

Within-cycle instantaneous frequency profiles report oscillatory waveform dynamics.

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Abstract: Non-sinusoidal waveform is emerging as an important feature of neuronal oscillations. However, the role of single cycle shape dynamics in rapidly unfolding brain activity remains unclear. Here, we develop an analytical framework that isolates oscillatory signals from time-series using masked Empirical Mode Decomposition to quantify dynamical changes in the shape of individual cycles (along with amplitude, frequency and phase) using instantaneous frequency. We show how phase-alignment, a process of projecting cycles into a regularly sampled phase-grid space, makes it possible to compare cycles of different durations and shapes. ‘Normalised shapes’ can then be constructed with high temporal detail whilst accounting for differences in both duration and amplitude. We find that the instantaneous frequency tracks non-sinusoidal shapes in both simulated and real data. Notably, in local field potential recordings of mouse hippocampal CA1, we find that theta oscillations have a stereotyped slow-descending slope in the cycle-wise average, yet exhibiting high variability on a cycle-by-cycle basis. We show how Principal Components Analysis allows identification of motifs of theta cycle waveform that have distinct associations to cycle amplitude, cycle duration and animal movement speed. By allowing investigation into oscillation shape at high temporal resolution, this analytical framework will open new lines of enquiry into how neuronal oscillations support moment-by-moment information processing and integration in brain networks.

1: Introduction

Frequency, phase and amplitude have long been reported as important features of neuronal oscillations with behavioural and electrophysiological relevance. Furthermore, neuronal oscillations show non-sinusoidal waveform shapes that span a wide range of spatial and temporal scales (Cole and Voytek, 2017). Whilst waveform shape is emerging as a fourth relevant feature of neuronal oscillations, many theories of neuronal oscillations currently assume sinusoidal waveforms. This might be due to the fact that characterising and quantifying non-sinusoidal waveforms remains a substantial analytic challenge (Amzica and Steriade, 1998; Cole and Voytek, 2017). To uncover the role of waveform dynamics in rapidly unfolding brain activity, there is a growing need for novel analysis methods that are able to characterise a wide range of waveform shape features at the single-cycle level.

Waveform shape related parameters, such as skewness or asymmetry, can be estimated from higher order Fourier spectra such as the bispectrum or bicoherence (Bartz et al., 2019; Elgar, 1987; Sheremet et al., 2016). These methods require relatively long data segments to have high frequency resolution, and therefore do not provide single-cycle estimates. Alternatively, a set of waveform features for individual cycles can be described by the relative durations of different quartiles of a cycle (Belluscio et al., 2012; Cole and Voytek, 2019; Trimper et al., 2014). Whilst this approach is tractable on single cycles, the extracted features must be defined *a priori* and are limited to the resolution of the selected cycle control points such as the extrema and zero-crossings.

The temporal dynamics in oscillatory frequency can be quantified for a given waveform by its instantaneous frequency computed from the differential of the signal's instantaneous phase (Boashash, 1992; Huang et al., 2009). Such instantaneous frequency estimates have been used previously in electrophysiology to explore dynamics in oscillatory peak frequency at high temporal resolution (Cohen, 2014; Liang et al., 2005; Nelli et al., 2017; Rudrauf et al., 2006). Crucially, any non-sinusoidal waveform features in an oscillation will lead to within-cycle instantaneous frequency modulations in which the frequency of an oscillation changes from moment-to-moment within a single cycle (Huang et al., 1998). The degree of non-linearity of an oscillation is related to the total amount of within-cycle frequency modulation (Huang et al., 2014; Tsai et al., 2016; Wang et al., 2012; Yeh et al., 2020).

We introduce a novel approach which creates smooth waveform shape profiles that describe non-sinusoidal features in single cycles with high temporal detail. To this end, we first operationalise waveform shape as the profile of instantaneous frequency across the cycle's instantaneous phase. We then identify when and how an ongoing cycle deviates from a sinusoidal waveform by identifying points in the cycle where instantaneous frequency departs from a flat profile. For example, a cycle with a wide peak has a relatively low instantaneous frequency around the peak, and a cycle with a fast-ascending edge will have a relatively high instantaneous frequency between the trough and peak. In order to allow between cycle comparisons, we also need to account for how different cycles of an oscillation will play out

44 at different speeds, leading to differences in extrema timing and overall duration. To
45 overcome these problems, we introduce the process of ‘phase-alignment’ which reregisters
46 the instantaneous frequency profiles onto a regularly sampled set of points in phase. To
47 obtain the instantaneous phase time course of each cycle, we use the Empirical Mode
48 Decomposition (EMD). EMD decomposes the time-series of interest into its oscillatory
49 modes (Intrinsic Mode Functions; IMFs) that retain the non-stationary and non-linear signal
50 features.

51

52 We outline and validate our novel approach in simulated data before applying it to theta-band
53 oscillations recorded in the local field potentials (LFPs) of the mouse hippocampal CA1
54 during active exploratory behaviour. The hippocampal theta rhythm has a characteristic non-
55 sinusoidal waveform shape (Belluscio et al., 2012; Buzsáki et al., 1985; Siapas et al., 2005)
56 that is modulated by movement (Sheremet et al., 2016) and changes in sleep or drug states
57 (Buzsáki et al., 1985). Using EMD to identify the theta rhythm, we show that phase-aligned
58 instantaneous frequency is able to robustly characterise a continuous waveform shape profile
59 for hippocampal theta. The results confirm the stereotyped fast-ascending and slow-
60 descending shape in the cycle-wise average. Additionally, it reveals but with high amounts of
61 shape variation across single cycles of theta, which we describe using a set of data-driven
62 shape ‘motifs’. Finally, we find that cycle-level shape motifs have differential associations
63 with theta amplitude, theta cycle duration and mouse movement speed. Overall, we
64 demonstrate that behaviourally relevant dynamics in single-cycle oscillatory waveforms can
65 be accurately and intuitively explored with phase-aligned instantaneous frequency profiles.

66 **2: Methods**

67

68 **2.1: Data and Code Availability Statement**

69

70 Code for the analyses in this paper are freely available online ([https://github.com/OHBA-](https://github.com/OHBA-analysis/Quinn2021_Waveform)
71 [analysis/Quinn2021_Waveform](https://github.com/OHBA-analysis/Quinn2021_Waveform)) and data are available from the MRC BNDU Data Sharing
72 Platform (<https://data.mrc.ox.ac.uk/data-set/instantaneous-frequency-profiles-theta-cycles>;
73 requires free registration). The analyses in this study were carried out in python 3.7 using
74 v0.4.0 of the EMD package (Quinn et al., 2021; <https://emd.readthedocs.io/>) and glmtools
75 v0.1.0 for General Linear Model design and fitting (<https://pypi.org/project/glmtools/>). The
76 wavelet transforms and principal components analysis were computed using SAILS v1.1.1
77 (Quinn and Hymers, 2020). The underlying python dependencies were numpy (Harris et al.,
78 2020) and scipy (SciPy 1.0 Contributors et al., 2020) for computation and matplotlib (Hunter,
79 2007) for visualisation.

80

81 **2.2: Masked Empirical Mode Methods**

82

83 The EMD is implemented as a sifting algorithm that incrementally extracts the highest
84 frequency features of a time-series into its oscillatory components known as IMFs (Huang et
85 al., 1998). Once identified, the IMF is subtracted from the signal and the sifting process
86 repeated to find the next fastest set of oscillatory dynamics. This process is iterated through
87 until only a residual trend remains in the dataset, constituting the very slowest dynamics of
88 the signal.

89

90 Transient or intermittent oscillatory signals can lead to a mix of different frequency
91 components appearing in a single IMF; an issue known as mode mixing (Deering and Kaiser,
92 2005; Wu and Huang, 2009). To reduce mode mixing, we use an adapted version of the mask
93 sift (Deering and Kaiser, 2005; Tsai et al., 2016). The mEMD involves the same core process
94 as the original sift outlined above. However, at each iteration, a masking signal is added to
95 the data before all the extrema (maxima and minima) in the masked signal are identified. An
96 upper and lower amplitude envelope is then estimated by interpolating between the maxima
97 and minima respectively. The average of the upper and lower envelope is subtracted from the
98 data and the extrema identification, envelope interpolation and subtraction repeated until the
99 average of the upper and lower envelopes are close to zero. The mask is then subtracted from
100 the signal to return the IMF.

101

102 The performance of the sift is limited by a number of factors such as accuracy in peak
103 detection and overshoot or edge effects in the envelope interpolation. As such, when applied
104 to real data, the sift may not always perfectly isolate individual oscillations. To prevent
105 envelope overshoot we used a monotonic PCHIP (Piecewise Cubic Hermite Interpolating
106 Polynomial; implemented in `scipy.interpolation.PChipInterpolator`) interpolation during the
107 sifting. The monotonic PCHIP interpolation avoids overshoot of the amplitude envelopes
108 where the data are not completely smooth. When compared to a cubic-spline interpolation,

109 this reduces instances where the envelopes become very large, or the upper and lower
110 envelopes cross over.

111

112 2.2.1: Frequency Transformation

113

114 The analytic form of each IMF was constructed using the Hilbert transform and the
115 instantaneous phase set as the angle of the analytic form on the complex plane (See figure 2C
116 for an example). To attenuate noise in the phase estimation, the unwrapped phase time-course
117 was smoothed using a Savitsky-Golay filter (`scipy.signal.savgol_filter`; order = 1, window
118 size = 3 samples). The instantaneous frequency (See figure 2D for an example) in Hertz is
119 then computed from the derivative of this unwrapped phase:

120

$$121 \quad IF = \frac{fs}{2\pi} \frac{d\phi(t)}{dt}$$

122

123 Where fs is the sampling frequency and $\phi(t)$ is the unwrapped instantaneous phase time-
124 course. The instantaneous amplitude time-course is computed as the absolute value of the
125 analytic form of each IMF.

126

127 The distribution of instantaneous amplitude values by frequency or time and frequency can
128 be computed from these instantaneous frequency metrics. A sparse matrix $H \in \mathbb{R}^{(T \times F)}$ is
129 filled with the instantaneous amplitudes value from the IMFs at their respective time and
130 frequency co-ordinates. This matrix is the Hilbert-Huang Transform (HHT; See figure 2E for
131 an example) and provides an alternative time-frequency transform to traditional Fourier based
132 methods such as the short-time Fourier transform and the wavelet transform (Huang et al.,
133 1998).

134

135 2.2.2: Cycle Detection

136

137 The next stage is to segment the IMFs into their constituent cycles and identify which cycles
138 will be included in further analysis. The start and end of theta cycles are located by the
139 differentials greater than six in the phase. The start and end point of cycles in this paper is the
140 ascending zero-crossing as this occurs at the point where the phase time-course wraps. Once
141 identified, some cycles will be ‘bad’ in the sense that the oscillation captured by the IMF is
142 not well represented, e.g., because the rhythm is not present over that time period or it is
143 poorly estimated, and will be excluded from subsequent analyses. This is important for
144 instantaneous frequency analyses as the differentiation step (see section 2.2.1) can be very
145 noise sensitive. Included cycles are identified from the wrapped instantaneous phase time-
146 course of the IMF containing the oscillation to be analysed. As the instantaneous phase
147 computation via the Hilbert transform returns a value for every sample regardless of whether
148 a prominent rhythm is present, only ‘good’ theta cycles are retained for further analysis. A
149 good cycle is defined as having a phase with a strictly positive differential (i.e., no phase

150 reversals) that starts with a value $0 \leq x \leq \pi/24$ and end within $2\pi - \pi/24 \leq x \leq 2\pi$ and 4 control
151 points (peak, trough, ascending edge and descending edge).

152

153 2.2.3: Control Point analysis

154

155 One approach to quantify oscillatory waveforms is to base analyses around parts of the cycle
156 which can be straightforwardly identified in all cycles (Belluscio et al., 2012; Cole and
157 Voytek, 2019). For example, the peak, trough, ascending zero-crossing and descending zero-
158 crossings of a cycle can define a set of control points describing the relative timing of
159 inflections and mid-points in a cycle. In the present analyses, the extrema (peaks and troughs)
160 are detected by finding zero-crossings in the differential of an IMF. This initial extrema
161 estimate is limited by the sampling frequency of the dataset and is refined using parabolic
162 interpolation (Rato et al., 2008). The zero-crossings are initially identified from sign changes
163 in the IMF time-course and refined by linear interpolation. The ratio of temporal durations
164 between these control points can describe large scale shape features. Finally, we compute the
165 peak-to-trough ratio and the ascent-to-descent ratios for each cycle (Cole and Voytek, 2019).

166

167 2.2.4: Phase-alignment

168

169 We present an alternative approach to control points that ensures that entire waveform
170 profiles can be combined and/or compared across cycles despite cycle-by-cycle differences in
171 progression rate and overall duration. To compare waveforms across cycles that play out at
172 different speeds, we use phase-alignment to register cycles onto a common grid. Phase
173 alignment is performed on the instantaneous phase of a cycle and a measure of interest, such
174 as the instantaneous frequency, which is observed at the same time-intervals. A linear one-
175 dimensional interpolation function is fitted between the instantaneous phase (as x values) and
176 the instantaneous frequency (as y values). The interpolation function is evaluated on a
177 template set of instantaneous phase values with a linear spacing between 0 and 2π , if any
178 points in the template fall outside the fitted range, the interpolator returns an extrapolated
179 value. This interpolated version of instantaneous frequency is then directly comparable across
180 cycles as each point in the phase will occur at the same time. We compute phase-alignment
181 using a linear interpolation across 48 fixed points across the zero to 2π phase range.

182

183 Once an instantaneous frequency profile has been phase aligned, we can visualise a
184 normalised waveform by projecting the frequency content back to a phase-time course. This
185 is achieved by re-normalising the instantaneous frequency from Hertz back to radians in
186 order to create a profile of successive phase differences. The phase time-course is then
187 reconstructed from the cumulative summation of these phase differences. An oscillatory
188 waveform can then be computed by taking the sine transform of this phase time-course. The
189 resulting waveform has an amplitude of one and a consistent time-axis for all cycles. This
190 ‘normalised waveform’ allows for visualisation of shape between cycles with different
191 durations and amplitudes.

192 2.2.5: Describing shape an instantaneous frequency mean vector

193

194 A simplified summary of a cycle's shape can be computed from a mean-vector of the phase
195 aligned instantaneous frequency according to the following equation:

196

197

$$MV = \overline{IF} e^{i\overline{\phi}}$$

198

199 Where $\overline{\phi}$ is the uniform phase-grid used in phase-alignment and \overline{IF} is the phase-aligned
200 instantaneous frequency. This is similar to the mean-vector approach to computing phase-
201 amplitude-coupling (Canolty et al., 2006). The mean-vector of a sinusoidal cycle will be zero
202 whilst a non-sinusoidal cycle will return a complex value whose angle indicates which phase
203 has the highest instantaneous frequency and the magnitude indicates the extent of the
204 frequency modulation through the cycle. This method provides a straightforward summary
205 but is only sensitive to unimodal deviations from a flat instantaneous frequency profile.

206

207 2.2.6: Principal Components shape motifs.

208

209 A more complete, data-driven approach to summarising shape from instantaneous frequency
210 uses Principal Components Analysis (PCA) to identify the principal modes of variation in
211 shape across the included cycles. Phase-aligned cycles are concatenated into a single matrix
212 of size [nphases x ncycles]. The second dimension of this matrix is reduced to by PCA. This
213 results in an [nphases x ncomponents] matrix of shape 'motifs' defined by the distribution of
214 component weights across phase and an [ncomponents x ncycles] matrix of PC scores
215 indicating the presence of each component motif in each individual cycle.

216

217 The component motif matrix defines the axes of variability in waveform shape across cycles.
218 The shapes captured along each of these axes are visualised by defining a set of PC scores
219 containing the maximum or minimum observed score for the PC in question and zeros for all
220 others. These scores can be projected back into the original data space to provide exemplar
221 instantaneous frequency profiles for both extremes of the PC axes. Finally, these exemplar IF
222 profiles can be projected back into the time-domain to generate a normalised waveform that
223 preserves the shape depicted in the exemplar IF profiles.

224

225 **2.3: Simulation analyses**

226

227 2.3.1: Schematic cycle generation

228

229 To illustrate the relationship between waveform shape, instantaneous phase and frequency, a
230 set of noise free oscillations were generated. First, a linearly progressing phase time-course is
231 generated and sinusoid is created by taking a sine-transform of this wrapped phase. Different
232 non-sinusoidal cycles are generated by modulating the unwrapped phase time-course by sine
233 and cosine waves at different phases and frequencies. The cycles with extrema and edge
234 asymmetry are generated by modulating the phase with a 1 Hz sine or cosine respectively.

235 The extrema curvature examples are generated by modulating the linear phase with a 2 Hz
236 sinusoid. From the computed cycle time-courses the instantaneous phase and instantaneous
237 frequency are re-estimated using the Hilbert transform. Finally, the waveform shape is
238 represented by phase-aligning the instantaneous frequency time-course of each cycle type
239 with its instantaneous phase.

240

241 2.3.2: Noisy Signal Generation

242

243 A more realistic noisy simulation was used for the results in figures 3, 4 and 5. A simulated
244 oscillation at 12 Hz was generated using an autoregressive oscillator with the following
245 transfer function:

246

$$247 \quad H(f) = \frac{1}{1 - 2r \cos \theta + r^2}$$

248

249 Where θ is the angular frequency of the oscillator (in rads/sec) and r is the magnitude of the
250 roots of the polynomial ($0 < r < 1$). For this simulation, we computed H for $r=0.95$ and θ
251 equivalent to 12 Hz and used its parameters to filter (a forwards and backwards filter using
252 `scipy.signal.filtfilt`) white noise. This generates a noisy sinusoidal oscillation which contains
253 random dynamics in the frequency and amplitude of each oscillatory cycle. Sixty seconds of
254 data were generated at 512 Hz.

255

256 This simulated oscillation was then modulated by one of two equations defined in equations
257 50.24 and 50.25 in section 50-6 in Volume 1 of Feynman's Lectures of Physics (Feynman et
258 al., 2011). The first equation defines a linear system that scales the signal by a constant
259 leaving the waveform shape unchanged.

260

$$261 \quad x_{out} = Kx_{in} + e(t)$$

262

263 Whilst the second equation defines a nonlinear system that includes a term inducing a change
264 in waveform shape as well as scaling the signal.

265

$$266 \quad x_{out} = K[x_{in} + \epsilon x_{in}^2] + e(t)$$

267

268 The nonlinearity in the second equation makes the peak of the oscillations shorter and widens
269 the trough. Both systems include an additive white noise term.

270

271 2.3.3: Noisy Signal Analysis

272

273 The simulations are separated into IMFs using the masked sift. The relatively straightforward
274 dynamics in this simulation allow a simplified mask-sift to be applied. To define the mask
275 frequencies, a first IMF is extracted using the standard sift routine. The number of zero-
276 crossings in this IMF defines the frequency of the initial mask with subsequent masks being

277 applied at half the frequency of the previous one. The mask amplitudes are set equivalent to
278 1-standard deviation of the previously extracted IMF.

279

280 The 12-Hz oscillation is isolated in the third IMF (IMF-3) of this mask sift. The frequency
281 transformation of this IMF is computed using the Hilbert transform (section 2.2.1). The
282 Hilbert-Huang transform (HHT) is computed using 64 frequency bins between 2 and 35 Hz.
283 A wavelet transform was computed using a 5-cycle Morlet basis computed using the same 64
284 frequencies as the HHT.

285

286 The timing of individual oscillatory cycles is identified using the phase jumps in the
287 instantaneous phase time course where the oscillatory amplitude was above a threshold of
288 0.04 and instantaneous frequency below 18Hz. We compute control-point ratios (section
289 2.2.3), time locked instantaneous frequency profile and phase aligned instantaneous
290 frequency profile (section 2.2.4) across all included cycles (see section 2.2.2) of IMF-3.

291

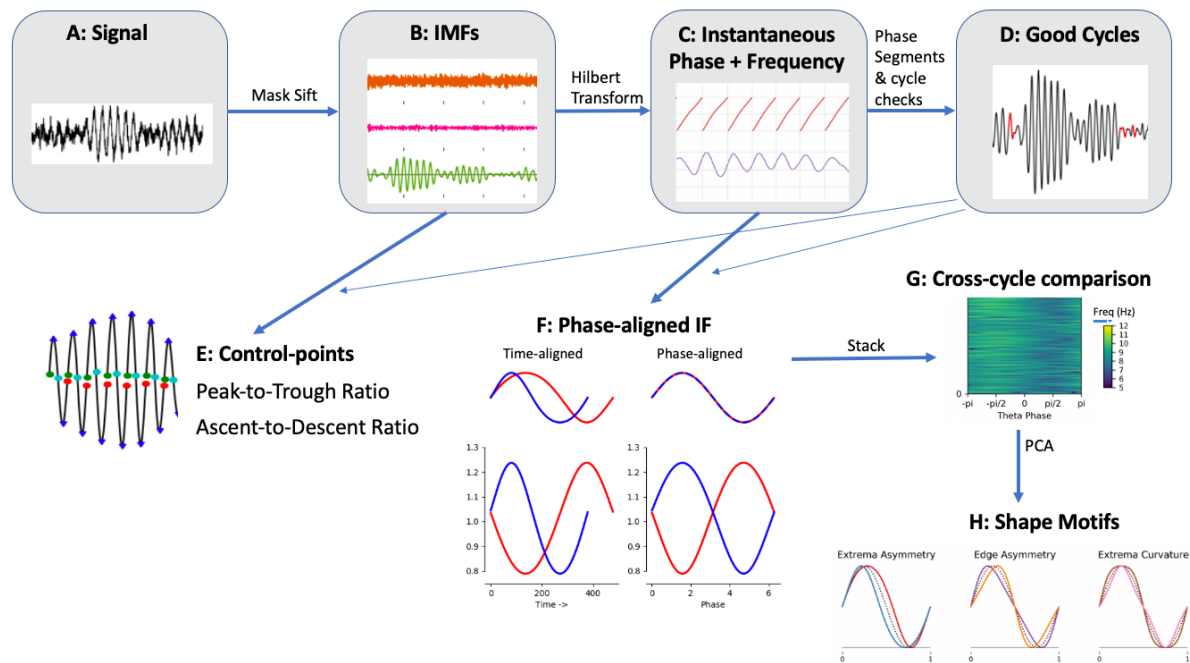
292 The sift, frequency transform, cycle detection and shape metrics are computed for both the
293 linear and non-linear systems defined in section 2.3.2 The difference between the linear and
294 non-linear systems is quantified using a separate t-test on the peak-to-trough and ascending-
295 to-descending control point ratios, and for each point in phase across the phase-aligned
296 instantaneous frequency profiles. Finally, the normalised waveforms are computed from the
297 average phase-aligned instantaneous frequency profiles from each system.

298

299 **2.4: Hippocampal theta analyses**

300

301 Local field potentials (LFPs) were recorded from the pyramidal layer of hippocampal CA1
302 using multi-channel tetrodes (Lopes-dos-Santos et al., 2018). Recordings were made during
303 open-field exploration in both familiar and novel environments across six recordings taken
304 over three recording days from each of three mice. Further data acquisition details can be
305 found in appendix 1.



306

307 **Figure 1: Overview of analysis applied to LFP recordings of hippocampal theta.**

308 A: The raw input LFP recording.

309 B: The raw signal is split into Intrinsic Mode Functions (IMFs) using a mask sift.

310 C: An instantaneous phase and instantaneous frequency time-course is estimated from the
311 theta IMF using the Hilbert Transform.

312 D: Cycle start and stop times are identified from jumps in the wrapped phase time-course and
313 'bad' cycles with distortions of reversals in phase are identified and removed.

314 E: Control points (peaks, troughs, ascending zero crossings and descending zero crossings)
315 are estimated from the good cycles within the theta IMF. Shape is then summarised using
316 peak-to-trough and ascending-to-descending duration ratios.

317 F: The instantaneous frequency of each good cycle is phase aligned to correct for variability
318 in cycle duration and internal cycle timings.

319 G: The phase-aligned cycles are stacked into a single array to allow for straightforward
320 comparisons between cycles.

321 H: A set of shape motifs are identified from the phase aligned IF using PCA.

322 2.4.1: Mask sift and Frequency transform.

323

324 LFP recordings were each separated into oscillatory components using the mask sift (Figure 1
325 A & B) with masks placed at $f_m = [350, 200, 70, 40, 30, 7, 1 \text{ Hz}]$. These masks were
326 selected to capture components with frequencies above the mask frequency down to around
327 $f_m * 0.7$ (Fosso and Molinas, 2017; Rilling and Flandrin, 2008). Keeping the masks constant
328 across recordings ensures that the frequency content of each IMF will be comparable across
329 recordings. For instance, we isolate the theta oscillation in IMF-6 using a mask frequency of
330 7 Hz. These mask frequency parameters were validated by rerunning the phase-aligned
331 instantaneous frequency analyses with jittered mask-frequencies (Supplemental section 8.2;
332 figure S2), showing that the theta waveform description is robust to reasonable changes to the
333 mask frequency values. These mask frequencies were effective in this set of CA1 LFP
334 recordings, but it is expected that a different set of mask frequencies would be needed to
335 analyse time-series containing different oscillatory dynamics. Next, a frequency
336 transformation was computed for each IMF using the Hilbert transform and the methods from
337 section 2.2.1 (Figure 1C).

338

339 2.4.2: Cycle detection

340

341 To ensure that the detected theta cycles are physiologically interpretable theta activity, we
342 identified cycles in each recording during times where the speed of movement of the mouse
343 was greater than 1 cm/second. As faster movement is associated with stronger theta
344 oscillations, this restriction increases the probability that our detected cycles represent
345 physiologically interpretable theta events. We additionally restricted analyses to cycles in
346 IMF-6 where cycle duration corresponded to 4-11 Hz frequency range (i.e., 312 and 113
347 samples respectively) and cycle amplitude was above the bottom 10% of the amplitude
348 distribution. Finally, cycles that failed the cycle inclusion checks outlined in section 2.2.2
349 were removed from analysis at this point (Figure 1D).

350

351 2.4.3: Cycle comparisons

352

353 We computed the temporally aligned instantaneous frequency profile, phase-aligned
354 instantaneous frequency profile and normalised waveform for each included theta cycle. The
355 average waveform shape within each dataset was estimated from the averaged phase aligned
356 instantaneous frequency and a group average constructed from the mean of the six individual
357 runs. Variability in waveform shape across single cycles in the group data is summarised
358 using the instantaneous frequency mean vector (section 2.2.5) and visualised as a distribution
359 in the complex plane in which the x-axis represents asymmetry between ascending and
360 descending edge frequency and the y-axis represents asymmetry between peak and trough
361 frequency. For comparison, we also identified the control points from each cycle of the theta
362 IMF and constructed the peak-to-trough and ascending-to-descending duration ratios (Cole
363 and Voytek, 2019).

364

365 2.4.4: Waveform Motifs and Relation to behaviour

366

367 We next look to explore the waveform shapes which are present in the phase-aligned
368 instantaneous frequency values. We use PCA (section 2.2.6) to identify the data-driven set of
369 shape components that explain the most variance in the shape of theta cycles in this dataset.
370 The first four principal components explaining 95% of variance defined our four shape motifs
371 and were retained for further analysis. The reproducibility of the PCA is validated across 500
372 split-half iterations assessing the proportion of variance explained by each PC and the
373 correspondence between the component shape in the two halves (Supplemental section 8.3;
374 figure S3).

375

376 The relationship between the shape motifs and a set of three cycle covariates (i.e., cycle
377 amplitude, cycle duration and mouse movement speed) was quantified using a General Linear
378 Model (GLM). The GLM was created with a design matrix containing the mean and the three
379 z-transformed covariates. These predictors were used to model the between cycle variability
380 in the principal component (PC) scores for each shape component in turn. This resulted in
381 four GLMs each fitting four parameter estimates. The t-statistic of each parameter estimate
382 was computed, and statistical significance established using a row-shuffle non-parametric
383 permutation scheme. 5000 permutations were computed for each PC-motif and dependent
384 variable before statistical significance determined at $p < 0.01$.

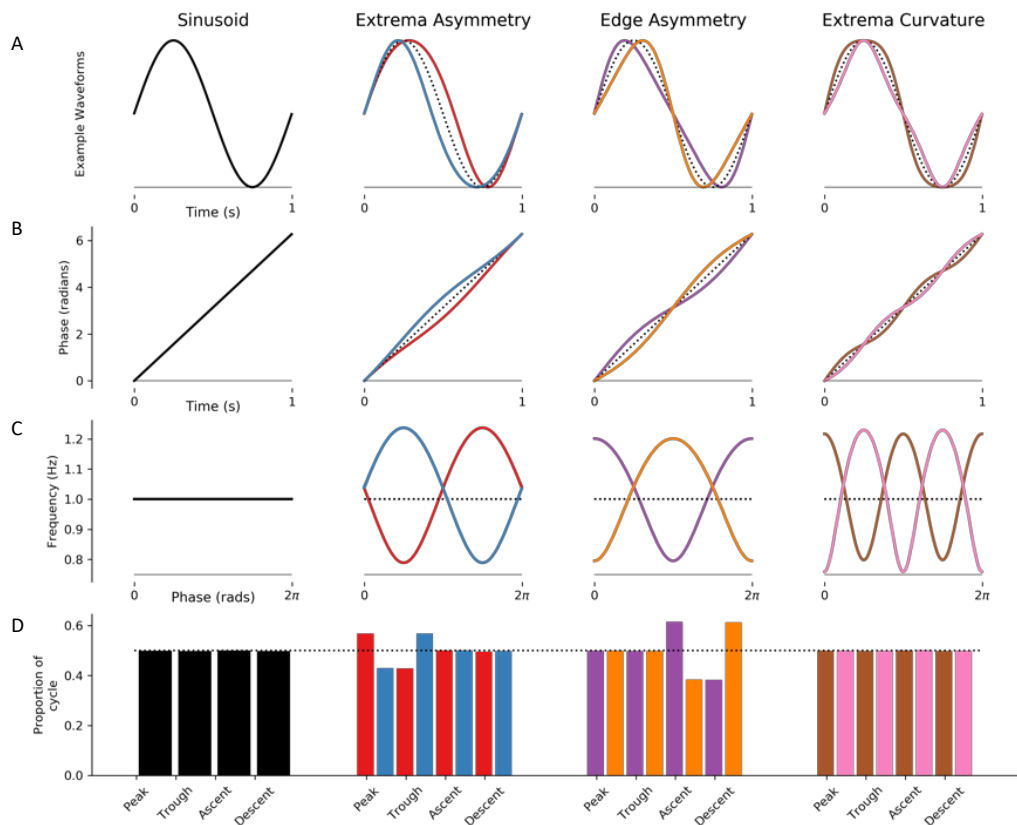
385 **3: Results**

386

387 **3.1: Instantaneous Frequency tracks waveform shape**

388

389 Figure 2 illustrates how instantaneous frequency reflects waveform shape in a set of noiseless
390 simulated cycles (see section 2.3.1). A sinusoidal cycle (Figure 2A) has a monotonically
391 progressing phase time-course which, in turn, has a flat instantaneous frequency profile.
392 Analysis of the duration of different segments reveals that the peak, trough, ascending edge
393 and descending edge all have the same duration. Cycles with a narrow peak, trough or
394 descending edge show corresponding changes in their instantaneous frequency (Figure 2 B &
395 C). Specifically, the longer duration, slower features correspond to a lower instantaneous
396 frequency. These instantaneous frequency profiles can describe a wide range of possible
397 shapes. For example, cycles in which both the peak and trough are widened or pinched lead
398 to instantaneous frequency profiles with multiple extrema (Figure 2: last column). While the
399 simple control-point metrics used here can track individual waveform features such as peak
400 or trough duration (Figure 2: bottom row), the quantification of more complex shapes would
401 require the definition of additional control points and shape metrics.



402

403 **Figure 2: Instantaneous Frequency changes with oscillatory waveform shape.**

404 The four columns illustrate examples of different simulated oscillatory cycles with distinct
405 waveform shapes.

406 A: The time-domain waveforms for each cycle. The first column shows a sinusoid, and the
407 remaining three columns show pairs of cycles with opposite waveform distortions (for
408 reference a sinusoid is shown as a dotted back line).

409 B: The instantaneous phase time course of the signals in the corresponding column.

410 C: The instantaneous frequency time course of the signals in the corresponding column.

411 D: The durations between different control points for each cycle the dotted line indicates the
412 expected duration for a sinusoid.

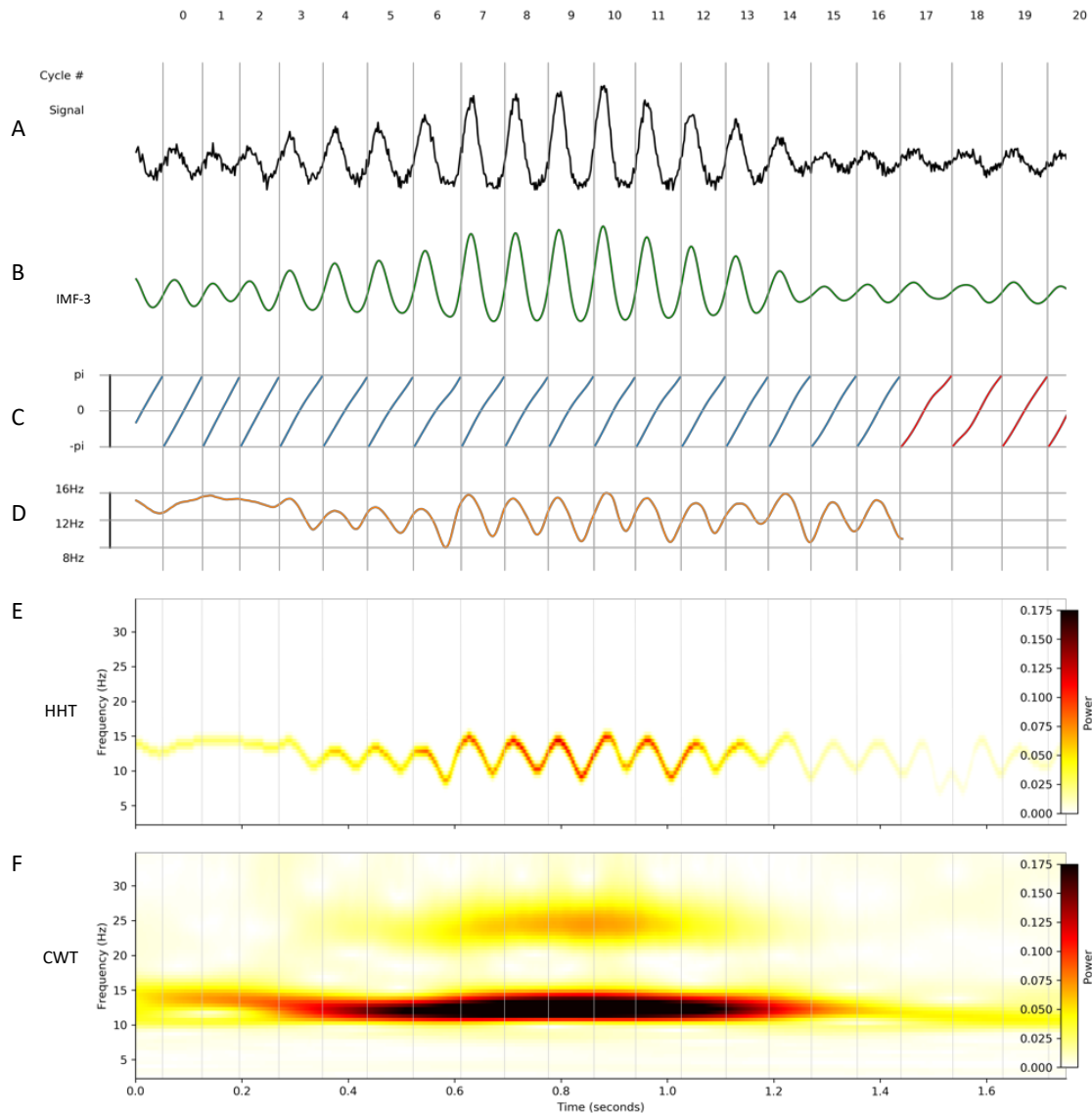
413 **3.2: Quantifying and comparing waveform shape in a simulated signal**

414

415 We next use simulations to illustrate how instantaneous frequency analyses can be conducted
416 on a noisy signal with a dynamic 12 Hz oscillation modified by a non-linearity which widens
417 the trough of each cycle (see section 2.3.2). This oscillation was isolated from the noisy
418 background using mask sift (Figure 3B) before the Hilbert transform was used to compute the
419 instantaneous phase time-course (Figure 3C). It is evident that the phase time-courses do not
420 progress linearly through all cycles; these deviations from monotonic phase progression are
421 quantified in the instantaneous frequency time-course (Figure 3D). Instantaneous frequency
422 sweeps within a single cycle reflect the non-sinusoidal shapes of the time-domain waveforms.
423 For this simulation, the instantaneous frequency tends to be higher during the first half of the
424 cycle and lower in the second half, reflecting the non-linearity that shortens the peak and
425 widens the trough of these oscillations.

426

427 The HHT of the simulated signal (Figure 3E) retains the high time-frequency resolution of
428 the instantaneous frequency time-course allowing within cycle frequency dynamics to be
429 visible. In contrast, while a standard 5-cycle Morlet wavelet transform identifies similar
430 power dynamics, variability in frequency within single cycles are not resolved (Figure 3F). A
431 further disadvantage is that the non-sinusoidal waveform shape of this simulation introduces
432 a 24Hz harmonic component into the wavelet transform.



433

434 **Figure 3: Instantaneous frequency analysis on a noisy non-sinusoidal oscillation.**

435

436 A simulated 12 Hz oscillation can be extracted from a noisy time-series and represented with
437 EMD, instantaneous frequency and the Hilbert-Huang Transform. Vertical grey lines denote
438 the starting times of individual cycles across the different panels.

439 A: The simulated noisy non-sinusoidal oscillation.

440 B: IMF-3 extracted from 'A' containing the simulated oscillation.

441 C: Instantaneous phase time-course of IMF-3. Cycles excluded from further analysis are
442 indicated in red. In this case, these cycles were below the amplitude threshold.

443 D: Instantaneous frequency time-course of IMF-3.

444 E: Hilbert-Huang Transform (HHT) of the simulated data segment.

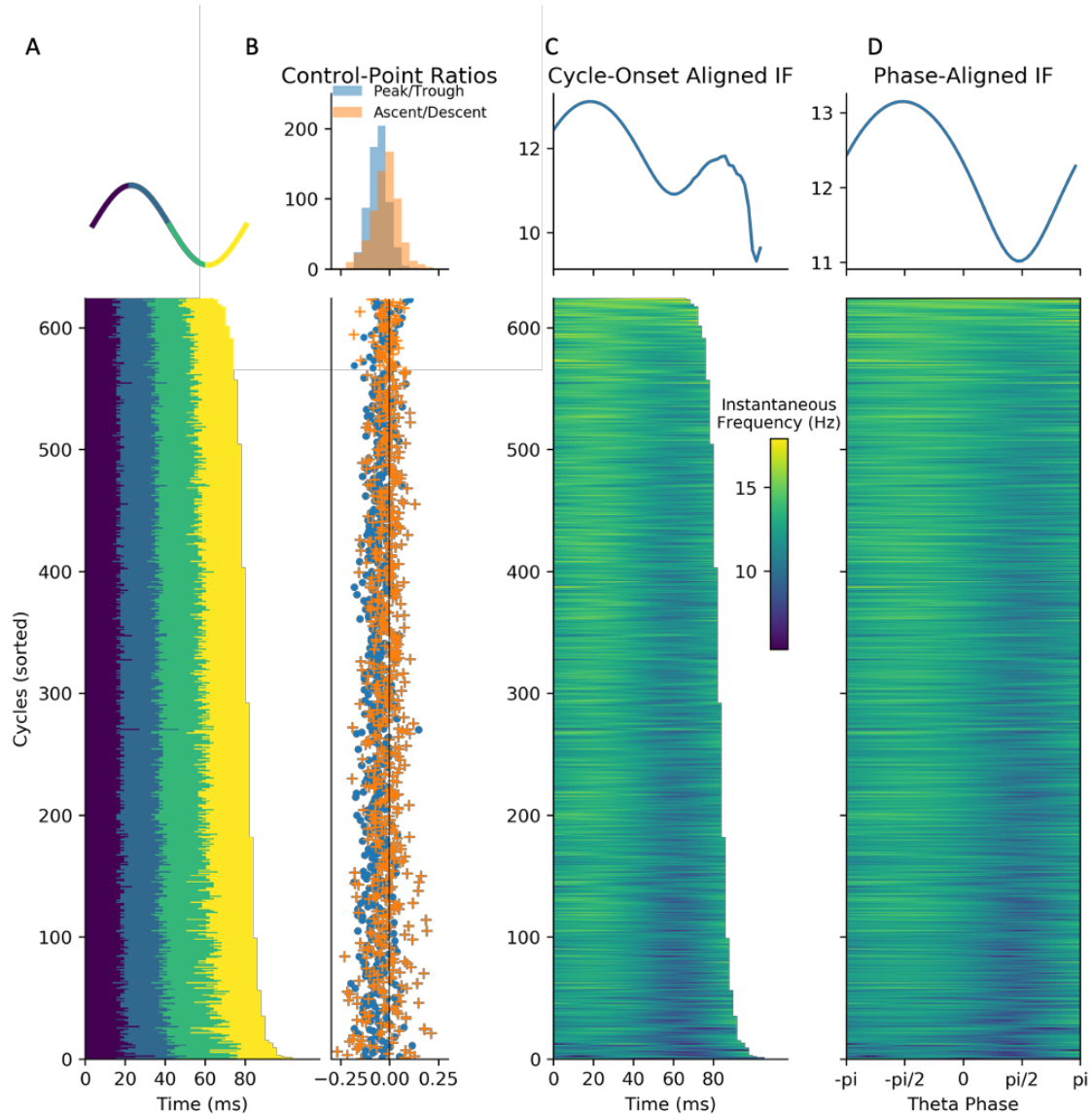
445 F: Continuous Wavelet Transform (CWT) of the simulated data segment.

446 Individual cycles of an oscillation play out at different rates leading to differences in the
447 timing of extrema within cycles and in overall cycle duration. These two sources of
448 variability hamper comparisons between individual oscillatory cycles. As outlined above
449 (section 2.2.3) one method to solve this issue is to discretise the cycle using a set of control
450 points before computing the proportion of time spent in different segments of the cycle
451 (Belluscio et al., 2012; Cole and Voytek, 2019). The cycle phase quartiles (ascending zero-
452 crossing, peak, descending zero-crossing and trough) of 500 cycles of the simulated signal is
453 shown in Figure 4A. The ratios of peak-to-trough duration and ascending-to-descending time
454 of these cycles suggests longer troughs and shorter peaks, whilst the ascending and
455 descending portions of the cycle are approximately equal in duration (Figure 4B).

456
457 An alternative approach for comparing cycles is to align the instantaneous frequency profiles
458 to one of the control points. For example, we aligned the 500 cycles to the ascending zero-
459 crossing and computed their time-locked average (Figure 4C). The time-locked instantaneous
460 frequency profile of these cycles is not flat, reflecting the presence of non-sinusoidal shape in
461 this simulated signal. However, the precise type of non-sinusoidal shape is ambiguous from
462 this average, due to variability in the location of different waveform features within single
463 cycles. In this case, the instantaneous frequency is highest around 10 samples after the
464 ascending zero-crossing; however, this time-lag might correspond to different points in the
465 waveform in different cycles. In addition, variability in the duration of cycles means that,
466 after a certain point, different numbers of cycles contribute to the average, making the
467 estimate unstable.

468
469 Here we present an approach that overcomes these shortcomings. In brief, phase-alignment
470 removes this ambiguity by visualising the instantaneous frequency of a cycle across a fixed
471 grid of points along its phase (see section 2.2.3). For instance, an oscillatory peak is
472 normalised to occur at the same phase value irrespective of the cycle's duration or shape.
473 This corresponds always to exactly one quarter of the phase of each cycle, but not necessarily
474 to one quarter of the duration of each cycle. By aligning the instantaneous frequency to the
475 phase, we remove the temporal distortions caused by varying shapes and cycle durations, and
476 express the shape with the phase-aligned instantaneous frequency values. The phase-aligned
477 instantaneous frequency of the simulated cycles (Figure 4D) now unambiguously shows the
478 increased frequency around the peak of the 12Hz oscillation and decreased frequency around
479 the trough. The average across the phase-aligned cycles is then a smooth representation of the
480 shape of the entire cycle.

481



482

483 **Figure 4: Methods for comparing waveform across cycles.**

484 Simulated data illustrating how waveform shape can be quantified using control point ratios,
485 instantaneous frequency and phase-alignment.

486 A : The durations between successive control points for each simulated cycle.

487 B : Top: distributions of peak-to-trough and ascent-to-descent durations. Bottom: The peak-
488 to-trough and ascent-to-descent ratios for each cycle.

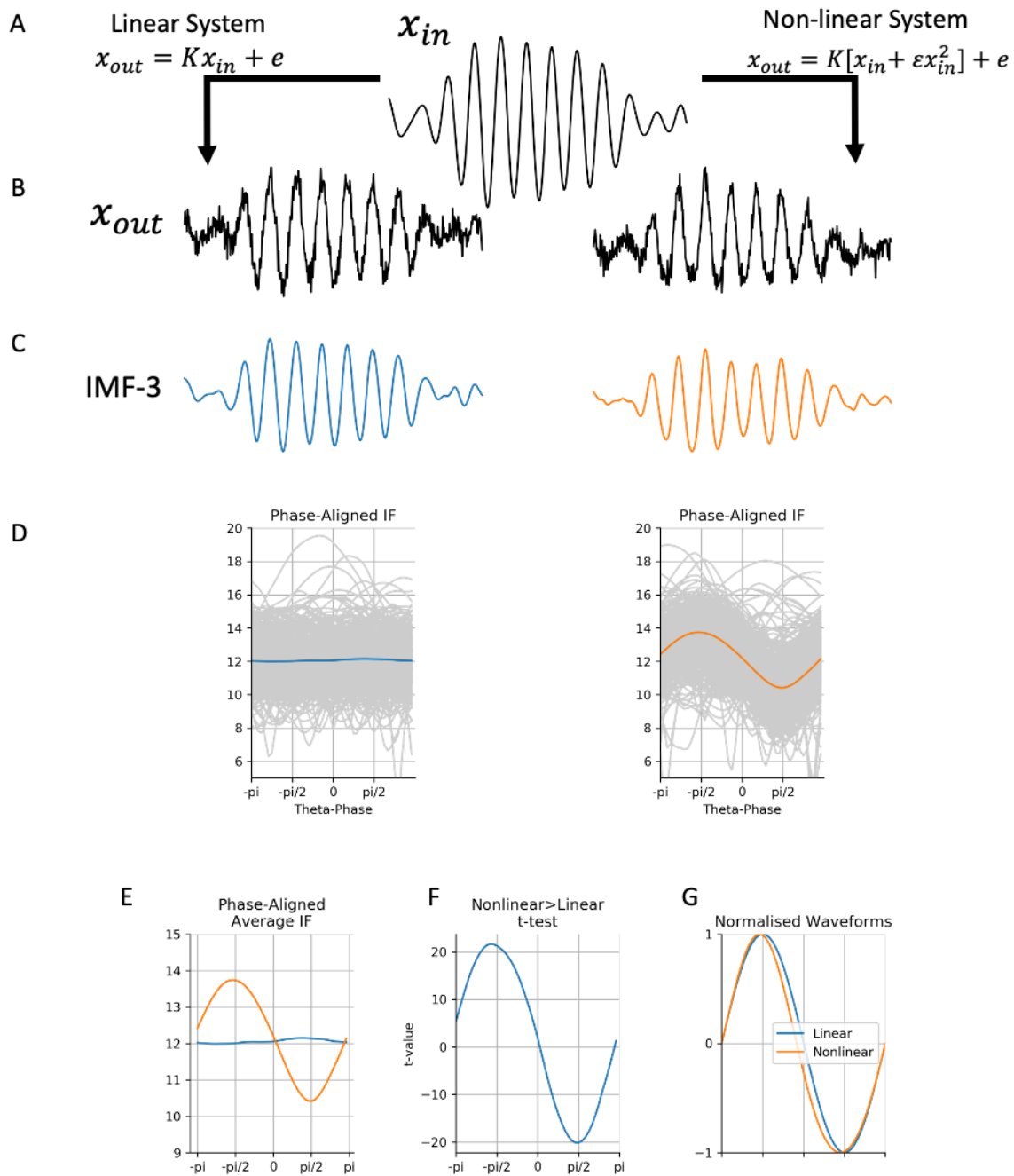
489 C : Top: the average temporally aligned instantaneous frequency profiles. Bottom the
490 temporally aligned instantaneous frequency profile for each cycle.

491 D : Top: the average phase-aligned instantaneous frequency profiles. Bottom: the phase-
492 aligned instantaneous frequency profile for each cycle.

493 Comparisons between sets of cycles is straightforward once waveform shape has been
494 estimated from instantaneous frequency and normalised through phase-alignment. This is
495 illustrated by contrasting the shape of a noisy a 12-Hz oscillation modulated by either a linear
496 or non-linear system (section 2.3.2; Figure 5A). The non-linear system has a wide-trough
497 shape whilst the linear system has a sinusoidal waveform (Figure 5B & C). The average
498 phase-aligned instantaneous frequency values for the linear system correspond to a flat line at
499 12 Hz throughout the cycle. In contrast the non-linear system has an increased instantaneous
500 frequency in the first half of the cycle and a decreased frequency in the second half (Figure
501 5D). We can compare the average instantaneous frequency profiles from both systems
502 (Figure 5E) and compute conventional t-statistics (Figure 5F) to quantify any differences in
503 waveform shape. In this simulation, we find that the non-linear system creates oscillations
504 with increased frequency peaks and decreased frequency troughs.

505

506 Finally, the phase-aligned average frequency profiles can be projected back into a
507 ‘normalised waveform’ to more intuitively visualise the type of non-sinusoidal distortions.
508 These normalised waveforms have a constant duration and an amplitude of one, but retain
509 any distortions in waveform shape quantified in the instantaneous frequency profile. For this
510 simulation, the normalised waveform reveals the pinched peak and the widened generated by
511 the non-linear system (Figure 5G).



512

513 **Figure 5: Comparing waveform shape in two simulated examples.**

514 A : A simulated oscillation is modulated by either a linear or a non-linear system.

515 B : Output oscillations of the two systems.

516 C : Oscillations recovered from the noisy simulation using EMD.

517 D : Average phase-aligned instantaneous frequency profiles for the two systems with
518 individual cycles in grey.

519 E : Average instantaneous frequency profiles from ‘D’ overlaid together.

520 F : t-values for a contrast between the instantaneous frequency in the two systems for each
521 point in phase.

522 G: Normalised waveforms for the two average frequency profiles in ‘E’.

523 **3.3: Characterising waveform shape in Hippocampal Theta**

524

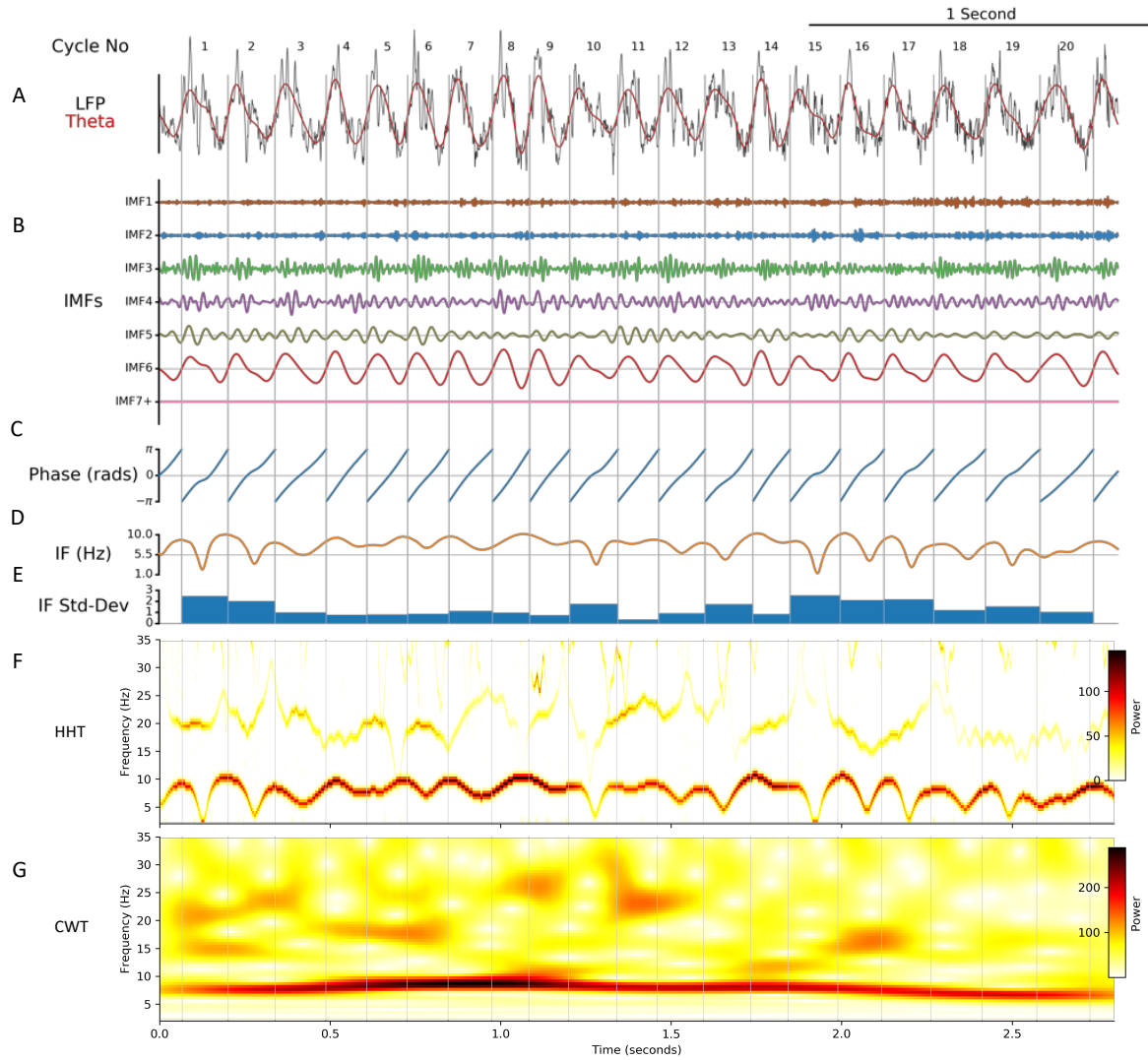
525 LFP data recorded from the mouse hippocampus were analysed to explore the utility of
526 phase-aligned instantaneous frequency as a measure of waveform shape. Figure 6A shows a
527 three-second LFP recording from the pyramidal layer of the mouse dorsal CA1 (black)
528 overlaid with the EMD-extracted theta IMF (red, Figure 6B shows all IMFs). In this case, the
529 theta oscillation was isolated into IMF 6 with minimal disruption to its amplitude or
530 waveform shape dynamics. Many of the theta cycles within this window have prominent non-
531 sinusoidal waveform shapes, which are qualitatively visible in both the raw data trace
532 (Buzsáki et al., 1986, 1985) and the EMD-extracted theta IMF. Importantly, the oscillatory
533 waveform shape varies between successive cycles, though the amplitude and duration of the
534 theta cycles are relatively consistent.

535

536 The instantaneous phase (Figure 6C) and instantaneous frequency (Figure 6D) were
537 computed from the theta oscillation in IMF-6. As with the simulation analysis, any within-
538 cycle dynamics in the instantaneous frequency naturally represent the waveform of each
539 cycle. This was summarised with the standard deviation of instantaneous frequency values
540 within each cycle (Figure 6E). As an illustration, cycles 5, 9 and 11 have relatively sinusoidal
541 shapes with flat instantaneous frequency profiles and low frequency variability. In contrast,
542 cycle 13 is relatively non-sinusoidal with a dynamic instantaneous frequency profile and high
543 frequency variability. The HHT provides a time-frequency description with sufficient
544 resolution to depict these within cycle instantaneous frequency sweeps (Figure 6F). In
545 contrast, a 5-cycle wavelet transform of the same data was not able to resolve these dynamics
546 (Figure 6G).

547

548 Looking at individual cycles illustrates how instantaneous frequency can characterise
549 waveform shape (Figure 7). Frequency increases and decreases correspond to slowing down
550 and speeding up of the cycle as its waveform shows non-sinusoidal behaviour. It is evident
551 that there are many observed shape profiles. For instance, the cycles labelled as *i* and *iii* in
552 figure 7 had frequencies that dip during the centre of the cycle, indicating an elongated, low
553 frequency descending edge. Cycle *v* had slowest frequency around $-\pi/2$ corresponding to a
554 wide peak. In contrast, cycle *vi* had relatively high frequency around $-\pi/2$ and lower
555 frequency at $+\pi/2$ leading to a short, pinched peak and an elongated trough. Overall, the
556 phase-aligned instantaneous frequency profiles and normalised waveforms provide a rich
557 description of oscillatory waveform, despite wide variability in cycle amplitude, duration and
558 shape.



559

560 **Figure 6: EMD analysis of a LFP segment containing hippocampal theta oscillations.**

561 A : A segment of a hippocampal LFP recording (black) overlaid with the extracted theta
562 oscillation (red).

563 B : IMFs extracted from this data segment using the mask-EMD. The theta oscillation is
564 isolated into the IMF-6.

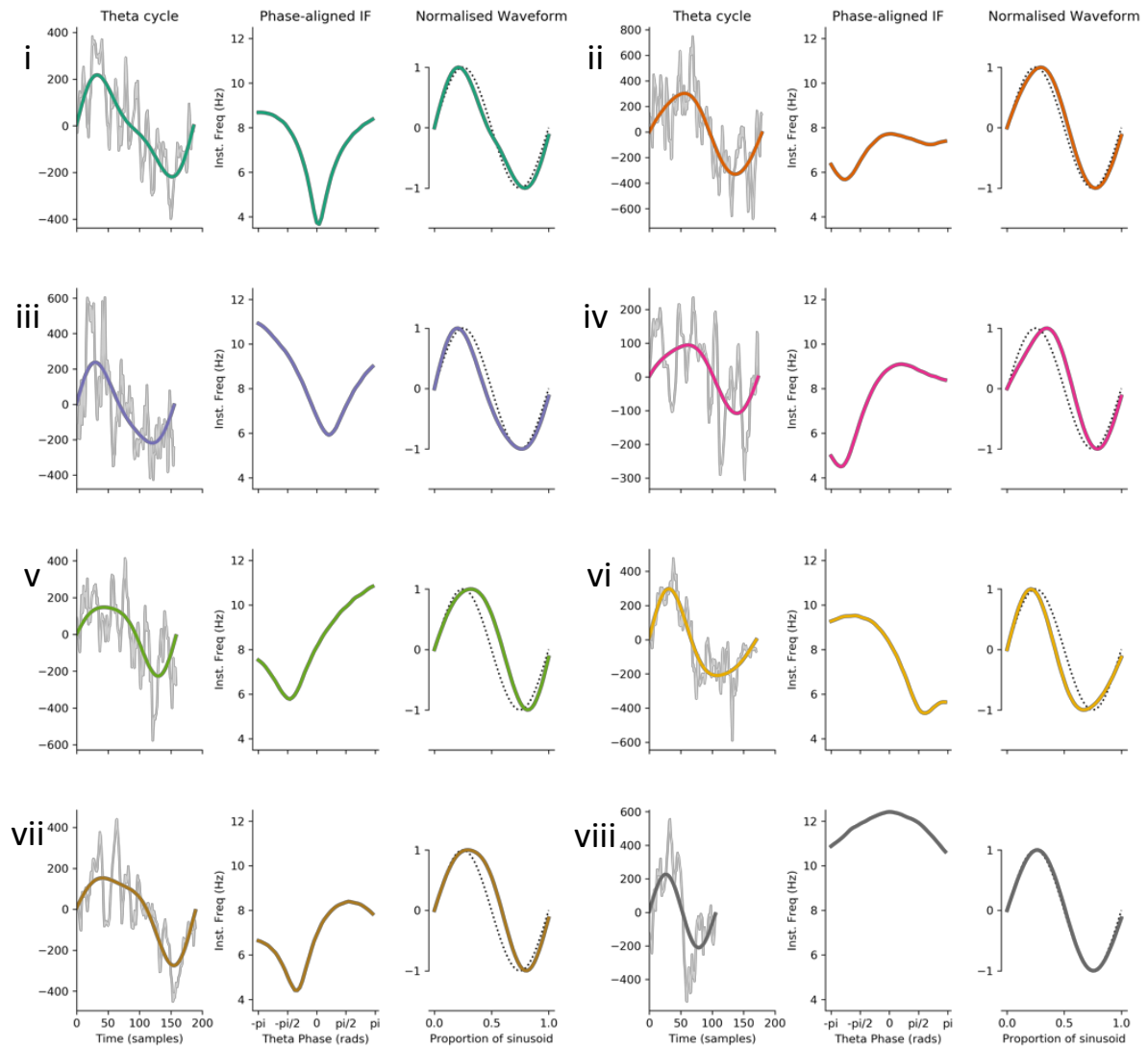
565 C : Instantaneous phase time-course of the theta IMF.

566 D : Instantaneous frequency time-course of the theta IMF.

567 E : Variability in instantaneous frequency for each theta cycle.

568 F : Hilbert-Huang Transform (HHT) of the LFP segment.

569 G : Continuous Wavelet Transform (CWT) of the LFP segment.



570

571 **Figure 7: Characterising shape in eight example cycles of hippocampal theta.**

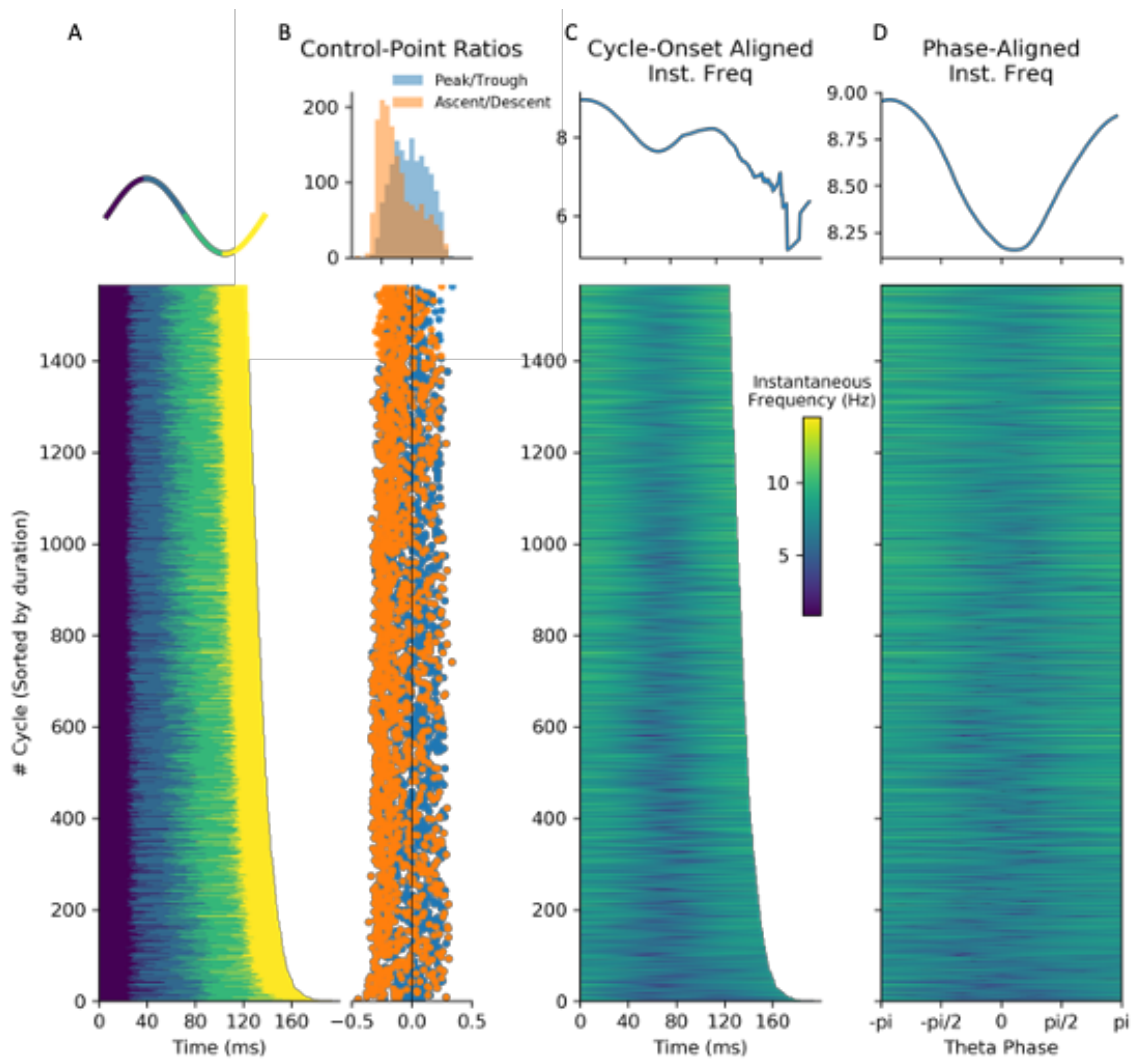
572

573 Eight representative theta cycles. For each example, the first subpanel shows the raw data
574 (grey) with the theta IMF super imposed (coloured line). The second subpanel shows the
575 phase aligned instantaneous frequency and third panel shows the normalised waveform
576 (coloured line) with a sinusoid for reference (black dotted line).

577 **3.4: Using phase-aligned instantaneous frequency to compare cycles.**

578

579 We next computed waveform shape from around 1500 hippocampal theta cycles from a
580 single recording using three different methods: control-point ratios (see section 2.2.6),
581 control-point locking, and phase-alignment (see section 2.2.3). The durations between
582 specified control points were computed for each cycle (Figure 8A) and the peak-to-trough
583 ratio and the ascent-to-descent ratio computed (Figure 8B). The peak-to-trough ratios are
584 evenly distributed around zero, whereas there is a bias in the ascent-to-descent ratios
585 suggesting that the descending edge of theta is longer than the ascending edge. The
586 instantaneous frequency profiles locked to the ascending zero-crossing show a wide variety
587 of shapes with a group average tendency for frequency to start around 9 Hz and to decrease
588 through the duration of the cycle (Figure 8C). As described above, this average effect is
589 challenging to interpret due to within-cycle variability in the timing of cycle features and
590 between cycle variability in total cycle duration. Our proposed phase-aligned instantaneous
591 frequency profiles (Figure 8D) resolve these ambiguities. This shows that theta cycle
592 instantaneous frequency in this single recording starts around 9 Hz at the ascending zero-
593 crossing, decreasing to around 8.1 Hz at the descending zero-crossing, before increasing
594 again to 9 Hz at the end of the cycle. This is consistent with a fast-ascending and slow
595 descending cycle shape revealed by the control point analysis and in previous literature. The
596 phase-aligned instantaneous frequency approach is able to show this effect as a continuous
597 shape profile for single cycles, which can be straightforwardly compared at the group level.



598

599 **Figure 8: Methods for quantifying waveform in hippocampal theta.**

600 A : Durations between successive control points for each simulated cycle.

601 B : Top: Distributions of peak-to-trough and ascent-to-descent durations. Bottom: Peak-to-

602 trough and ascent-to-descent ratios for each cycle.

603 C : Top: Average of temporally aligned instantaneous frequency profiles. Bottom:

604 Temporally aligned instantaneous frequency profile for each cycle.

605 D : Top: Average of phase-aligned instantaneous frequency profiles. Bottom: Phase-aligned

606 instantaneous frequency profile for each cycle.

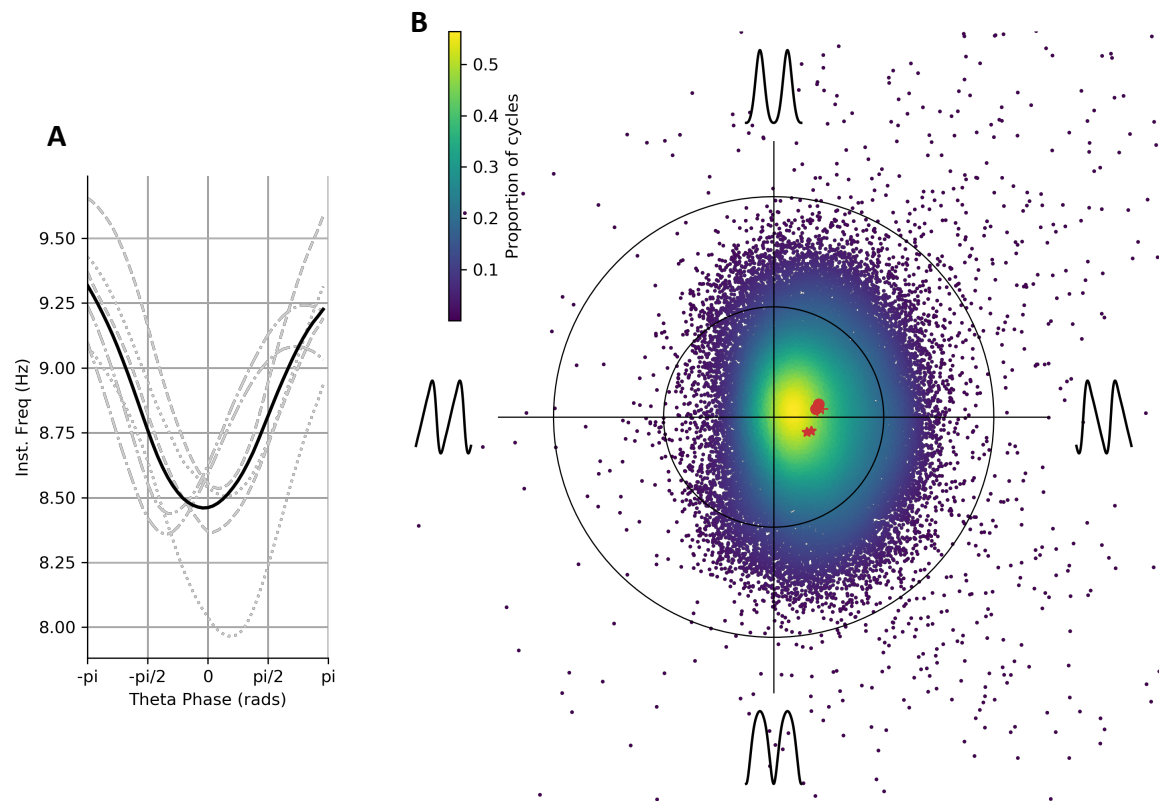
607 **3.5: Theta has a stereotyped asymmetric shape with wide variability over cycles**

608

609 We next summarised the average waveform across theta cycles from six recordings taken
610 from three mice. The average phase-aligned instantaneous-frequency profile is computed for
611 each recording and for the whole dataset. The overall group-level average waveform had a
612 cosine-type profile centred around an average instantaneous frequency of ~8.6 Hz (Figure
613 9A, average in black and individual recording sessions in grey). On average, the
614 instantaneous frequency peaked within the cycle around 9 Hz at the ascending zero-crossing
615 and drops to just below 8.4 Hz between the peak and descending zero-crossing. These results
616 are consistent with previous studies showing an asymmetry between the fast-rising and slow-
617 decaying halves of a theta cycle (Belluscio et al., 2012; Buzsáki et al., 1986; Cole and
618 Voytek, 2019). All six recordings across three animals showed a shape with a maximum
619 frequency around the ascending zero-crossing and a minimum on the descending edge,
620 though there was some variability in whether the lowest frequency was closer to the peak or
621 trough.

622

623 To visualise the variability in waveform shape across cycles and recording sessions, we
624 performed a complementary analysis using the instantaneous frequency mean-vector to see
625 the distribution of single-cycle waveforms across a simplified 2-dimensional shape-space
626 (Figure 9B; section 2.2.4). The distribution had a non-zero mean on the x-axis for all
627 recordings, indicating that the highest frequencies in a cycle are typically at the ascending
628 edge, consistent with the average in Figure 9A and with previous literature on the theta cycle
629 (Belluscio et al., 2012). Though the overall mean shift in the distribution of cycles is robust,
630 there is substantial cycle-to-cycle variability indicated by the width of the distribution.



631

632 **Figure 9: Average waveform shape and variability in shape across cycles in**
633 **hippocampal theta.**

634 A : Average of phase-aligned instantaneous frequency profiles for each of the six separate
635 recording sessions across three mice. The different dashed line styles indicate the runs from
636 the different mice and the solid black line represents the average across all six recordings.

637 B : Each individual cycle is projected into a simplified 'shape-space' to visualise the overall
638 variability in waveform shape around the average. Individual recording averages are shown
639 in red with different symbols representing the three animals.

640 **3.6: Distinct waveform motifs are differentially related to behavioural and**
641 **electrophysiological states.**

642

643 To further describe the variability in waveform shape across cycles and characterise its
644 relation to movement speed, theta amplitude and theta cycle duration we identify a set of
645 waveform shape ‘motifs’ using PCA. The PC values define a set of shape motifs that are each
646 expressed to different degrees in the observed cycles (see section 2.2.5 and 2.4.4). The first
647 four components describing 96% of variance are retained for further analysis (Figure 10).
648 The waveform shape represented by each PC motif is summarised by the normalised
649 waveforms (Figure 10A). These normalised waveforms are computed from the instantaneous
650 frequency profiles of the PC component-vectors (Figure 10B) projected onto the extreme
651 ends of the PC score distribution (Figure 10C). The shape of each individual cycle can then
652 be described by a set of four PC scores, relating to the amount of each component which it
653 contains.

654

655 PC-1 (63.31% of variance) describes a continuum of shape from a sharp peak and wide
656 trough through to a wide peak and sharp trough. This shape is similar to the y-axis in the
657 mean vector distribution in Figure 9B. In contrast, PC-2 (22.56% of variance) describes
658 shapes ranging between an elongated ascending edge and an elongated descending edge,
659 similar to the x-axis of Figure 9B. The remaining components describe more complex shapes
660 with relatively small contributions to the variance explained. PC-3 (6.86% of variance)
661 captures shapes with a left or right ‘tilt’ around their extrema and PC-4 (3.33% of variance)
662 describes shapes with a sharper or flatter curvature around the extrema.

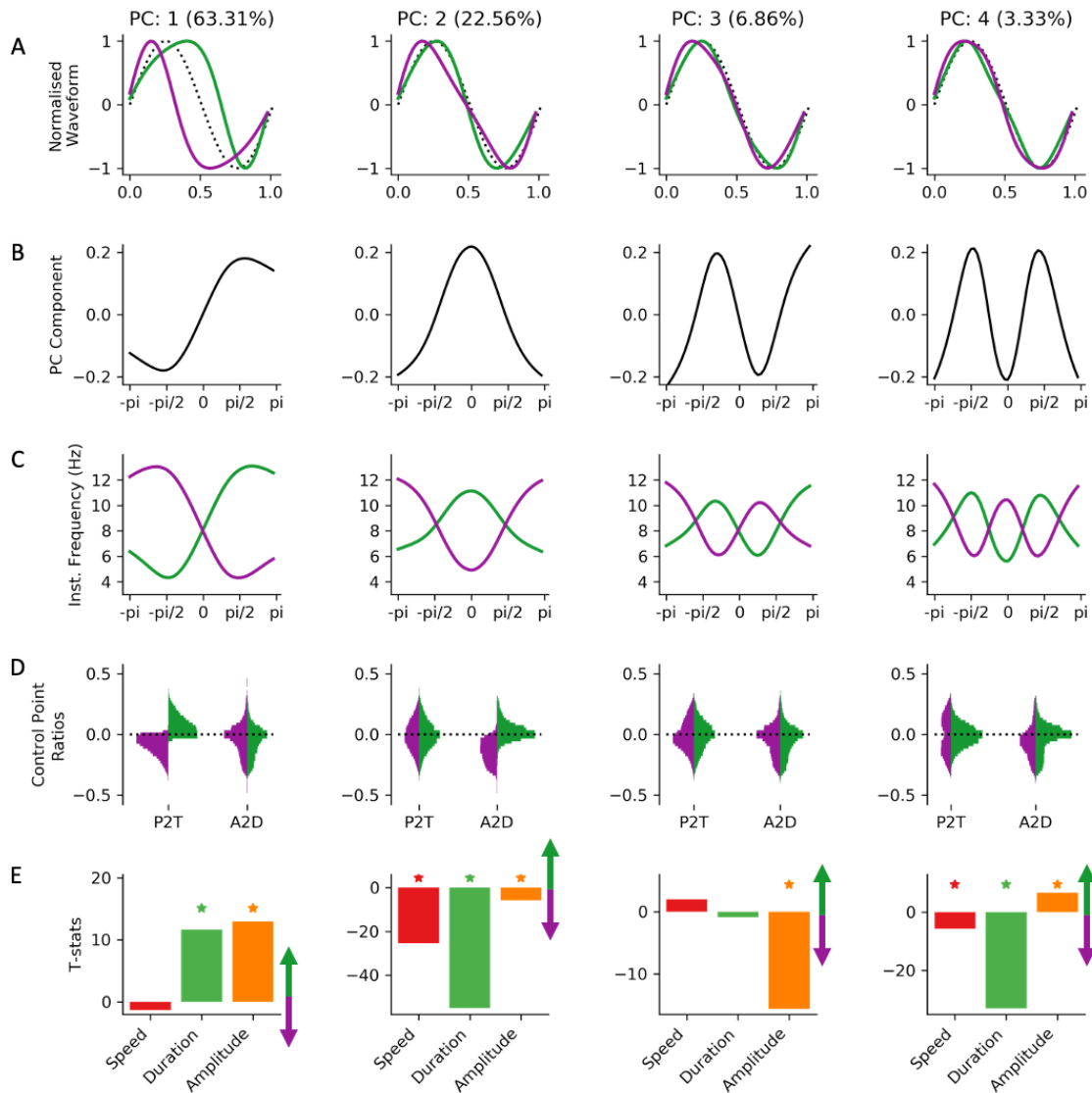
663

664 The control-point-based ascending-to-descending ratio and peak-to-trough ratio are computed
665 for each cycle. For each PC, these values are partitioned into cycles with positive or negative
666 PC scores (relating to distinct ends of the shape continuum for that component) and their
667 distributions plotted in Figure 10D. The peak-to-trough ratios are clearly separated in the two
668 ends of PC-1 whilst the ascending-to-descending ratios are similar for cycles with a positive
669 or negative score in PC-1. This is consistent with the normalised waveforms summarising
670 PC-1 in Figure 10A. PC-2 also shows the expected separation of ascending-to-descending
671 ratios by PC score, whilst the peak-to-trough ratios are unchanged. Whilst PC-3 and PC-4
672 describe around 10% of shape variability, they are not characterised by the control point
673 analyses. Neither peak-to-trough ratios nor ascending-to-descending ratios are changed by PC
674 score for PC-3 or PC-4. These shape profiles are robustly identified by the phase-aligned
675 instantaneous frequency method but are not distinguished by these control point-based
676 metrics as the shape distortions in PC-3 and PC-4 occur between the four specified control
677 points.

678

679 A general linear model was used to quantify the relationship between the different shape
680 motifs and theta amplitude, theta duration and mouse movement speed. This regression is
681 computed separately for each PC and the resulting parameter estimates converted into t-
682 statistics. PC-1 codes for changes in average instantaneous frequency across the cycle with a

683 small shape distortion around the descending edge. This PC has a strong relationship with
684 cycle duration and amplitude but no significant covariation with movement speed. Longer
685 and higher amplitude cycles tend to have more positive scores in PC-1 relating to wide peak
686 shapes. PC-2 has a strong relationship with duration and movement speed. Specifically,
687 cycles with elongated descending edges have longer cycle durations and are more likely to
688 occur during fast animal movement. PC-3 shows significant covariance with cycle amplitude.
689 High amplitude cycles tend to have shapes in which instantaneous frequency is relatively
690 high just before the peak or trough. Finally, PC-4 varies strongly with duration and weakly
691 with movement speed. Cycles with flatter curvatures around the extrema have longer cycle
692 durations and are less likely to occur during faster animal movement.



693

694 **Figure 10: Shape motifs in hippocampal theta and their relation to movement speed.**

695

696 A : The normalised waveforms for the first four shape motifs identified from a PCA across
 697 all phase-aligned instantaneous frequency profiles. Waveforms for positive PC scores are
 698 shown in purple and waveforms for negative scores shown in green with a sinusoid for
 699 reference (black dotted line).

700 B : PC for each motif.

701 C : Instantaneous frequency profiles of each shape motif created by multiplying the PC shape
 702 in 'B' with the maximum or minimum observed PC score for that PC and adding the mean.
 703 Purple profiles represent the positive end of the score distribution and green profiles represent
 704 the negative end.

705 D : Control-point ratios for cycles split by the sign of the PC score. Purple profiles represent
 706 the positive end of the score distribution and green profiles represent the negative end.

707 E : t-value of a GLM modelling the PC score for each motif as a function of movement
 708 speed, theta cycle duration and theta cycle amplitude. Asterisks indicate statistical
 709 significance at $p > 0.01$ as identified by non-parametric permutations.

710 **4: Discussion**

711

712 Non-sinusoidal waveforms are often visible by eye in raw LFP traces of electrophysiological
713 datasets, yet discovering and quantifying these non-sinusoidal and non-linear features present
714 substantial analytic challenges. We utilise within-cycle variability in instantaneous frequency
715 to describe distortions in waveform (Huang et al., 2009). Further, we introduce phase-
716 alignment as a solution to comparing full-resolution waveforms between cycles of different
717 durations. Taken together, we establish that the phase-aligned instantaneous frequency profile
718 of an oscillation provides a flexible framework for complete characterisation of oscillatory
719 waveform shape. We demonstrate the utility of this approach by applying it to simulated data
720 and LFP recordings of theta oscillations of behaving mice.

721

722 In real data, we observed that theta oscillations have, on average, a fast-ascending and slow
723 descending waveform, in line with previous reports (Belluscio et al., 2012; Buzsáki et al.,
724 1986, 1985; Cole and Voytek, 2019). Though this average shape is robust across many
725 cycles, recording sessions and animals; the shape of individual cycles is highly variable. We
726 characterise this variability using PCA to identify a range of shape components, or ‘shape
727 motifs’, which maximally explain the variability in the dataset. The first two PCs quantify the
728 relative durations of the peak and trough (PC-1) and the ascending to descending edge (PC-
729 2). These PCs broadly map onto the features described by the peak-to-trough and ascending-
730 to-descending control point ratios. We show that these theta shape PCs have distinct patterns
731 of covariation with movement speed, theta amplitude and theta cycle duration. Critically, we
732 show that though PC-2 describes less variability overall, it most clearly co-varies with
733 movement speed.

734

735 PC-3 and PC-4 capture more complex waveform shapes. We show that the curvature around
736 the extrema of the waveform shape (PC-4) is wider in theta cycles occurring during faster
737 animal movement. This shape is naturally described by instantaneous frequency but not
738 visible to standard ascending-to-descending and peak-to-trough control-point ratios. More
739 generally, if the waveform shape of interest is known *a priori*, it is possible to construct
740 specific control point-based measures so that the waveform shape can be identified. For
741 instance, waveform sharpness can be explored by looking at the differential between the
742 extrema and the samples 5ms before and after (Cole et al., 2017). However, in real data, we
743 may not know the waveform shape of interest *a priori*, implying that many separate metrics
744 may need to be computed for each cycle. In contrast, the phase-aligned instantaneous
745 frequency can quantify any waveform shape as a within-cycle instantaneous frequency sweep
746 without pre-specifying the features which may be of interest.

747

748 The present results demonstrate that single-cycle dynamics in oscillations can be
749 meaningfully estimated using phase-aligned instantaneous frequency, and that specific shape
750 motifs are differentially related to the wider electrophysiological (theta amplitude and
751 duration) and behavioural (movement speed) context. Future models of theta function may
752 consider these dynamics in waveform shape that deviate from a canonical sinusoidal theta

753 template. Given that many sub-processes occur preferentially at different parts of the theta
754 cycle (Klausberger and Somogyi, 2008), we hypothesise that shape distortion may indicate or
755 reflect a change in the underlying theta-phase-nested sub-processes.

756

757 The outlined approach requires that each cycle is smooth in both its waveform and phase
758 profiles, as any jumps or discontinuities will lead to noisy or even negative instantaneous
759 frequency estimates. If the cycle is smooth, we can characterise very large distortions in
760 waveform shape as within-cycle dynamics in instantaneous frequency. Finally, we assume
761 that the features being analysed are well described as oscillations. If the features are non-
762 sinusoidal and non-oscillatory, such as spiking activity, then descriptions using the language
763 of frequency may not be appropriate. With these improvements and caveats in hand, this
764 approach is readily generalisable to other datasets and provides a flexible framework for
765 investigating waveform shape oscillating systems.

766

767 In conclusion, the full-cycle waveform of single cycles of hippocampal theta can be
768 quantified and explored with phase-aligned instantaneous frequency. We use this approach to
769 confirm the characteristic fast-ascending waveform of theta oscillations; and to additionally
770 reveal that this is highly variable on the single-cycle level. Moreover, we are able to link this
771 variability with behavioural and electrophysiological states, suggesting that waveform shape
772 is a relevant feature of neuronal oscillations alongside frequency, phase and amplitude.
773 Finally, whilst we have illustrated this approach with hippocampal theta oscillations, it is
774 likely that this methodology will readily generalise to neuronal oscillation in other brain
775 regions, frequency bands and contexts.

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777

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791

792 **6: Author contributions**

793

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803 Supervision: ACN and MWW.

804 Funding Acquisition: DD, ACN and MWW

805 7: References

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937 **8: Supplemental Methods**

938

939 **Experimental model and subject details**

940 Animals used were male adult (4–7 months old) C57BL/6J mice (Charles River, UK). All
941 animals had free access to water and food in a dedicated housing facility with a 12/12 h
942 light/dark cycle. They shared a cage with their littermates until the surgery. All experiments
943 involving animals were conducted according to the UK Animals (Scientific Procedures) Act
944 1986 under personal and project licenses issued by the Home Office following ethical review.
945

946 **Microdrive implantation**

947 Animals were implanted with a 10-12 tetrode microdrive during a surgical procedure
948 performed under deep anesthesia using isoflurane (0.5–2 %) and oxygen (2 l/min), with
949 analgesia (0.1 mg/kg vetergesic) provided before and after. Tetrodes were constructed by
950 twisting together four insulated tungsten wires (12 µm diameter, California Fine Wire) and
951 shortly heating them to bind them together in a single bundle. Each tetrode was attached to a
952 M1.0 screw to enable their independent movement. The drive was implanted under
953 stereotaxic control in reference to bregma (Lopes-dos-Santos et al., 2018). Tetrodes were
954 initially implanted above the CA1 pyramidal layer and their exposed parts were covered with
955 paraffin wax. The drive was then secured to the skull using dental cement. For extra stability,
956 stainless-steel anchor screws had first been inserted into the skull. Two of the anchor screws,
957 which were inserted above the cerebellum, were attached to 50 µm tungsten wires (California
958 Fine Wire) and served as ground and reference electrodes during the recordings. The
959 placement of the tetrodes in dorsal CA1 was confirmed by the electrophysiological profile of
960 the local field potentials in the hippocampal ripple frequency band.
961

961

962 **Recording procedures**

963 Recordings commenced following full recovery from the surgery. Each animal was
964 connected to the recording apparatus and familiarized with a high-walled box containing
965 home cage bedding and with one open-field enclosure (the familiar enclosure) over a period
966 of approximately seven days. During this period, tetrodes were gradually lowered to the
967 stratum oriens of the hippocampal CA1. On the morning of each recording day, tetrodes were
968 further lowered into the pyramidal cell layer in search of multi-unit spiking activity and
969 sharp-wave/ripple events (van de Ven et al., 2016). Tetrodes were not moved for at least 1.5
970 h before recordings started. For each recording day, the animal was exposed to various open-
971 field enclosures including the familiar, which the animal had repeatedly been exposed to
972 before, and a novel enclosure the animal had never seen before. The open-field enclosures
973 differed in shape and in the cue-cards that lined some of the walls. The present study includes
974 a total of six LFP recordings from three mice (including three familiar enclosure and three
975 novel enclosure sessions). At the end of each recording day, tetrodes were raised to the
976 stratum oriens to avoid damaging the pyramidal layer overnight.
977

977

978 Multichannel data acquisition and position tracking
979 The extracellular signals from the electrodes were buffered on the head of the animal (unity
980 gain op-amps, Axona Ltd) and transmitted over a single strand of litz wire to a dual stage
981 amplifier and band pass filter (gain 1000, pass band 0.1 Hz to 5 kHz; Sensorium Inc.,
982 Charlotte, VT), or (in other setups) the electrode signals were amplified, multiplexed, and
983 digitized using a single integrated circuit located on the head of the animal (RHD2164, Intan
984 Technologies, Los Angeles; pass band 0.09 Hz to 7.60 kHz). The amplified and filtered
985 electrophysiological signals were digitized at 20 kHz and saved to disk along with the
986 synchronization signals from the position tracking. LFPs were further down sampled to 1250
987 Hz for all subsequent analyses. In order to track the location of the animal three LED clusters
988 were attached to the electrode casing and captured at 39 frames per second by an overhead
989 color camera.

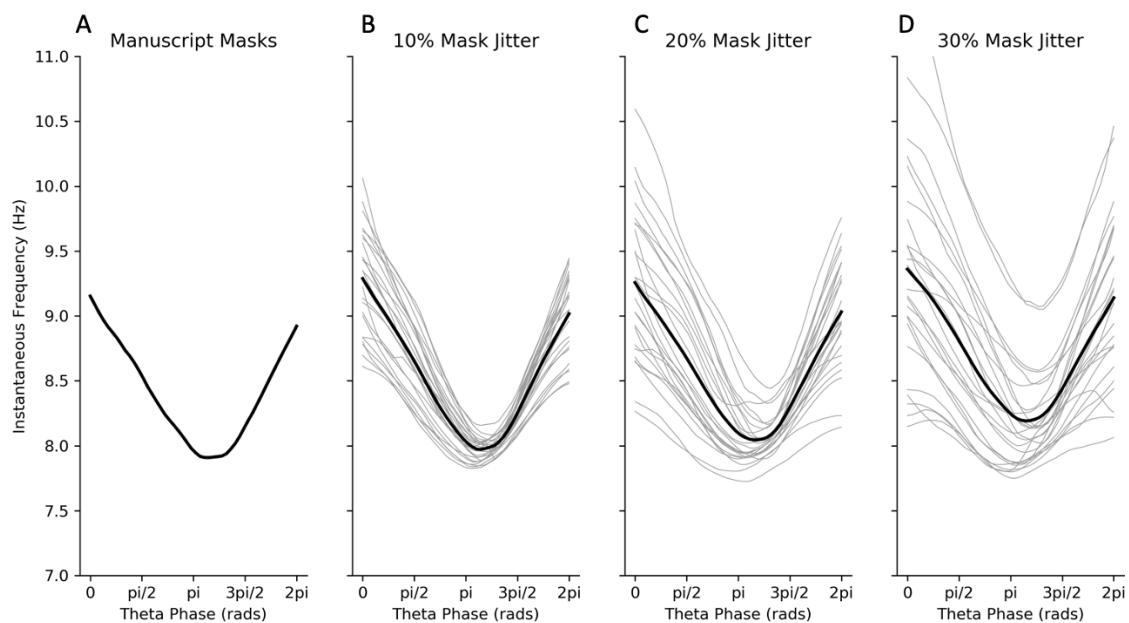
990 8.2: Mask sift parameter robustness

991

992 We ran a supplemental analysis to ensure that our LFP phase-aligned instantaneous
993 frequency results were robust to moderate changes in the masking parameters. 5 minutes of
994 data from a single data recording were repeatedly sifted with jittered masked frequencies. 25
995 sifts were computed for each of three mask frequency jitter values of 10%, 20% and 30%.
996 For example, in the 10% condition each iteration each mask frequency was randomised to a
997 value drawn from a uniform distribution between 90% and 110% of the original frequency.
998 The results showed that jitter of 10% has a small effect on the phase-aligned instantaneous
999 frequency values, though the centre frequency and shape profile remain consistent across all
1000 iterations. Jitters of 20% and 30% have larger effects on both centre frequency and shape on
1001 individual iterations, suggesting that some mask frequency combinations are having a large
1002 impact on the results. Despite this, the average across all iterations remains strikingly similar.
1003 These results indicate that the main waveform results are robust to moderate changes to the
1004 masking parameters.

1005

1006



1007

1008 **Figure S1 – Phase-aligned instantaneous frequency values across a range of jittered**
1009 **mask frequencies.**

1010 A: Instantaneous frequency for mask frequencies used in main analysis

1011 B: Instantaneous frequency for mask frequencies jittered by +/- 10%. Individual iterations are
1012 shown in grey and the average in black.

1013 C: As B for jitter of +/- 20%

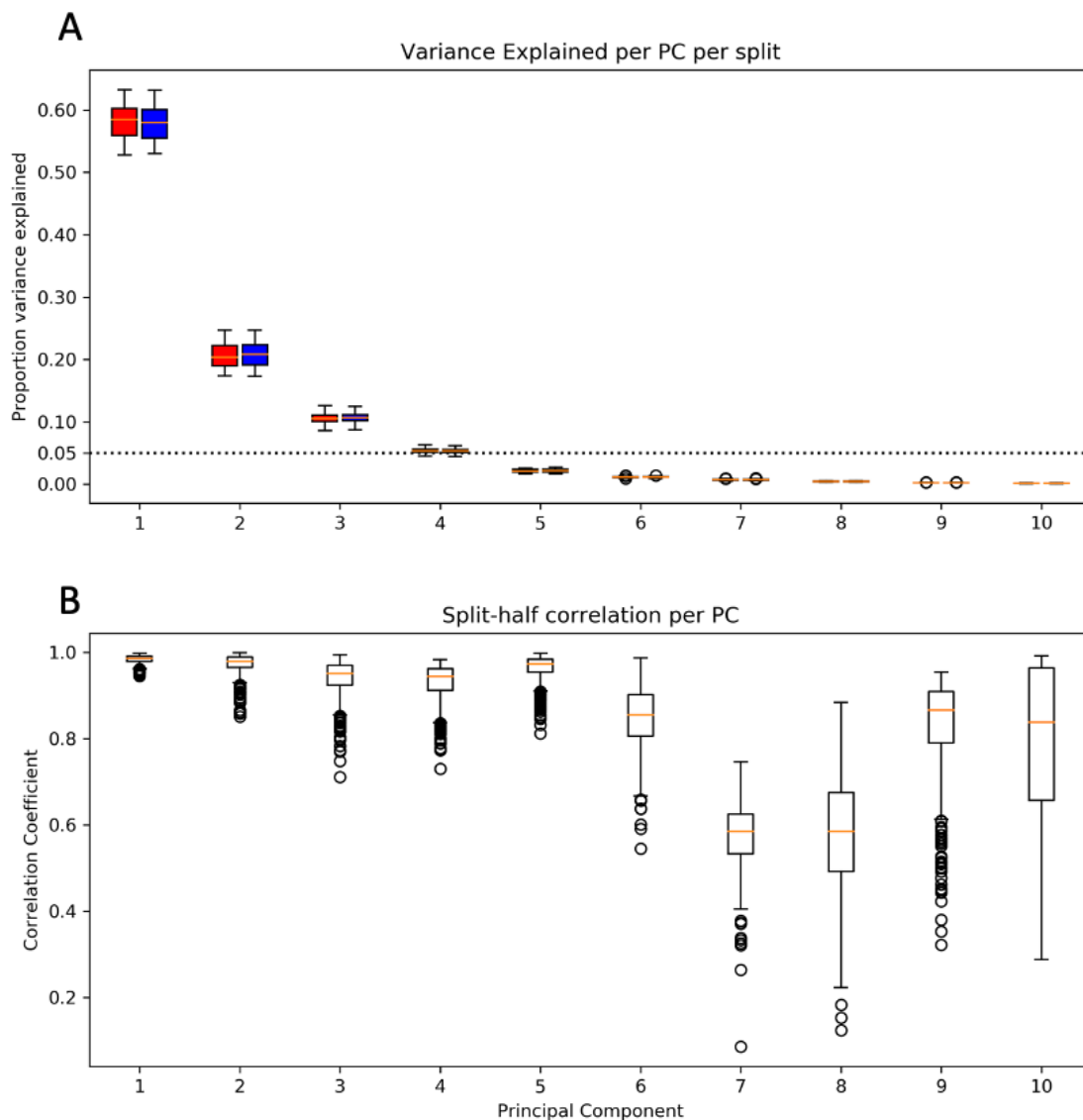
1014 D: As B for jitter of +/- 30%

1015 8.3: PCA component selection and reproducibility

1016

1017 The principal components analysis results were validated by computing the split half
1018 reproducibility of the PC components across 500 splits. The correlation of component shapes
1019 between the separate halves and the proportion of variance explained was computed for each
1020 split. The distribution of explained variance for each mode was highly reproducible across
1021 the 500 splits and the first four components explained more than 5% of overall variance
1022 (Figure S2A). The PC component shapes were also highly reliable for the first four
1023 components. The average correlation between components for the two halves of each split
1024 was over $r=0.95$ for the first 5 PCs. Based on these comparisons we carried the first four
1025 components forward for further analyses.

1026



1027

1028 **Figure S2 – the variance explained and split-half correlation distributions across 500**
1029 **split half iterations.**

1030 A: the variance explained by each component for each half over the 500 splits. The first half
1031 is in red and the second half in blue.

1032 B: The correlation in the component shape between the first and second half of the 500 splits.