

# Synchronization principles of gamma rhythms in monkey visual cortex

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**Neural synchronization<sup>1-5</sup> in the gamma-band (25-80Hz) can enhance and route information flow during sensory<sup>6-8</sup> and cognitive processing<sup>2,9-11</sup>. However, it is not understood how synchronization between neural groups is robustly achieved and regulated despite of large variability in the precise oscillation frequency<sup>10,12-16</sup>. A common belief is that continuous frequency matching over time is required for synchronization and that thus rhythms with different frequencies cannot establish preferred phase-relations. Here, by studying gamma rhythms in monkey visual area V1, we found that the temporal variation of the frequency difference was to the contrary essential for synchronization. Gamma rhythms synchronized by continuously varying their frequency difference in a phase-dependent manner. The synchronization level and the preferred phase-relation were determined by the amplitude and the mean of the frequency difference variations. Strikingly, stronger variation of the frequency difference led to stronger synchronization. These observations were reproduced by a biophysical model of gamma rhythms<sup>8,17-19</sup> and were explained within the theory of weakly coupled oscillators<sup>20-25</sup>. Using a single and general equation, we derived analytical predictions that precisely matched our V1 gamma data across different stimulus conditions. Our work reveals the principles of how gamma rhythms synchronize, where phase-dependent frequency variations play a central role. These frequency variations are characteristic for the intermittent synchronization regime, a non-stationary regime naturally occurring between the states of complete synchrony and asynchrony. This regime allows for synchronization between rhythms of variable frequencies, which is essential for achieving robust synchronization in the complex and noisy networks of the brain.**

35 To study the principles of neural synchronization we developed an experimental approach using  
36 V1 gamma oscillations. V1 gamma is subject to excellent experimental control, as it emerges  
37 locally and retinotopically, with a preferred oscillation frequency that can be readily manipulated  
38 by stimulus input properties<sup>14,16</sup>. Further, its generative mechanism is one of the best understood  
39 in the brain<sup>26,2</sup>. We first asked how synchronization within V1 was influenced by frequency  
40 differences, and by distance between recording sites. To this aim, we recorded from 2 to 3  
41 laminar probes simultaneously in cortical area V1 of two macaques (M1 and M2) (Fig.1A). We  
42 investigated distances on the order of magnitude of V1 horizontal connectivity<sup>27</sup> hence probes  
43 were separated by between 1 and 6mm. Using laminar probes enabled us to reduce volume  
44 conduction by calculating current-source density (CSD) as a network signal. Using CSD, we  
45 estimated the instantaneous frequency, phase and phase difference of gamma signals. The  
46 monkeys fixated centrally while a whole-field static grating was shown with spatially variable  
47 contrast. Gamma power was induced in layers 2-4 and in the deepest layer (Fig.1B, Fig.S2). V1  
48 locations showed increased gamma frequency with increased local contrast (Fig.1C). The  
49 frequency difference correlated with a difference in the MUA spike rate between probes  
50 (Fig.S4). Correspondingly, neurons recorded from different probes, in whose receptive fields  
51 (RFs) different contrasts were placed, showed different mean gamma frequencies (linear  
52 regression, single contact level, M1:  $R^2=0.31$ , M2:  $R^2= 0.25$ , both  $p<10^{-10}$ ), giving us control to  
53 parametrically vary the frequency difference between probes.

54 We will first show the key results through three illustrative examples. In the first example, we  
55 chose two cortical locations separated by a relatively large distance of ~5mm, presented with a  
56 visual contrast difference of 17% (Fig.1D). Their frequency difference was 5Hz as shown by  
57 their non-overlapping power spectra (Fig.1E). This would imply that the phase difference would  
58 not be constant, but would advance at a phase precession rate of  $2\pi$  every 200ms, which could be  
59 expected to preclude synchronization. However, the frequency difference was not constant.  
60 Instead, the instantaneous frequency difference was modulated as a function of phase difference  
61 (Fig.1F) with a modulation amplitude of 1Hz. At the smallest frequency difference (4Hz, yellow  
62 point) the phase precession was slowest, at  $2\pi$  every 250ms, meaning that the oscillators stayed  
63 relatively longer around that phase difference. As a result, the probability distribution of phase  
64 differences over time (Fig.1G) was non-uniform giving a phase-locking value<sup>28</sup> (PLV) of 0.11.  
65 The peak of the distribution, the ‘preferred phase’, was at 1.3rad, in line with the minimum of the  
66 instantaneous frequency modulation function. In the second example, we chose a pair with a  
67 similar frequency difference of 4.8Hz and a closer distance (~2.5mm, Fig.1H). The instantaneous  
68 frequency modulation amplitude was larger with 1.8Hz modulation amplitude (Fig.1J) with a  
69 modulation minimum around 3Hz at the preferred phase. Because phase precession at the  
70 preferred phase was slower, the phase difference distribution was narrower indicating higher  
71 synchrony (PLV=0.32, Fig.1K) with a peak centered at a different phase (0.78rad). In the third  
72 example the cortical distance remained the same and the frequency difference was reduced  
73 (2.8Hz) by eliminating the contrast difference (Fig.1M, the remaining frequency difference  
74 might be due eccentricity, see Fig.S4). The frequency modulation amplitude did not change

75 however, with a lower mean difference, the modulation minimum was close to zero (1Hz,  
76 Fig.1N), thus the associated phase difference (0.48rad) could be maintained for relatively longer  
77 periods and the phase difference probability distribution was even narrower (PLV=0.51, Fig.1O).  
78 The three examples were representative for the 1079 recorded contact pairs in monkey M1 and  
79 887 contact pairs in monkey M2.

80 We now show how the observed behavior can be accounted for within the mathematical  
81 framework of the theory of weakly coupled oscillators<sup>20-25</sup>, where V1 populations can be  
82 approximated as oscillators, ‘weakly coupled’ by horizontal connections. According to the  
83 theory, the synchronization of two coupled oscillators can be predicted from the forces they exert  
84 on each other as a function of their instantaneous phase difference. This interaction function is  
85 referred to as the phase-response curve (PRC). Accordingly, the phase precession of two given  
86 cortical V1 locations is reduced to:

$$87 \quad (1) \quad \dot{\theta} = \Delta\omega + \varepsilon G(\theta) + \eta$$

88 where  $\dot{\theta}$  is the time derivative of the phase difference  $\theta$  (the rate of phase precession),  $\Delta\omega$  the  
89 detuning (the frequency difference),  $\varepsilon$  the interaction strength,  $G(\theta)$  is defined as the mutual  
90 PRC, and  $\eta$  the phase noise, where  $\eta \sim N(0, \sigma)$ . Phase noise is here defined as variation,  
91 unrelated to interaction, that likely occurs for neural oscillators due to inherent instabilities of the  
92 generation mechanism<sup>13,12</sup> and due to other complex interactions occurring in cortical networks.  
93 For convenience, we express  $\omega$ ,  $\varepsilon$  and  $\eta$  in units of Hz (1Hz=2 $\pi$ \*rad/s). The time derivative  $\dot{\theta}$  is  
94 also expressed in Hz (instantaneous frequency, IF). Equation 1 was solved analytically to study  
95 changes in the phase-difference probability distribution, here characterized by the PLV and the  
96 mean (preferred) phase difference, as a function of detuning  $\Delta\omega$  and interaction strength  $\varepsilon$ .

97 The theory predicts that the PLV and the mean phase difference result from an interplay between  
98 the detuning and interaction strength (Fig.S1). When detuning is smaller than the interaction  
99 strength ( $\Delta\omega < \varepsilon$ ), the PLV is high and the mean phase difference is small. When detuning is  
100 larger than the interaction strength ( $\Delta\omega > \varepsilon$ ), the PLV is low and the mean phase difference is  
101 large. With stronger interaction strength  $\varepsilon$ , larger detuning  $\Delta\omega$  can be ‘tolerated’, leading to a  
102 triangular shaped region of high PLV in the  $\Delta\omega - \varepsilon$  space termed the “Arnold tongue”<sup>21,24,29</sup>.  
103 Oscillators start to phase precess due to detuning or due to destabilization by phase noise<sup>24</sup>.  
104 Oscillators show linear phase precession if uncoupled. If coupled, the phase precession is non-  
105 linear and modulated by phase difference through the PRC, leading to instantaneous frequency  
106 modulations ( $\Delta IF(\theta)$ ); a regime called intermittent synchronization<sup>24,30,31,29</sup>. In this regime, there  
107 is a preferred phase difference at which the instantaneous frequency difference is minimized.

108 The power of using the theory is the possibility to make precise predictions of the PLV and mean  
109 phase difference as a function of  $\Delta\omega$  and  $\varepsilon$ . According to equation 1, the time-averaged  
110 modulation of the instantaneous frequency difference by phase difference,  $\overline{\Delta IF(\theta)}$  (Fig.1F,J,N),  
111 directly relates to the deterministic term  $\Delta\omega + \varepsilon G(\theta)$ , as noise is averaged out. We used the

112 observed  $\Delta\overline{IF}(\theta)$  to estimate the underlying synchronization properties. We estimated a single  
113  $G(\theta)$  function and  $\sigma$  value for a given whole dataset assuming stability of underlying PRCs and  
114 of the noise sources, whereas  $\Delta\omega$  and  $\varepsilon$  was estimated for each contact pair and condition.  
115 Specifically, the  $G(\theta)$  was estimated by the modulation shape of the  $\Delta\overline{IF}(\theta)$  put to unity (see  
116 supplementary materials). We estimated interaction strength  $\varepsilon$  by the modulation amplitude of  
117 their  $\Delta\overline{IF}(\theta)$  and detuning  $\Delta\omega$  by the average value of their  $\Delta\overline{IF}(\theta)$  computed over  $[-\pi, \pi]$ . The  
118 remaining parameter  $\sigma$  was approximated by finding the  $\sigma$  value for equation 1 that could  
119 reproduce the observed overall instantaneous frequency variability (full description the theory of  
120 weakly coupled oscillators and parameter estimation in supplementary materials).

121 Before addressing whether the theory captures V1 gamma synchronization, we first tested our  
122 approach in a computational model of gamma oscillations in which the underlying network  
123 parameters were known. To that aim, we investigated simulations of two coupled Pyramidal-  
124 Interneuron gamma (PING) spiking networks<sup>8,17-19</sup> (Fig.2A). The gamma frequency was  
125 modulated by input drive<sup>17,18,32</sup> ( $R^2=0.98$ , Fig.2B), whereas the interaction strength  $\varepsilon$  was  
126 modulated by inter-network synaptic connectivity strength ( $R^2=0.97$ , green arrows, Fig.2A). For  
127 each network we estimated a population signal from which we extracted their instantaneous  
128 phase difference (Fig.2C). By reconstructing the  $G(\theta)$  (see example of  $\Delta\overline{IF}(\theta)$  in Fig.2D) and  
129 estimating phase noise variance ( $\sigma=15\text{Hz}$ ), we could solve equation 1 and we found that the  
130 theory (Fig.2F-I) accurately predicted the PLV (model accuracy:  $R^2=0.93$ ) and the mean phase  
131 difference (model accuracy:  $R^2=0.94$ ). Mapping the gamma PLV and the mean phase difference  
132 in the  $\Delta\omega - \varepsilon$  parameter space yielded the predicted Arnold tongue.

133 We then tested whether the theory predicted the in vivo data with equal success. In the same  
134 manner as with PING modeling data, we estimated for each monkey the  $G(\theta)$  (see representative  
135  $\Delta\overline{IF}(\theta)$  examples in Fig.3A and F) and the phase noise variance (M1: $\sigma=19\text{Hz}$ , M2: $\sigma=20\text{Hz}$ ). The  
136 interaction strength  $\varepsilon$  was found to be inversely correlated to the cortical distance between probes  
137 (M1:  $R^2=0.41$ , M2:  $R^2=0.29$ , both  $p < 10^{-10}$ , Fig.S5E) in line with anatomy of horizontal  
138 connectivity<sup>27</sup>. Quantitative predictions were derived of V1 gamma synchronization for different  
139  $\Delta\omega$  and  $\varepsilon$  values. We found that gamma PLV closely followed the analytical predictions as a  
140 function of  $\Delta\omega$  and  $\varepsilon$  (model accuracy for population averages: M1:  $R^2=0.88$ , M2:  $R^2=0.90$ ; for  
141 single contact data: M1:  $R^2=0.18$ , M2:  $R^2=0.32$ , see in Fig.3B/G). The PLV was dependent on  
142 both interaction strength  $\varepsilon$  and detuning  $\Delta\omega$  (Fig.S7), and showed the predicted Arnold tongue in  
143 both M1 and M2 (Fig.3C/H). The mean phase differences (dots in Fig.3D/I) were also well  
144 predicted by the analytical model as a function of  $\Delta\omega$  and  $\varepsilon$  (gray line; model accuracy for  
145 population averages: M1:  $R^2=0.94$ , M2:  $R^2=0.88$ , for single-contact data: M1:  $R^2=0.56$ , M2:  
146  $R^2=0.27$ ). The phase difference was largely determined by detuning  $\Delta\omega$  and more weakly by  
147 interaction strength  $\varepsilon$  (Fig.S7). The phase spread (Fig.3E/J) had a range of nearly  $-\pi/2$  to  $\pi/2$  in  
148 both M1 and M2 as predicted by the shape of  $G(\theta)$  (Fig.S1). We confirmed the phase locking  
149 value and phase difference analysis in spike-CSD and spike-spike measurements (Fig.S8).

150 The present study shows that gamma synchronization in PING networks and in awake monkey  
151 V1 adheres to theoretical principles of weakly coupled oscillators<sup>20–25</sup>, thereby providing insight  
152 into the dynamic principles underlying neural synchronization. Crucially, we observed phase-  
153 dependent modulations of the instantaneous frequency difference in both PING model and V1  
154 recording data. These modulations are characteristic for the intermittent synchronization  
155 regime<sup>24,29–31</sup> which naturally arises in frequency-variable and noisy oscillator networks.. These  
156 observations show that a fixed and common frequency is not per se required for  
157 synchronization<sup>10,14</sup>. To the contrary, these non-stationary frequency modulations reflect the  
158 essential process of synchronization and, furthermore, allow the experimental estimation of the  
159 interaction function<sup>20–25</sup> and of the regulative parameters underlying gamma synchronization. We  
160 found that two parameters mainly regulated synchronization: the detuning  $\Delta\omega$  (mean frequency  
161 difference) and the interaction strength  $\varepsilon$  (amplitude of frequency modulations). This was  
162 highlighted in the mapping of the Arnold tongue<sup>21,24,29</sup>, a predicted synchronization region within  
163 the parameter space of detuning and interaction strength. In our experiment, detuning was  
164 dependent on the local contrast difference<sup>14,16</sup>, known to change neural excitation in V1<sup>33</sup>, while  
165 the interaction strength was dependent on the underlying horizontal connectivity strength, here  
166 varied by cortical distance<sup>27</sup>. These properties suggest V1 gamma as a relevant mechanism for  
167 sensory processing<sup>6,7</sup> as local gamma synchronization will be informative about the sensory  
168 input<sup>8</sup> and informative about the underlying structure of connectivity. Gamma frequency is  
169 indeed modulated by various sensory stimuli<sup>2,14,16,15,34</sup> and by cognitive manipulations<sup>26,2,10</sup>.  
170 Importantly, in line with previous findings<sup>14,35–37</sup>, we found V1 gamma synchrony to be local,  
171 restricted by horizontal connectivity that extends only few mm across the cortex<sup>27</sup>, and hence not  
172 likely to reflect whole perceptual objects. Our findings reconcile several studies that have given  
173 different theoretical interpretations to observations of frequency variation in gamma<sup>14,16,12</sup>, and  
174 are relevant for understanding synchronization with and across cortical areas where differences  
175 in preferred frequency were also observed<sup>9,10</sup>. Given the generality of the observed  
176 synchronization principles, they are likely to apply to other brain regions and frequency bands.

177

## 178 **Methods summary:**

179 Experimental recording: We recorded in two adult male rhesus monkeys, implanted with a  
180 chamber above early visual cortex, positioned over V1/V2. A head post was implanted to head-  
181 fix the monkeys during the experiment. The monkey's task was to passively gaze on a fixation  
182 point while a whole-field static square-wave grating was shown. We simultaneously recorded  
183 from multiple locations in monkey V1 while the monkey viewed luminance gratings in which  
184 spatially varying contrast set the frequencies of local gamma rhythms<sup>14,16</sup>. V1 recordings were  
185 made with 16-contact laminar Plexon U-probes (Plexon Inc.). We recorded the local field  
186 potential (LFP) and multi-unit spiking activity (MUA). We aligned the neural data from the  
187 different laminar probes according to their cortical depth and excluded contacts coming from



188 deep V2. All the procedures were in accordance with the European council directive  
189 2010/63/EU, the Dutch ‘experiments on animal acts’ (1997) and approved by the Radboud  
190 University ethical committee on experiments with animals (Dier-Experimenten-Commissie,  
191 DEC).

192 Theoretical and computational modelling: Using the formulism of the theory of weakly coupled  
193 oscillators we investigated the phase-locking as well as the mean phase difference of two  
194 mutually coupled noisy phase-oscillators with variable intrinsic frequency difference (detuning)  
195 and interaction strength. The stochastic differential equation was solved analytically<sup>24</sup>. The  
196 analytical results correctly predicted the numerical simulations. In addition, we simulated two  
197 coupled excitatory-inhibitory spiking networks generating gamma oscillations using Izhikevich-  
198 type neuronal model. The detuning between the networks was altered by changing the difference  
199 in excitatory input drive. The interaction strength was altered by changing the cross-network  
200 synaptic connection strength.

201 Signal analysis and statistics: To investigate dynamical changes in the gamma phase and  
202 frequency over time we estimated the instantaneous gamma phase and frequency using the  
203 singular spectrum decomposition of the signal (SSD, see  
204 <https://project.dke.maastrichtuniversity.nl/ssd/>). The accuracy of the theoretical predictions for  
205 the experimental data was quantified as the explained variance  $R^2$ . We also used a multiple  
206 regression approach to quantify the dependence of synchronization on amplitude, detuning and  
207 interaction strength. Full methods and additional supportive results are provided in the  
208 *Supplementary Information*.

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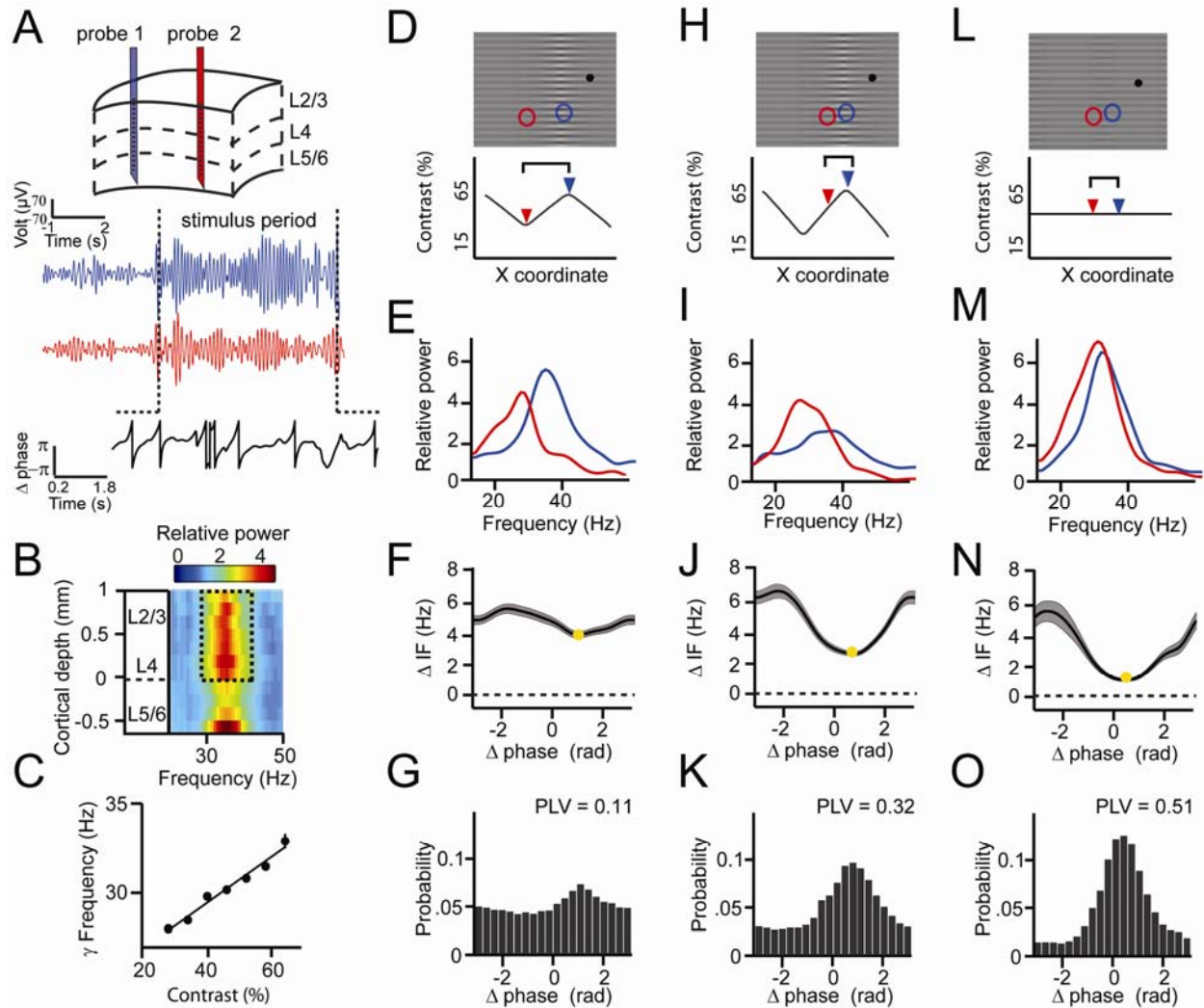
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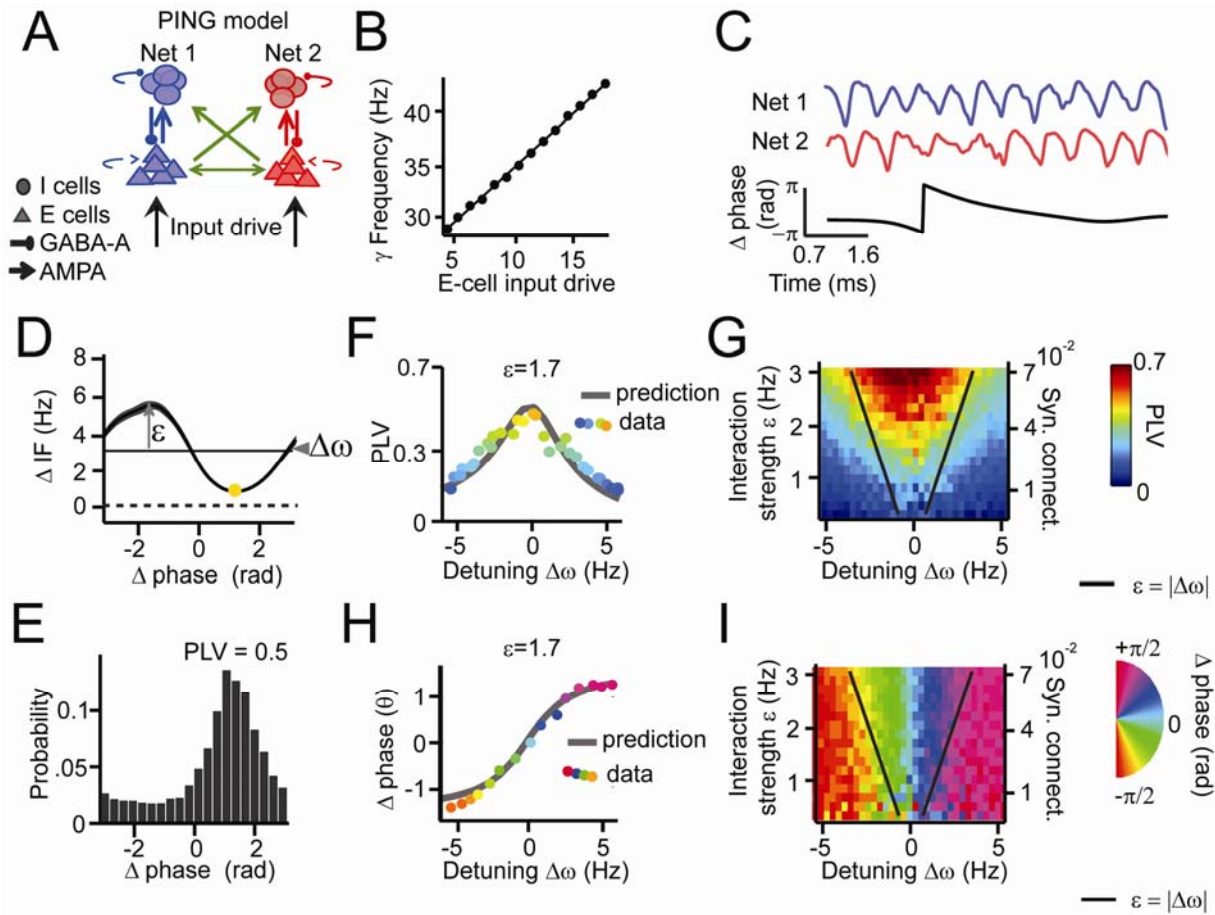


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300 **Fig.1.** Experimental paradigm and intermittent synchronization. (A) Recordings preparation and  
 301 example CSD (blue and red) traces from which phase difference (black) trace was extracted. The  
 302 gradient of the black trace indicates the rate of phase precession. (B) Spectral power relative to  
 303 baseline as a function of V1 cortical depth (36.5% contrast, population average, M1) dashed box  
 304 indicates gamma in the layers taken for main analysis (C) Local contrast modulated gamma  
 305 frequency (population average, M1). (D-G) Example 1 showing synchronization despite  
 306 frequency difference. (D) Stimulus grating and fixation spot (black dot). Two receptive fields  
 307 (RF) from different probes are superimposed (blue and red circles). Below, black line gives  
 308 contrast over space, arrowheads mark RF positions. (E) Power spectra of the two probes showing  
 309 different peak frequencies. (F) Instantaneous frequency difference ( $\Delta IF$ ), equivalent to phase  
 310 precession rate, as a function of phase difference. Yellow dot indicates the modulation minimum  
 311 equivalent to preferred phase difference, shading is  $\pm SE$  (G) The phase difference probability  
 312 distribution and phase-locking value (PLV). (H-K) Example 2; probes were closer, gamma peak  
 313 frequency difference was similar. Conventions as in D-G. (L-O) Example 3; same distance,

314 reduced frequency difference. Compare F, J, N RF distance determine IF modulation amplitude,  
 315 gamma frequency difference determines preferred phase difference.

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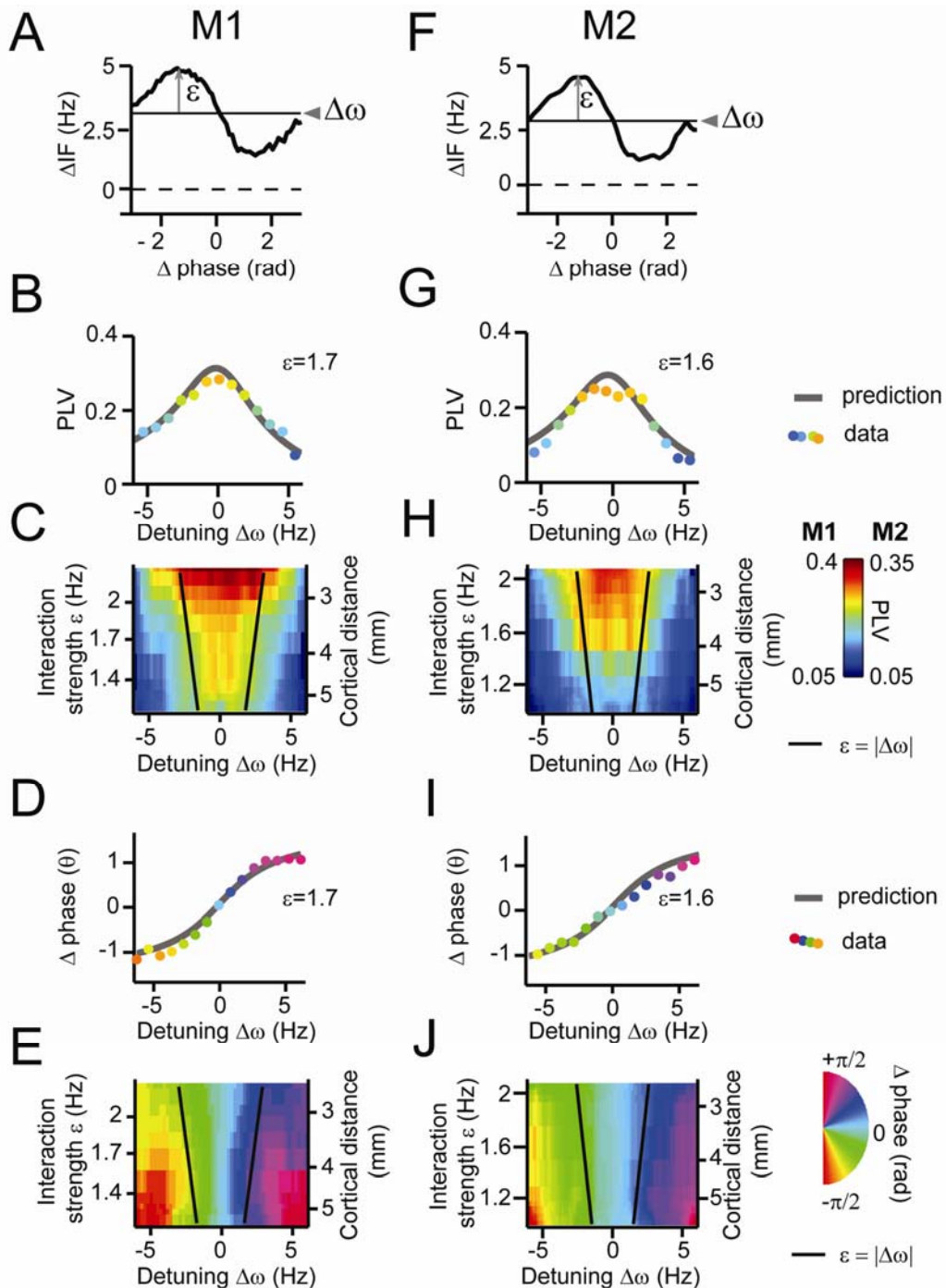


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321 **Fig.2.** Applying the theory of weakly coupled oscillators to coupled PING networks. (A) Two  
 322 coupled pyramidal-interneuron gamma (PING) networks (Net 1 and Net 2). (B) The frequency of  
 323 gamma in a single network depends on input strength. (C) Simulation output example network  
 324 signals (red and blue) and phase difference  $\theta$  (black). (D) An example  $\Delta \overline{IF}(\theta)$  modulation used  
 325 to estimate the interaction strength  $\varepsilon$  and detuning value  $\Delta\omega$ . The shape of the modulation  
 326 indicates the  $G(\theta)$ . (E) The corresponding phase difference probability distribution. (F) The  
 327 simulation PLV at different detuning values  $\Delta\omega$  (dots colored by PLV) at a single interaction  
 328 strength value ( $\varepsilon = 1.7$ ) was well predicted by the model (gray line). (G) The PLV at many  
 329 interaction strengths and detuning values mapped the Arnold tongue. Black lines mark the  
 330 predicted Arnold tongue borders in the noise-free case ( $\varepsilon = |\Delta\omega|$ ). (H-I) As (F-G), but for preferred  
 331 phase difference  $\theta$ . Color code of dots in F and H as in G and I, respectively.

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335 **Fig.3.** Predicting V1 gamma synchronization. (A-E) Results from M1. (A) An example  $\overline{\Delta IF}(\theta)$   
 336 modulation used for estimating  $\epsilon$  and  $\Delta\omega$  for monkey M1 (left) and M2 (right). The shape of the  
 337 modulation indicates the  $G(\theta)$ . (B) Observed PLV (dots) and analytical prediction (gray line) as a

338 function of detuning  $\Delta\omega$  for one level of interaction strength ( $\varepsilon=1.7$ ). **(C)** Combining different  
339 detuning  $\Delta\omega$  and interaction strengths  $\varepsilon$  we observed a triangular synchronization region, the  
340 Arnold tongue. Black lines mark the predicted Arnold tongue border as expected from the noise-  
341 free case ( $\varepsilon=|\Delta\omega|$ ) **(D)** Analytical prediction (gray) and experimentally observed preferred phase  
342 differences (dots colored by phase difference) as a function of detuning  $\Delta\omega$  for one level of  
343 interaction strength ( $\varepsilon=1.7$ ). **(E)** Similar to C), but now plotting the preferred phase difference.  
344 **(F-J)** As (A-E) but for M2 population data. Color coding of dots in B, G, D, I is as indicated in  
345 color scales in panels just below them.  
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