 amplitude and frequency on  $\bar{g}_{CAN}$  and average synaptic weight for  $Ca_{Syn}$  and  $Ca_{V}$  networks with high  $(P_{Syn} = 1)$  and low  $(P_{Syn} = 0.05)$  connection probabilities, and high  $(g_{Ca} =$ 247  $0.1 nS, P_{Ca} = 0.1$ , medium ( $g_{Ca} = 0.01 nS, P_{Ca} = 0.01$ ) and low ( $g_{Ca} = 0.001 nS, P_{Ca} = 0.01 nS, P_{Ca} = 0.01 nS$ ) 248 249 0.005) strengths of calcium sources (Figures 5 and 6). We found that changing the synaptic 250 connection probability and changing the strength of the calcium sources has no effect on the 251 general relationship between  $\bar{g}_{CAN}$  and the amplitude or frequency of bursts in the  $Ca_V$  or  $Ca_{SVN}$ 252 networks. In other words, the general effect of increasing  $\bar{g}_{CAN}$  on amplitude and frequency is 253 conserved in both networks regardless of the synaptic connection probability or strength of the 254 calcium sources. Increasing the strength of the calcium sources does, however, affects the range 255 of possible  $\bar{g}_{CAN}$  values where both networks produce rhythmic activity.

To summarize, in the  $Ca_V$  model, increasing  $\bar{g}_{CAN}$  increases frequency, through increased excitability but has no effect on amplitude. In contrast, in the  $Ca_{Syn}$  model, increasing  $\bar{g}_{CAN}$  slightly decreases frequency and increases amplitude. In this case, increasing  $\bar{g}_{CAN}$  acts as a mechanism to increase the inspiratory drive potential and recruit previously silent neurons. Additionally, these features of the  $Ca_V$  and  $Ca_{Syn}$  models are robust and conserved across a wide range of network parameters.

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Figure 5 Robustness of amplitude and frequency effects to changes in  $\bar{g}_{CAN}$  and synaptic strength in the  $Ca_V$ network for 'high'(left), 'medium' (middle) and 'low' (right) conductance of the voltage-gated calcium channel  $I_{Ca}$  as well as 'high'(top) and 'low' (bottom) network connection probabilities. Amplitude and frequency are indicated by color (scale bar at right). Black regions indicate tonic network activity.

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Figure 6 Robustness of amplitude and frequency effects to changes in  $\bar{g}_{CAN}$  and synaptic strength in the  $Ca_{Syn}$ network for 'high'(left), 'medium' (middle) and 'low' (right) calcium conductance in synaptic currents as well as 'high'(top) and 'low' (bottom) network connection probabilities. Amplitude and frequency are indicated by color (scale bar at right). Black regions indicate tonic network activity.

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#### 281 Intracellular Calcium Transients Activating I<sub>CAN</sub> Primarily Result from Synaptically-

282 Activated Sources

In experiments where  $I_{CAN}$  was blocked by bath application of FFA or 9-phenanthrol (Koizumi et al., 2018), the amplitude of network oscillations was strongly reduced and their frequency remained unchanged. Our model revealed that the effects of  $I_{CAN}$  blockade on amplitude and frequency depend on the source(s) of intracellular calcium (see Figs. 1 and 2). If the calcium influx is exclusively voltage-gated, our model predicts that  $I_{CAN}$  blockade will 288 have no effect on amplitude but reduce the frequency. In contrast, if the calcium source is 289 exclusively synaptically gated, our model predicts that blocking  $I_{CAN}$  will strongly reduce the 290 amplitude and slightly increase the frequency. Therefore, a multi-fold decrease in amplitude, 291 seen experimentally, is consistent with the synaptically driven calcium influx mechanisms, 292 while constant bursting frequency may be due to calcium influx through both voltage- and 293 synaptically gated channels. Following predictions above, to reproduce experimental data, we 294 incorporated both mechanisms in the model and inferred their individual contributions by 295 finding the best fit. We found that the best match is observed (Fig. 7) if synaptically mediated 296 and voltage gated calcium influxes comprise about 95% and 5% of the total calcium influx, 297 respectively.



300 Figure 7. Experimental and simulated pharmacological blockade of I<sub>CAN</sub> by (A & C) 9-phenanthrol and (B & D) 301 flufenamic acid (FFA). Both voltage-gated and synaptic sources of intracellular calcium are included. 302 Experimental blockade of I<sub>CAN</sub> (black) by 9-phenanthrol and FFA significantly reduce the (A & B) amplitude of 303 network oscillations while having little effect on (C & D) frequency. The black line represents the mean and the 304 gray is the S.E.M. of experimental JXII output recorded from neonatal rat brainstem slices in vitro, reproduced 305 from Koizumi et al., 2018. Simulated blockade of  $I_{CAN}$  (red) closely matches the reduction in (A & B) amplitude 306 of network oscillations and slight decrease in (C & D) frequency seen with 9-phenanthrol and FFA. Simulated 307 and experimental blockade begins at the vertical dashed line. Blockade was simulated by exponential decay of 308  $\bar{g}_{CAN}$  with the following parameters: 9-phenanthrol:  $\gamma_{Block} = 0.85$ ,  $\tau_{Block} = 357s$ ; FFA:  $\gamma_{Block} = 0.92$ ,  $\tau_{Block} = 0.92$ 309 415s. The network parameters are:  $\bar{g}_{Ca} = 0.00175 (nS)$ ,  $P_{Ca} = 0.0275$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.096 (nS)$ . 310

#### 311 $I_{NaP}$ -dependent and $[Ca]_i$ - $I_{CAN}$ Sensitive Intrinsic Bursting

312 In our model, we included  $I_{NaP}$ ,  $I_{CAN}$  as well as voltage-gated and synaptic mechanisms of  $Ca^{2+}$  influx. Activation of  $I_{CAN}$  by  $Ca_{Syn}$  is the equivalent mechanism used in other 313 314 computational group-pacemaker models (Rubin et al., 2009; Song et al., 2015). Burst 315 generation and termination in our model, however, are dependent on  $I_{NaP}$  (Butera et al., 1999a). 316 We investigated the sensitivity of intrinsic bursting in our model to  $I_{NaP}$  and calcium channel 317 blockade (Fig. 8). Intrinsic bursting was identified in neurons by zeroing the synaptic weights 318 to simulate synaptic blockade.  $I_{NaP}$  and  $I_{Ca}$  blockade was simulated by setting  $\bar{g}_{NaP}$  and  $\bar{g}_{Ca}$ 319 to 0. We found that after decoupling the network ( $W_{max} = 0$ ) a subset of neurons remained rhythmically active (7%) and that these were all neurons with a high  $I_{NaP}$  conductance. In these 320 321 rhythmically active neurons, bursting was abolished in all neurons by  $I_{NaP}$  blockade. Interestingly,  $I_{Ca}$  blockade applied before  $I_{NaP}$  blockade abolished intrinsic bursting in 2 of the 322 323 7 neurons and  $I_{NaP}$  blockade applied afterwards abolished intrinsic bursting in the remaining 324 5 neurons. Although only one rhythmogenic ( $I_{NaP}$ -based) mechanism exists in this model, 325 bursting in a subset of these neurons is calcium sensitive. In calcium-sensitive bursters, Ca<sup>2+</sup> 326 blockade abolishes bursting by reducing the intracellular calcium concentration and, hence, 327  $I_{CAN}$  activation, which ultimately reduces excitability.



330 Figure 8  $I_{NaP}$ -dependent and  $Ca^{2+}$  sensitive intrinsic bursting. (A) From left to right, intrinsic bursters are first 331 identified by blocking synaptic connections. Then, calcium sensitive neurons are silenced and identified by  $I_{Ca}$ 332 blockade. The remaining neurons are identified as sensitive to  $I_{NaP}$  block. Top traces show the network output and 333 Cell ID vs.  $\bar{g}_{NaP}$  scatter plots identify silent and bursting neurons under each condition. (B)  $I_{NaP}$  blockade after 334 synaptic blockade eleminates bursting in all neurons. Therefore, all intrinsic bursters are  $I_{NaP}$  dependent. (C) 335 Identification of calcium sensitive and  $I_{NaP}$ -dependent as well as calcium insensitive and  $I_{NaP}$ -dependent intrinsic 336 bursters. Notice that only the neurons with the highest value of  $\bar{g}_{NaP}$  are intrinsic bursters and that a subset of 337 these neurons are sensitive to calcium blockade but all are dependent on  $I_{NaP}$ . The network parameters are:  $\bar{g}_{Ca}$  = 338  $0.00175 (nS), P_{Ca} = 0.0275, P_{Syn} = 0.05 \text{ and } W_{max} = 0.096 (nS).$ 

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#### 340 The Rhythmogenic Kernel

Our simulations have shown that the primary role of  $I_{CAN}$  is amplitude but not oscillation frequency modulation with little or no effect on network activity frequency. Here we examined the neurons that remain active and maintain rhythm after  $I_{CAN}$  blockade (Fig. 9). We found that the neurons that remain active are primarily neurons with the highest  $\bar{g}_{NaP}$  and that bursting in these neurons is dependent on  $I_{NaP}$ . Some variability exists and neurons with relatively low  $\bar{g}_{NaP}$  value can remain active due to synaptic interactions while a neuron with a slightly higher  $\bar{g}_{NaP}$  without sufficient synaptic input may become silent. These neurons, that remain active

348 after compete blockade of  $I_{CAN}$ , form a  $I_{NAP}$ -dependent kernel of a rhythm generating circuit.

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Figure 9  $I_{CAN}$  blockade reveals an  $I_{NaP}$ -dependent rhythmogenic kernel. The top traces show the network output at baseline, after  $I_{CAN}$  blockade and  $I_{NaP}$  blockade. The bottom Cell ID vs.  $\bar{g}_{NaP}$  scatter plots identify silent and bursting neurons in each conditon. Notice that only neurons with relitively high  $\bar{g}_{NaP}$  remain active after  $I_{CAN}$ block. The network parameters used are:  $\bar{g}_{Ca} = 0.00175$  (*nS*),  $P_{Ca} = 0.0275$ ,  $P_{Syn} = 0.05$  and  $W_{max} =$ 0.096 (*nS*).

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# Network Activity Amplitude and Neuronal Spiking Frequency Correlate with Intracellular Calcium Transients

359 Dynamic calcium imaging has been utilized to assess activity of individual pre-BötC excitatory 360 neurons and populations of these excitatory neurons in vitro during pharmacological inhibition 361 of  $I_{CAN}$ /TRPM4 (Koizumi et al., 2018). In such imaging-based analyses of excitatory network 362 activity, the sources of calcium transients are not precisely known, but the calcium transients 363 are assumed to correlate with the spiking frequency of neurons and provide measurements that 364 are correlated with circuit activity as a whole, as the experimental data suggest. We compared 365 the network activity characterized by the average intracellular calcium concentration and the 366 network firing. We found that for a single network burst, the average intracellular calcium 367 concentration and network firing are highly correlated (Fig. 10A).

We also compared intracellular calcium transients and spiking frequency in individual pacemaker and follower neurons (Fig. 10B). Pacemaker neurons drive network oscillations and begin firing before the rest of the network. As a result, in the pacemaker neurons the actionpotential generation precedes calcium influx. Consequently, in pacemaker neurons the  $[Ca]_i$ and firing rate correlate poorly in the first portion of the burst. In follower neurons bursting is dependent on synaptic input and recruitment through  $I_{CAN}$  activation. Thus, in follower neurons, which make up most neurons, the  $[Ca]_i$  and action potential firing rate are highly

#### 375 correlated.



377 Figure 10 Correlations between network and neuronal spiking activity and intracellular calcium influx. A) 378 Comparison of network activity amplitude and the average intracellular calcium concentration across a single 379 network burst (left). Network activity amplitude and the average network intracellular calcium concentration are 380 highly correlated (right). B) Comparison of spiking frequency and intracellular calcium concentration for a single 381 burst in a typical follower and pacemaker neuron (left). Spiking frequency is highly correlated with the 382 intracellular calcium transient in follower neurons (right, red). In the pacemaker neuron calcium influx occurs 383 after spiking is initiated, consequently the correlation between spiking frequency and the intracellular calcium 384 concentration is poor at the beginning of the bursts (right, black). The network parameters used are:  $\bar{g}_{Ca}$  = 385 0.00175 (*nS*),  $P_{Ca} = 0.0275$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.096$  (*nS*).

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## 387 Low Network Connection Probability Increases the Variability of the $\Delta[Ca]_i$ in 388 Individual Neurons During $I_{CAN}$ Blockade

In our model, synaptically mediated calcium influx into the cell is proportional to the total synaptic current through its membrane. The synaptic current is determined by the number of synapses, the strength of the connections (synaptic weights) and the firing rates of the pre-

392 synaptic neurons. The synaptic connection probability  $(P_{Syn})$  together with the total number of 393 neurons in the network determines the average number of connections each neuron receives. 394 We examined the relationship between connection probability and the change in  $[Ca]_i$  during 395 simulated  $I_{CAN}$  blockade (Fig. 11). We found that the  $P_{Syn}$  has no effect on the relationship 396 between amplitude, frequency or calcium transients at the network level provided that the 397 synaptic strength remains constant  $(N \cdot P_{Syn} \cdot \frac{1}{2}W_{max} = const)$  (Fig. 11A, B). Additionally, regardless of  $P_{Svn}$ , the network activity amplitude and average intracellular calcium 398 399 concentration are highly correlated.  $P_{Svn}$  does however affects the change in the peak  $[Ca]_i$  in individual neurons. In a network with a high connection probability  $(P_{Syn} = 1)$  the synaptic 400 401 current/calcium transient is nearly identical for all neurons and therefore the change in  $[Ca]_i$ 402 during  $I_{CAN}$  blockade is approximately the same for each neuron (Fig. 11C). In a sparsely 403 connected network the synaptic current and calcium influx are more variable and reflect the 404 heterogeneity in spiking frequency of the pre-synaptic neurons (Fig. 11D). Interestingly, in a network with low connection probability ( $P_{Syn} < 0.1$ ), the peak  $[Ca]_i$  transient in some 405 406 neurons increases when  $I_{CAN}$  is blocked (Fig. 11E). 407





409 Figure 11 Low network connection probability increases the variability of changes in the peak  $[Ca]_i$  in individual 410 neurons during I<sub>CAN</sub> blockade. (A & B) Effect of I<sub>CAN</sub> blockade on network activity amplitude, network calcium 411 amplitude and frequency for network connection probabilities A) P = 1 and B) P = 0.05. (C & D) Effect of  $I_{CAN}$ 412 blockade on changes in the magnitude of peak cellular calcium transients for network connection probabilities C) 413  $P_{Syn} = 1$  and D)  $P_{Syn} = 0.05$ . E) Maximum, minimum and average change in the peak intracellular calcium 414 transient of individual neurons as a function of synaptic connection probability. All curves in A through E are 415 normalized to their baseline values. Synaptic weight was adjusted to keep the average synaptic strength  $(N \cdot P_{Syn} \cdot$ 416  $\frac{1}{2}W_{max} = const$ ) constant. Notice that lowering the synaptic connection probability increases the variability in the 417 peak intracellular calcium concentration during ICAN blockade. Interestingly, for connection probabilities below

418 approximately 5%, blocking  $\bar{g}_{CAN}$  can increase the peak calcium transient in a small subset of neurons. The 419

network parameters used are:  $\bar{g}_{Ca} = 0.00175$  (*nS*) and  $P_{Ca} = 0.0275$  and  $W_{max} = var$ .

420

#### 4. Discussion 421

422 Establishing cellular and circuit mechanisms generating the rhythm and amplitude of 423 respiratory oscillations in the mammalian brainstem pre-BötC has remained an unsolved 424 problem of wide-spread interest in neurophysiology since this structure, essential for breathing 425 to support mammalian life, was discovered nearly three decades ago (Smith et al., 1991). Our 426 objective in this theoretical study was to re-examine and further define contributions of two of 427 the main currently proposed neuronal biophysical mechanisms operating in pre-BötC 428 excitatory circuits, specifically mechanisms involving  $I_{CAN}$  activated by neuronal calcium 429 fluxes and voltage-dependent  $I_{NaP}$  in the circuit neurons. While these sodium- and calcium-430 based mechanisms have been studied extensively over the past two-decades and shown 431 experimentally to be integrated in pre-BötC circuits, their actual roles in circuit operation are 432 continuously debated and unresolved. Both mechanisms have been proposed to be 433 fundamentally involved in rhythm generation either separately or in combination, as plausibly 434 shown from previous theoretical modeling studies. Furthermore, the process of rhythm 435 generation in pre-BötC circuits must be associated with an amplitude of excitatory circuit 436 activity sufficient to drive downstream circuits to produce adequate respiratory motor output. 437 Biophysical mechanisms involved in generating excitatory population activity amplitude have 438 also not been established. Our analysis is motivated by the recent experimental observations 439 obtained from neonatal rodent slices isolating pre-BötC circuits in vitro that inhibition of the 440 endogenously active  $I_{CAN}$ /TRPM4 strongly reduces the amplitude of network oscillations 441 within pre-BötC circuits but have little effect on oscillation frequency (Koizumi et al., 2018). 442 This is contrary to the proposed  $I_{CAN}$ -based models for rhythm generation in the isolated pre-443 BötC and indicates a fundamentally different functional organization of pre-BötC circuits, in 444 terms of oscillatory frequency and amplitude generation, that needs to be defined.

445 We accordingly analyzed the role of  $I_{CAN}$  and possible sources of intracellular calcium 446 transients activating this conductance and found that the effect of simulated  $I_{CAN}$  blockade on 447 amplitude and frequency is highly dependent on the source(s) of intracellular calcium, which 448 is also a central issue to be resolved. In the case where  $Ca_{Syn}$  is the primary intracellular 449 calcium source, I<sub>CAN</sub> blockade generates a large reduction in network activity amplitude. In 450 contrast, when  $Ca_V$  is the only intracellular calcium source,  $I_{CAN}$  blockade has little effect on 451 network activity amplitude and primarily affects the population bursting frequency that is

caused by decreased excitability. Additionally, we show that activation of  $I_{CAN}$  by  $Ca_{Syn}$ 452 453 functions as a mechanism to augment the inspiratory drive potential and amplitude of 454 population activity, and that this effect is similar to increasing the synaptic coupling strength 455 within the network. Therefore, in the case of  $Ca_{Syn}$ , blockade of  $I_{CAN}$  reduces the inspiratory 456 drive potential causing de-recruitment of non-rhythmogenic follower neurons and reduction of 457 network activity amplitude. In a model where  $I_{CAN}$  is activated by both  $Ca_V$  and  $Ca_{Syn}$  with 458 contributions of 5% and 95% respectively, we show that simulated blockade of  $I_{CAN}$  generates 459 a large reduction in network population activity amplitude and a slight decrease in frequency. 460 This closely reproduces experimental blockade of  $I_{CAN}$ /TRPM4 by either 9-phenanthrol or FFA 461 (Fig. 7). Finally, we showed that the change in the peak calcium transients for individual neurons during  $I_{CAN}$  blockade match experimental data particularly at relatively low connection 462 probabilities ( $P_{Svn} \sim < 0.1$ ). 463

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### 465 Role of *I*<sub>CAN</sub> in the pre-BötC Respiratory Network

466 The hypothesis that  $I_{CAN}$  is involved in generation of the inspiratory rhythm is based on 467 experimental observations from in vitro mouse medullary slice preparations (Pena et al., 2004; 468 Thoby-Brisson & Ramirez, 2001), and in silico modeling studies (Jasinski et al., 2013; Rubin, 469 Hayes, et al., 2009; Toporikova & Butera, 2011). Theories of  $I_{CAN}$ -dependent bursting rely on 470 intracellular  $Ca^{2+}$  signaling mechanisms that have not been well defined.

471 Two models of  $I_{CAN}$ -dependent rhythmic bursting in vitro have been proposed and are 472 referred to as the "dual pacemaker" and "group pacemaker" models. In the dual pacemaker model, two types of pacemaker neurons are proposed that are either  $I_{NaP}$ -dependent (riluzole 473 sensitive) or  $I_{CAN}$ -dependent ( $Cd^{2+}$  sensitive) intrinsic bursters (see Rybak et al., 2014 for 474 475 review). In this model network, oscillations are thought to originate from these pacemaker 476 neurons which through excitatory synaptic interactions synchronize bursting and drive activity 477 of follower neurons within the pre-BötC. Although pacemaker neurons sensitive to neuronal  $Ca^{2+}$  flux blockade through  $Ca_V$  have been reported (Pena et al., 2004; Thoby-Brisson & 478 Ramirez, 2001), the source and mechanism driving intracellular  $Ca^{2+}$  oscillations has not been 479 described. Computational models of  $I_{CAN}$ -dependent pacemaker neurons rely on controversial 480 481 mechanisms for burst initiation and terminations, e.g. IP3-dependent  $Ca^{2+}$  oscillations 482 (Toporikova & Butera, 2011; Del Negro et al., 2010), that have been questioned from recent 483 negative experimental results (Beltran-Parrazal et al., 2012; Toporikova et al., 2015). 484 In versions of the "group pacemaker" model (Rubin et al., 2009; Feldman & Del Negro,

2006; Del Negro et al., 2010) network oscillations are initiated through recurrent synaptic 485 excitation that trigger postsynaptic  $Ca^{2+}$  influx. Subsequent  $I_{CAN}$  activation generates 486 membrane depolarization (inspiratory drive potential) to drive neuronal bursting. Synaptically 487 triggered  $Ca^{2+}$  influx and the contribution of  $I_{CAN}$  to the inspiratory drive potential of individual 488 489 pre-BötC neurons are experimentally supported (Mironov, 2008; Pace et al., 2007), however 490 the mechanism of burst termination remains unclear. Again, the computational group-491 pacemaker models that have been explored (Rubin, et al., 2009) rely on as yet unproven 492 mechanisms for burst termination, and in some cases lack key biophysical features of the pre-493 BötC neurons such as voltage-dependent frequency control and expression of  $I_{NaP}$ .

494 In our model, we showed that blockade of either  $I_{CAN}$  or synaptic interactions produce 495 qualitatively equivalent effects on network population activity amplitude and frequency when 496 the calcium transients are primarily generated from synaptic sources (Fig.4). Consequently, our 497 model predicts that blockade of  $I_{CAN}$  or synaptic interactions in the isolated pre-BötC in vitro 498 will produce comparable effects on amplitude and frequency. This is the case as Johnson et al. 499 (1994) showed that gradual blockade of synaptic interactions by low calcium solution 500 significantly decreases network activity amplitude while having little effect of frequency, 501 similar to the experiments where the  $I_{CAN}$  channel TRPM4 is blocked with 9-phenanthrol 502 (Koizumi et al., 2018).

503 Overall, our new model simulations for the isolated pre-BötC excitatory network 504 suggest that the role of  $I_{CAN}$ /TRPM4 activation is to amplify excitatory synaptic drive in 505 generating the amplitude of inspiratory population activity, independent of the biophysical 506 mechanism generating inspiratory rhythm. We note that the recent experiments have also 507 shown that in the more intact brainstem respiratory network that ordinarily generates patterns 508 of inspiratory and expiratory activity, endogenous activation of  $I_{CAN}$ /TRPM4 appears to 509 augment the amplitude of both inspiratory and expiratory population activity, and hence these 510 channels are fundamentally involved in inspiratory-expiratory pattern formation (Koizumi et 511 al., 2018).

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#### 513 Calcium Transients as Correlates of Activity

514 Neuronal calcium transients can arise from voltage-gated calcium sources, driven by 515 action potentials, and serve as correlates of neuronal activity. We analyzed the correlation 516 between calcium transients and inspiratory activity of individual inspiratory neurons as well as 517 the entire network, particularly since dynamic calcium imaging has been utilized to assess 518 activity of individual and populations of pre-BötC excitatory neurons in vitro during 519 pharmacological inhibition of  $I_{CAN}$ /TRPM4 (Koizumi et al., 2018). In our model, most of the 520 calcium influx is synaptically-triggered and may occur within a given neuron in the absence of 521 action potentials. We show that intracellular calcium transients are highly correlated with 522 network and cellular activity. This is true across individual neuron bursts and when comparing 523 changes in peak values of neuronal firing and intracellular calcium transients across the 524 duration of an  $I_{CAN}$  blockade simulation. The correlation at the onset of bursting in pacemaker 525 neurons are an exception. In these neurons, the correlations between the intracellular calcium 526 concentration and the instantaneous firing rate across a single burst are not apparent at the onset 527 of this burst. This is because pacemaker neurons start spiking before the rest of the network, 528 which precedes synaptically triggered calcium influx.

529 Additionally, we examined the relative change in the peak calcium transients in single 530 neurons as a function of  $I_{CAN}$  conductance. We show that in a subset of neurons the peak 531 calcium transient increases with reduced  $I_{CAN}$ . This result is surprising but is supported by the 532 recent calcium imaging data (Koizumi et al., 2018). This occurs in neurons that receive most 533 of their synaptic input from pacemaker neurons and our analyses suggest this is possible in 534 sparse networks, i.e. with relatively low connection probability. In pacemaker neurons,  $I_{CAN}$ 535 blockade leads to a reduction of their excitability resulting in an increased value of  $I_{NaP}$ 536 inactivation gating variable at the burst onset. Thus, during the burst, the peak action potential 537 frequency and the synaptic output from these neurons is increased with  $I_{CAN}$  blockade. 538 Consequently, neurons that receive synaptic input from pacemaker neurons will see an increase 539 in their peak calcium transients. In most neurons, however, synaptic input is received primarily 540 from follower neurons. Since  $I_{CAN}$  blockade de-recruits follower neurons, the synaptic input 541 and subsequent calcium influx in most decreases. Therefore, our model predicts that in a sparse 542 network, blocking  $I_{CAN}$  results in very diverse responses at the cellular level with overall 543 tendency to reduce intracellular calcium transients such that the amplitude of these transients 544 averaged over the entire population decreases during  $I_{CAN}$  blockade while their frequency is 545 unchanged. This is consistent with the experimental calcium imaging data (Koizumi et al., 546 2018).

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#### 548 Synaptic Calcium Sources

549 Our model suggests that calcium transients in the pre-BötC are coupled to excitatory 550 synaptic input, i.e. pre-synaptic glutamate release and binding to post-synaptic glutamate

receptors triggers calcium entry. The specific mechanisms behind this process is unclear, however it is likely dependent on specific types of ionotropic or metabotropic glutamate receptors.

554 There are three subtypes of ionotropic glutamate receptors, N-methyl-D-aspartate 555 (NMDA), Kainate (KAR), and  $\alpha$ -amino-3-hydroxy-5-methyl-4-isoxazolepropionic acid 556 (AMPA), all of which are expressed in the pre-BötC (Paarmann, Frermann, Keller, & 557 Hollmann, 2000) and have varying degrees of calcium permeability. NMDA and AMPA are 558 unlikely candidates for direct involvement in synaptically mediated calcium influx in the pre-559 BötC. Pharmacological blockade of NMDA receptors has no significant effect on the amplitude 560 or frequency of XII motor output (Morgado-Valle & Feldman, 2007; Pace et al., 2007) and 561 AMPA receptors in the pre-BötC show high expression of the subunit GluR2, which renders the AMPA ion channel pore impermeable to  $Ca^{2+}$  (Paarmann et al., 2000). It is possible, 562 563 however, that AMPA mediated depolarization may trigger calcium influx indirectly through 564 the voltage-gated calcium channel activation on the post-synaptic terminal. The contribution 565 of the latter to synaptically triggered calcium influx is likely small as pharmacological blockade 566 of L-, N- and P/Q-type calcium channels have no significant effect on XII motor output from 567 the pre-BötC (Morgado-Valle et al., 2008).

568 Calcium permeability through KAR receptors is dependent on subunit expression. The 569 KAR subunit GluK3 is highly expressed in the pre-BötC (Paarmann et al., 2000) and is calcium 570 permeable (Perrais et al., 2009) making it a possible candidate for synaptically mediated 571 calcium entry. Furthermore, GluK3 is insensitive to tonic glutamate release and only activated 572 by large glutamate transients (Perrais et al., 2009). Consequently, GluK3 may only be activated 573 when receiving synaptic input from a bursting presynaptic neuron which would presumably 574 generate large glutamate transients. The role of GluK3 in the pre-BötC has not been 575 investigated.

576 Metabotropic glutamate receptors (mGluR) indirectly activate ion channels through G-577 protein mediated signaling cascades. Group 1 mGluRs which include mGluR1 and mGluR5 578 are typically located on post-synaptic terminals (Shigemoto et al., 1997) and activation of group 579 1 mGluRs is commonly associated with calcium influx through calcium permeable channels 580 (Berg et al., 2007; Endoh, 2004; Mironov, 2008) and calcium release from intracellular calcium 581 stores (Pace et al., 2007).

In the pre-BötC, mGluR1/5 are thought to contribute to calcium influx by triggering the release of calcium from intracellular stores (Pace et al., 2007) and/or the activation of the transient receptor potential C3 (TRPC3) channel (Ben-Mabrouk & Tryba, 2010). Blockade of 585 mGluR1/5 reduces the inspiratory drive potential in pre-BötC neurons and reduces XII motor 586 output (Pace et al., 2007) which is consistent with the effects of  $I_{CAN}$ /TRPM4 blockade (Pace 587 et al., 2007). TRPC3 is a calcium permeable channel (Thebault et al., 2005) that is associated 588 with calcium signaling (Hartmann et al., 2011), store-operated calcium entry (Kwan et al., 589 2004), and synaptic transmission (Hartmann et al., 2011). TRPC3 is activated by diacylglycerol 590 (DAG) (Clapham, 2003) which is formed after synaptic activation of mGluR1/5. TRPC3 is 591 highly expressed in the pre-BötC and was hypothesized to underlie I<sub>CAN</sub> activation in the pre-592 BötC (Ben-Mabrouk & Tryba, 2010) and other brain regions (Amaral & Pozzo-Miller, 2007; 593 Zitt et al., 1997). Furthermore, TRPC3 and  $I_{CAN}$  have been shown to underlie slow excitatory 594 post synaptic current (sEPSC) (Hartmann et al., 2008; Hartmann et al., 2011). This is consistent 595 with our model since  $I_{CAN}$  activation is dependent on synaptically triggered calcium entry, and 596 the calcium dynamics are slower than the fast AMPA based current  $I_{Syn}$ . Therefore, in our 597 model,  $I_{CAN}$  decays relatively slowly and, hence, can be treated as a sEPSC. 598

In the pre-BötC, the effect of TRPC3 blockade by 3-pyrazole on network amplitude is 599 remarkably similar to blockade of TRPM4 (Koizumi et al., 2018). This suggests that the 600  $I_{CAN}$ /TRPM4 activation may be dependent on/coupled to TRPC3. A possible explanation is 601 that TRPC3 mediates synaptically-triggered calcium entry. It is also likely that TRPC3 plays a 602 role in maintaining background calcium concentration levels. We tested this hypothesis by 603 simulating the blockade of synaptically-triggered calcium influx while simultaneously 604 lowering the background calcium concentration (Fig. 12). These simulations generated large 605 reductions in activity amplitude with no effect on frequency which are consistent with data 606 from experiments where TRPC3 is blocked using 3-pyrazole (Koizumi et al., 2018). This 607 indirectly suggests that TRPC3 is critical for synaptically-triggered calcium entry and 608 subsequent  $I_{CAN}$  activation.



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**Figure 12** Comparison of experimental (black) and simulated (red) TRPC3 blockade (by  $Ca_{Syn}$  block) on network activity amplitude (**A**) and frequency (**B**). Simulated and experimental blockade begins at the vertical dashed line. The black line represents the mean and the gray band represents the mean S.E.M. of experimental JXII output recorded from neonatal rat brainstem slices in vitro, reproduced from Koizumi et al., 2018. Blockade was simulated by exponential decay of  $P_{Ca}$  with the following parameters: 3-Pyrazole:  $\gamma_{Block} = 1.0$ ,  $\tau_{Block} = 522.5s$ . The network parameters are:  $\bar{g}_{Ca} = 0.00175$  (*nS*),  $P_{Ca} = 0.0275$ ,  $P_{Syn} = 0.05$  and  $W_{max} = 0.096$  (*nS*).

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#### 618 I<sub>NaP</sub>-dependent Rhythmogenic Kernel

619  $I_{NaP}$  is a conductance present ubiquitously in pre-BötC inspiratory neurons, and is 620 established to underlie intrinsic oscillatory neuronal bursting in the absence of excitatory 621 synaptic interactions in neurons with sufficiently high  $I_{NaP}$  conductance densities (Koizumi & 622 Smith, 2008; Koizumi et al., 2008). Accordingly, we randomly incorporated this conductance 623 in our model excitatory neurons from a uniform statistical distribution to produce heterogeneity 624 in  $I_{NaP}$  conductance density across the population. Our simulations indicate that the circuit 625 neurons mostly with relatively high  $I_{NaP}$  conductance values underlie rhythm generation and 626 remain active after compete blockade of  $I_{CAN}$  in our model network, thus forming a  $I_{NaP}$ -627 dependent rhythmogenic kernel, including some neurons with intrinsic oscillatory bursting 628 behavior when synaptically uncoupled.

Recently it has become apparent that there is functional heterogeneity within pre-BötC excitatory circuits, including distinct subpopulations of neurons involved in generating periodic sighs (Toporikova et al., 2015; Li et al., 2016), and the subpopulations generating regular inspiratory activity. Activity of both of these subpopulations in the pre-BötC isolated 633 in vitro is proposed to be dependent on activation of  $I_{NaP}$  (Toporikova et al., 2015). Our 634 experimental and modeling results suggest that within the normal inspiratory population, there 635 are subpopulations distinguished by their role in rhythm versus amplitude generation due to 636 biophysical properties: there is a  $I_{CAN}$ /TRPM4-dependent recruitable population of excitatory 637 neurons for burst amplitude generation and the  $I_{NaP}$ -dependent rhythmogenic kernel 638 population. The spatial arrangements of these two synaptically connected excitatory 639 populations within the pre-BötC are currently unknown, and it remains an important 640 experimental problem to identify the cells constituting the rhythmogenic kernel and their 641 biophysical properties. This should now be possible, since our analysis and experimental 642 results suggest that the rhythmically active neurons of the kernel population can be revealed 643 and studied after pharmacologically inhibiting the  $I_{CAN}$ /TRPM4-dependent inspiratory burst-644 generating population.

645 Recently a "burstlet theory" for emergent network rhythms has been proposed to 646 account for inspiratory rhythm and pattern generation in the isolated pre-BötC in vitro (Kam et 647 al., 2013; Del Negro et al., 2018). This theory postulates that a subpopulation of excitatory 648 neurons generating small amplitude oscillations (burstlets) functions as the inspiratory rhythm 649 generator that drives neurons that generate the larger amplitude, synchronized inspiratory 650 population bursts. This concept emphasizes that subthreshold neuronal membrane oscillations 651 need to be considered and that there is a neuronal subpopulation that functions to independently 652 form the main inspiratory bursts. This is similar to our concept of distinct excitatory 653 subpopulations generating the rhythm versus the amplitude of inspiratory oscillations. 654 Biophysical mechanisms generating rhythmic burstlets and the large amplitude inspiratory 655 population bursts in the burstlet theory are unknown. We have identified a major  $Ca^{2+}$ -656 dependent conductance mechanism for inspiratory burst amplitude (pattern) generation and 657 show theoretically how this mechanism may be coupled to excitatory synaptic interactions and 658 is independent of the rhythm-generating mechanism. We also note that a basic property of  $I_{NaP}$ 659 is its ability to generate subthreshold oscillations and promote burst synchronization (Butera et 660 al., 1999b; Bacak et al., 2016). However, in contrast to our proposal for the mechanisms 661 operating in the kernel rhythm-generating subpopulation,  $I_{NaP}$  with its favorable voltage-662 dependent and kinetic autorhythmic properties- is not proposed to be a basic biophysical 663 mechanism for rhythm generation in the burstlet theory (Del Negro et al., 2018).

664 We emphasize that the above discussions regarding the role of  $I_{NaP}$  pertain to the 665 isolated pre-BötC including in more mature rodent experimental preparations in situ where

inspiratory rhythm generation has also been shown to be dependent on  $I_{NaP}$  (Smith et al., 2007). The analysis is more complex when the pre-BötC is embedded within interacting respiratory circuits in the intact nervous system generating the full complement of inspiratory and expiratory phase activity, where rhythmogenesis is tightly controlled by inhibitory circuit interactions and the contribution of  $I_{NaP}$  kinetic properties alone in setting the timing of inspiratory oscillations is diminished (Smith et al., 2007; Rubin et al., 2009; Richter & Smith, 2014).

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### 674 Conclusions

675 Based on our new data-driven computational model, distinct biophysical mechanisms are 676 involved in generating the rhythm and amplitude of inspiratory oscillations in the isolated pre-677 BötC excitatory circuits. Inspiratory rhythm generation arises from a group of  $I_{NaP}$ -dependent 678 excitatory neurons, including cells with intrinsic oscillatory bursting properties, that form a 679 rhythmogenic kernel. Rhythmic synaptic drive from these neurons triggers post-synaptic 680 calcium transients,  $I_{CAN}$  activation, and subsequent membrane depolarization which drives 681 bursting in the population of non-rhythmogenic follower neurons. We showed that activation 682 of  $I_{CAN}$  by synaptically-driven calcium influx functions as a mechanism that amplifies the 683 excitatory synaptic input to generate the inspiratory drive potential and population activity 684 amplitude in these non-rhythmogenic neurons. Consequently, blockade of  $I_{CAN}$  causes a robust 685 decrease in overall network activity amplitude via de-recruitment of these follower neurons 686 without perturbations of the inspiratory rhythm, which is consistent with the results with 687 experimental blockade of I<sub>CAN</sub>/TRPM4 channels. Our model provides a theoretical explanation 688 for these recent paradigm-shifting experimental results that  $I_{CAN}$  is not fundamentally involved 689 in generating the inspiratory rhythm and gives new insights into the functional operation of 690 pre-BötC excitatory circuits.

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#### 693 Materials and Methods

#### 694 Model Description

695 The model describes a network of N = 100 synaptically coupled excitatory neurons. Simulated

696 neurons are comprised of a single compartment described using a Hodgkin Huxley formalism.

- 697 For each neuron, the membrane potential  $V_m$  is given by the following current balance equation:
- 698

699 
$$C_m \frac{dV_m}{dt} + I_{Na} + I_K + I_{Leak} + I_{NaP} + I_{CAN} + I_{Ca} + I_{Syn} = 0$$

700

where  $C_m$  is the membrane capacitance,  $I_{Na}$ ,  $I_K$ ,  $I_{Leak}$ ,  $I_{NaP}$ ,  $I_{CAN}$ ,  $I_{Ca}$  and  $I_{Syn}$  are ionic currents through sodium, potassium, leak, persistent sodium, calcium activated non-selective cation, voltage-gated calcium, and synaptic channels, respectively. Description of these currents, synaptic interactions, and parameter values are taken from (Jasinski et al., 2013). The channel currents are defined as follows:

706

707 
$$I_{Na} = \bar{g}_{Na} \cdot m_{Na}^3 \cdot h_{Na} \cdot (V_m - E_{Na})$$

708

 $I_K = \bar{g}_K \cdot m_K^4 \cdot (V_m - E_K)$ 

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712

/11	$I_{Leak} = g_{Leak} \cdot (V_m)$	$L = L_{Leak}$

- 713  $I_{NaP} = \bar{g}_{NaP} \cdot m_{NaP} \cdot h_{NaP} \cdot (V_m E_{Na})$
- 714

715 
$$I_{CAN} = \bar{g}_{CAN} \cdot m_{CAN} \cdot (V_m - E_{CAN})$$

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717 
$$I_{Ca} = \bar{g}_{Ca} \cdot m_{Ca} \cdot h_{Ca} \cdot (V_m - E_{Ca})$$

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 $I_{Syn} = g_{Syn} \cdot \left( V_m - E_{Syn} \right)$ 

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For I_{Na}, I_K, I_{NaP}, and I_{Ca}, the dynamics of voltage-dependent gating variables m_i, and h_i are
defined by the following differential equation:
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727 
$$\tau_{\eta}(V) \cdot \frac{d\eta}{dt} = \eta_{\infty}(V) - \eta; \quad \eta \in \{m_i, h_i\}$$

728

729 where steady state activation/inactivation  $\eta_{\infty}$  and time constant  $\tau_{\eta}$  are given by:

where  $\bar{g}_i$  is the maximum conductance,  $E_i$  is the reversal potential,  $m_i$  and  $h_i$  are voltage dependent gating variables for channel activation and inactivation, respectively, and  $i \in$  $\{Na, K, Leak, NaP, CAN, Ca, Syn\}$ . The parameters  $\bar{g}_i$  and  $E_i$  are given in Table 1.

731 
$$\eta_{\infty}(V) = \left(1 + e^{-(V - V_{\eta_{1/2}})/k_{\eta}}\right)^{-1}$$

732

733 
$$\tau_{\eta}(V) = \tau_{\eta_{max}} / \cosh \left( (V - V_{\tau \eta_{1/2}}) / k_{\tau_{\eta}} \right).$$

734

For the voltage-gated potassium channel, steady state activation 
$$m_{K\infty}(V)$$
 and time constant

736  $\tau_{mK}(V)$  are given by:

737 
$$m_{K\infty}(V) = \frac{\alpha_{\infty}(V)}{\alpha_{\infty}(V) + \beta_{\infty}(V)}$$

- 738
- 739  $\tau_{mK}(V) = 1/(\alpha_{\infty}(V) + \beta_{\infty}(V))$
- 740

741 where

- 742  $\alpha_{\infty}(V) = A_{\alpha} \cdot (V + B_{\alpha})/(1 exp(-(V + B_{\alpha})/\kappa_{\alpha}))$ 743
- 744  $\beta_{\infty}(V) = A_{\beta} \cdot exp(-(V+B_{\beta})/\kappa_{\beta}).$
- 745

The parameters  $V_{\eta_{1/2}}$ ,  $V_{\tau\eta_{1/2}}$ ,  $\kappa_{\eta}$ ,  $\kappa_{\tau\eta}$ ,  $\tau_{\eta_{max}}$ ,  $A_{\alpha}$ ,  $A_{\beta}$ ,  $B_{\alpha}$ ,  $B_{\beta}$ ,  $\kappa_{\alpha}$ , and  $\kappa_{\beta}$  are given in Table 1. *I<sub>CAN</sub>* activation is dependent on the intracellular calcium concentration  $[Ca]_{in}$  and is given by: 748

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749 
$$m_{CAN} = 1/(1 + (Ca_{1/2}/[Ca]_{in})^n).$$

750

The parameters  $Ca_{1/2}$  and *n*, given in Table 1, represent the half-activation calcium concentration and the Hill Coefficient, respectively.

Calcium enters the neurons through voltage-gated calcium channels  $(Ca_V)$  and/or synaptic channels  $(Ca_{Syn})$ , where a percentage  $(P_{Ca})$  of the synaptic current  $(I_{Syn})$  is assumed to consist of Ca<sup>2+</sup> ions. A calcium pump removes excess calcium with a time constant  $\tau_{Ca}$  and sets the minimum calcium concentration  $Ca_{min}$ . The dynamics of  $[Ca]_{in}$  is given by the following differential equation:

758

759 
$$\frac{d[Ca]_{in}}{dt} = -\alpha_{Ca}(I_{Ca} + P_{Ca} \cdot I_{syn}) - ([Ca]_{in} - Ca_{min})/\tau_{Ca}.$$

760

761 The parameters  $\alpha_{Ca}$  is a conversion factor relating current and rate of change in  $[Ca]_{in}$ , see

762 Table 1 for parameter values.

763 The synaptic conductance of the  $i^{th}$  neuron  $(g_{Syn}^i)$  in the population is described by the

- 764 following equation:
- 765

766 
$$g_{Syn}^{i} = g_{Tonic} + \sum_{j,n} w_{ji} \cdot C_{ji} \cdot H(t - t_{j,n}) \cdot e^{-(t - t_{j,n})/\tau_{Syn}}$$

767

where  $w_{ji}$  is the weight of the synaptic connection from cell *j* to cell *i*, *C* is a connectivity matrix ( $C_{ji} = 1$  if neuron *j* makes a synapse on neuron *i*, and  $C_{ji} = 0$  otherwise), *H*(.) is the Heaviside step function, *t* is time,  $\tau_{Syn}$  is the exponential decay constant and  $t_{j,n}$  is the time at which an action potential *n* is generated in neuron *j* and reaches neuron *i*.

772

773**Table 1. Model parameter values**. The channel kinetics, intracellular  $Ca^{2+}$  dynamics and the774corresponding parameter values, were derived from previous models (see (Jasinski et al., 2013)775and the references therein).

Channel	Parameters
I <sub>Na</sub>	$\bar{g}_{Na} = 150.0 \ nS, E_{Na} = 55.0 \ mV,$
	$V_{m_{1/2}} = -43.8 \text{ mV}, k_m = 6.0 \text{ mV},$
	$V_{\tau m_{1/2}} = -43.8 \ mV, \ k_{\tau_m} = 14.0 \ mV, \ \tau_{m_{max}} = 0.25 \ ms,$
	$V_{h_{1/2}} = -67.5 \text{ mV}, k_h = -10.8 \text{ mV},$
	$V_{\tau h_{1/2}} = -67.5 \ mV, k_{\tau_h} = 12.8 \ mV, \tau_{h_{max}} = 8.46 \ ms$
$I_K$	$\bar{g}_K = 160.0 \ nS, E_K = -94.0 \ mV,$
	$A_{\alpha} = 0.01, B_{\alpha} = 44.0 \ mV, \kappa_{\alpha} = 5.0 \ mV$
	$A_{\beta} = 0.17, B_{\beta} = 49.0 \ mV, \kappa_{\beta} = 40.0 \ mV$
I <sub>Leak</sub>	$\bar{g}_{Leak} = 2.5 \ nS$ , $E_{Leak} = -68.0 \ mV$ ,
I <sub>NaP</sub>	$\bar{g}_{NaP} \in [0.0, 5.0]  nS,$
	$V_{m_{1/2}} = -47.1 \text{ mV}, k_m = 3.1 \text{ mV},$
	$V_{\tau m_{1/2}} = -47.1 \ mV, \ k_{\tau_m} = 6.2 \ mV, \ \tau_{m_{max}} = 1.0 \ ms,$
	$V_{h_{1/2}} = -60.0 \text{ mV}, k_h = -9.0 \text{ mV},$
	$V_{\tau h_{1/2}} = -60.0 \ mV, \ k_{\tau_h} = 9.0 \ mV, \ \tau_{h_{max}} = 5000 \ ms$

I <sub>CAN</sub>	$\bar{g}_{CAN} = 1.0 \ nS, \ E_{CAN} = 0.0 \ mV,$
	$Ca_{1/2} = 0.00074 \ mM, n = 0.97$
I <sub>Ca</sub>	$\bar{g}_{Ca} = 0.01  nS, E_{Ca} = R \cdot T/F \cdot \ln([Ca]_{out}/[Ca]_{in}),$
	$R = 8.314 J/(mol \cdot K), T = 308.0 K,$
	$F = 96.485 \ kC/mol, \ [Ca]_{out} = 4.0 \ mM$
	$V_{m_{1/2}} = -27.5 \text{ mV}, k_m = 5.7 \text{ mV}, \tau_m = 0.5 \text{ ms},$
	$V_{h_{1/2}} = -52.4 \text{ mV}, k_h = -5.2 \text{ mV}, \tau_h = 18.0 \text{ ms}$
Ca <sub>in</sub>	$\alpha_{Ca} = 2.5 \cdot 10^{-5} mM/fC, P_{Ca} = 0.01, Ca_{min} = 1.0 \cdot 10^{-10} mM, \tau_{Ca} = 50.0 mS$
I <sub>Syn</sub>	$g_{Tonic} = 0.31  nS, E_{Syn} = -10.0  mV, \tau_{Syn} = 5.0  ms$

777

778 To account for heterogeneity of neuron properties within the network, the persistent sodium 779 current conductance,  $\bar{g}_{NaP}$ , for each neuron was assigned randomly based on a uniform 780 distribution over the range [0.0, 5.0] nS which is consistent with experimental measurements 781 (Rybak et al., 2003; Koizumi & Smith, 2008; Koizumi et al., 2008). The weight of each 782 synaptic connection was uniformly distributed over the range  $w_{ji} \in [0, W_{max}]$  where  $W_{max}$ 783 ranged from 0.0 to 1.0 nS depending on the network connectivity and specific simulation. The 784 elements of the network connectivity matrix,  $C_{ii}$ , are randomly assigned values of 0 or 1 such 785 that the probability of any connection between neuron i and neuron i being 1 is equal to the 786 network connection probability  $P_{Svn}$ . We varied the connection probability over the range  $P_{Syn} \in [0.05, 1.0]$ , however, a value of  $P_{Syn} = 0.05$  was used in most simulations. 787

788 789

#### Data Analysis and Definitions

790 The time of an action potential was defined as when the membrane potential of a neuron 791 crosses -35mV in a positive direction. The network activity amplitude and frequency were 792 determined by identifying peaks and calculating the inverse of the interpeak interval in 793 histograms of network spiking. Network histograms of the population activity were calculated 794 as the number of action potentials generated by all neurons per 50 ms bin per neuron with units 795 of *spikes/s*. The number of recruited neurons is defined as the peak number of neurons that 796 spiked at least once per bin during a network burst. The average spike frequency of recruited 797 neurons is defined as the number of action potentials per bin per recruited neuron with units of 798 spikes/s. The average network resting membrane potential was defined as the average 799 minimum value of  $V_m$  in a 500 ms window following a network burst. The average inactivation 800 of the persistent sodium current at the start of each burst was defined by the maximum of the

801 average value of  $h_{NaP}$  in a 500 ms window before the peak of each network burst. The average 802 inactivation of the persistent sodium current at the end of each burst was defined by the 803 maximum of the average value of  $h_{NaP}$  in a 500 ms window after the peak of each network 804 burst. Synaptic strength is defined as the number of neurons in the network multiplied by the 805 connection probability multiplied by the average weight of synaptic connections  $(N \cdot P_{Svn} \cdot$ 806  $\frac{1}{2}W_{max}$ ). Pacemaker neurons were defined as neurons that continue bursting intrinsically after 807 complete synaptic blockade. Follower neurons were defined as neurons that become silent after 808 complete synaptic blockade. The inspiratory drive potential is defined as the envelope of 809 depolarization that occurs in neurons during the inspiratory phase of the network oscillations 810 (Morgado-Valle et al. 2008).

811

## 812 Characterization $I_{CAN}$ in regulating network activity amplitude and frequency in $Ca_V$ 813 and $Ca_{Syn}$ Models

814 To characterize the role of  $I_{CAN}$  in regulation of network activity amplitude and 815 frequency we slowly increased the conductance  $(\bar{g}_{CAN})$  in our simulations from zero until the 816 network transitioned from a rhythmic bursting to a tonic (non-bursting) firing regime. To 817 ensure that the effect(s) are robust, these simulations were repeated over a wide range of 818 synaptic weights, synaptic connection probabilities, and strengths of the intracellular calcium 819 transients from  $Ca_V$  or  $Ca_{Syn}$  sources. Changes in network activity amplitude were further 820 examined by plotting the number of recruited neurons and the average action potential 821 frequency of recruited neurons versus  $\bar{g}_{CAN}$ .

822

#### 823 Simulated Pharmacological Manipulations

In simulations that are compared with experimental data, both  $Ca_V$  and  $Ca_{Syn}$  calcium sources are included. Pharmacological blockade of  $I_{CAN}$  was simulated by varying the conductance,  $\bar{g}_{CAN}$  according to a decaying exponential function

827 
$$\bar{g}_{CAN}(t) = g_{CAN}^{max} - \gamma_{block} \cdot \left(1 - e^{-t/\tau_{block}}\right).$$

The percent block  $\gamma_{block}$ , decay constant  $\tau_{block}$  and the maximum  $I_{CAN}$  conductance  $g_{CAN}^{max}$ were adjusted to match the experimental changes in network amplitude. The synaptic weight of the network was chosen such that at  $\bar{g}_{CAN} = 0$  the network activity amplitude was close to 20% of maximum. To reduce the computational time, the duration of  $I_{CAN}$  block simulations was one tenth of the total of experimental durations. For comparison, the plots of normalized change in amplitude and frequency of the simulations were stretched over the same time-period as experimental data. Increasing the simulation time had no effect on our results (data not

835 shown).

836

### 837 Comparison with Calcium Imaging Data

838 To allow comparisons with network and cellular calcium imaging data, we analyzed rhythmic

calcium transients from our simulations. Single cell calcium signals are represented by  $[Ca]_i$ .

840 The network calcium signal was calculated as the average intracellular calcium concentration

841 in the network  $(\sum_{i=1}^{N} [Ca]_i/N)$ .

842

#### 843 Integration Methods

All simulations were performed locally on an 8-core Linux-based operating system or on the high-performance computing cluster Biowulf at the National Institutes of Health. Simulation software was custom written in C++. Numerical integration was performed using the exponential Euler method with a fixed step-size ( $\Delta t$ ) of 0.025*ms*. In all simulations, the first 50 s of simulation time was discarded to allow for the decay of any initial condition-dependent transients.

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#### 852

## 853 5. Acknowledgments

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