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Fitness costs of noise in biochemical reaction networks and the evolutionary limits of cellular robustness

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Keywords: gene expression noise; drift barrier hypothesis; robustness; noise propagation; allometric scaling; coarse-grained cell model; linear noise approximation; Michaelis-Menten reaction; non-linear stochastic model

28 **ABSTRACT**

29 Gene expression is inherently noisy, but little is known about whether noise affects cell  
30 function or, if so, how and by how much. Here I present a theoretical framework to  
31 quantify the fitness costs of gene expression noise and identify the evolutionary and  
32 synthetic targets of noise control. I find that gene expression noise reduces fitness by  
33 slowing the average rate of nutrient uptake and protein synthesis. This is a direct  
34 consequence of the hyperbolic (Michaelis-Menten) kinetics of most biological reactions,  
35 which I show cause “hyperbolic filtering”, a process that diminishes both the average rate  
36 and noise propagation of stochastic reactions. Interestingly, I find that transcriptional  
37 noise directly slows growth by slowing the average translation rate. Perhaps surprisingly,  
38 this is the largest fitness cost of transcriptional noise since translation strongly filters  
39 mRNA noise, making protein noise largely independent of transcriptional noise,  
40 consistent with empirical data. Translation, not transcription, then, is the primary target of  
41 protein noise control. Paradoxically, selection for protein-noise control favors increased  
42 ribosome-mRNA binding affinity, even though this increases translational bursting.  
43 However, I find that the efficacy of selection to suppress noise decays faster than linearly  
44 with increasing cell size. This predicts a stark, cell-size-mediated taxonomic divide in  
45 selection pressures for noise control: small unicellular species, including most  
46 prokaryotes, face fairly strong selection to suppress gene expression noise, whereas larger  
47 unicells, including most eukaryotes, experience extremely weak selection. I suggest that  
48 this taxonomic discrepancy in selection efficacy contributed to the evolution of greater  
49 gene-regulatory complexity in eukaryotes.

50

51 **ARTICLE SUMMARY**

52 Gene expression is a probabilistic process, resulting in random variation in mRNA and  
53 protein abundance among cells called “noise”. Understanding how noise affects cell  
54 function is a major problem in biology. Here I present theory demonstrating that gene  
55 expression noise slows the average rate of cell division. Furthermore, by modeling  
56 stochastic gene expression with non-linearity, I identify novel mechanisms of cellular  
57 robustness. However, I find that the cost of noise, and therefore the strength of selection  
58 favoring robustness, decays faster than linearly with increasing cell size. This may help  
59 explain the vast differences in gene-regulatory complexity between prokaryotes and  
60 eukaryotes.

61

62 **INTRODUCTION**

63 The probabilistic nature of chemical reactions, along with the small volumes of living  
64 cells, create substantial molecular randomness called “noise” (Berg 1978; McAdams and  
65 Arkin 1997). Gene expression noise, for example, is measured as variation in mRNA or  
66 protein abundance within a single cell over time (Taniguchi et al. 2010), or, equivalently,  
67 among cells in an isogenic cell population (Elowitz et al. 2002). Because the abundance  
68 of most mRNAs and proteins are on the order of 1 to 1000 copies per cell (Taniguchi et  
69 al. 2010), a random birth or death of a single copy can instantaneously change the cellular  
70 concentration by 0.1-100%. These random fluctuations are amplified by bursting, the  
71 synthesis of multiple mRNAs or proteins in brief pulses (Raser and O’Shea 2004; Raser  
72 and O’Shea 2005; Raj et al. 2006; Pedraza and Paulsson 2008; Thattai and van  
73 Oudenaarden 2001). Using fluorescent probes and fluorescently tagged proteins,

74 researchers have quantified the noise statistics of mRNA and/or proteins in *E. coli*,  
75 budding yeast, and mammalian cell lines, demonstrating the ubiquity of gene expression  
76 noise and its considerable variability across the genome and among taxa (Elowitz et al.  
77 2002; Blake et al. 2003; Golding et al. 2005; Raj et al. 2006; Zenklusen, Larson, and  
78 Singer 2008; Taniguchi et al. 2010; Newman et al. 2006; Balázsi, van Oudenaarden, and  
79 Collins 2011).

80 A central question in cell biology and evolution is whether or not cellular noise  
81 has a functional consequence. Answering this question is necessary for establishing  
82 rational design criteria for synthetic genetic circuits and for providing evolutionary  
83 explanations for empirical patterns of noise among genes and taxa. For example, why  
84 does *E. coli* have an average mRNA Fano factor (variance/mean) of  $F = 1.6$  (Taniguchi et  
85 al. 2010), which is just above the theoretical minimum of  $F = 1$ , whereas mammalian  
86 cells have 40-100 times stronger mRNA noise (Raj et al. 2006; Balázsi,  
87 van Oudenaarden, and Collins 2011; Dar et al. 2012)? While the mechanistic causes of  
88 this discrepancy are becoming clearer (Suter et al. 2011), the evolutionary forces  
89 permitting such a discrepancy remain murky.

90 Mounting empirical evidence suggests that natural selection has acted to limit  
91 gene expression noise in many cases (Metzger et al. 2015; Newman et al. 2006; Alemu et  
92 al. 2014; Batada and Hurst 2007; Fraser et al. 2004; Lehner 2008), implying that noise is  
93 often costly to cells. Yet identifying and quantifying this cost remains a major challenge.  
94 Perhaps the most widely invoked cost of cellular noise is imprecision, the lack of fine  
95 control over cellular processes. Evolutionarily, this cost of noise is typically modeled as  
96 random phenotypic deviations from the adaptive optimum, which is opposed by

97 stabilizing selection (Wagner et al. 1997; Draghi and Whitlock 2015). Several molecular  
98 mechanisms linking imprecision to suboptimal fitness deviations have been proposed.  
99 The dosage-control hypothesis proposes that noise disrupts control over the  
100 stoichiometric balance of interacting proteins (Fraser et al. 2004; Lehner 2008).  
101 Consistent with this hypothesis, proteins that form complexes have lower average noise  
102 than other proteins (Fraser et al. 2004; Lehner 2008; Bar-Even et al. 2006), but models  
103 currently do not predict the magnitude of this effect so that quantitative correspondence  
104 between predictions and observations remain unknown. More commonly, researchers  
105 focus on how noise disrupts the fine control over cell states regulated by bistable  
106 (ON/OFF) switches (Balázsi, van Oudenaarden, and Collins 2011; Csete and Doyle  
107 2002). Noise in the input parameter can cause the cell to flip stochastically from the  
108 desired to undesired state, or to respond slowly to environmental cues, creating a  
109 maladaptive mismatch between the cell state and the selecting environment. However,  
110 the least noisy genes, those presumably experiencing the strongest selection to suppress  
111 noise, are constitutively expressed and involved in core cell functions, not control over  
112 bistability (Fraser et al. 2004; Lehner 2008; Bar-Even et al. 2006). This cost is also  
113 difficult to quantify in any general way because its effects are highly context dependent,  
114 to the extent that in some contexts the cost of imprecision can change signs, becoming  
115 evolutionarily beneficial (Acar, Mettetal, and van Oudenaarden 2008; Rao, Wolf, and  
116 Arkin 2002; Kussell and Leibler 2005; Eldar and Elowitz 2010; Kaern et al. 2005; Thattai  
117 and Van Oudenaarden 2004).

118 As an alternative, here I demonstrate a general, readily quantifiable and consistent  
119 cost of cellular noise. I show that noise in substrate abundance slows the average rate of

120 product formation in Michaelis-Menten-type (i.e., non-cooperative bimolecular) chemical  
121 reactions, which constitute most reactions in the cell including those most closely tied to  
122 fitness: nutrient uptake and protein synthesis. This result is a straightforward consequence  
123 of the non-linearity of bimolecular reactions. While it is well known that noise alters the  
124 mean of nonlinear stochastic systems, this effect is typically ignored, giving rise to the  
125 widespread use of linear approximations (e.g., the “linear noise approximation” (LNA)  
126 (van Kampen 2007)), particularly in studies of gene expression noise (Thattai and van  
127 Oudenaarden 2001; Shahrezaei and Swain 2008; Paulsson 2004). The LNA is well-  
128 justified when applied to investigating steady-state concentrations in noisy chemical  
129 networks, since the difference between linear and non-linear predictions at steady-state  
130 are on the order of a single molecule. However, fitness in living cells is determined by  
131 rates—not concentrations--integrated over an enormous number of reactions over the  
132 lifetime of a cell and its genetic lineage. For example, an *E. coli* cell must carry out a  
133 minimum of  $2.5 \times 10^6$  translation reactions to produce a single daughter cell. A minor  
134 slowing of reaction rates within cells, if heritable, may pose a non-negligible selective  
135 cost. From this perspective, the effects of noise on fitness cannot be ignored, and indeed,  
136 I find that they can be substantial in small cells.

137

## 138 **RESULTS**

139 The goal of the paper is to quantify the strength of selection favoring the attenuation of  
140 gene expression noise, and to identify the targets of noise attenuation. The paper proceeds  
141 as follows. First, I show that noise slows the average rate of biochemical reactions that  
142 have non-linear, hyperbolic kinetics. Then, I extend this result to networks of coupled

143 non-linear reactions in order to determine how noise effects total flux through  
144 biochemical networks, as well as how the network architecture itself affects the  
145 propagation of noise. This enables us to quantify the cost of noise, measured as the loss in  
146 average rate of end-product formation, and to quantify the noise statistics of the system. I  
147 then apply the model to gene expression, focusing on the consequences of non-linear,  
148 hyperbolic translation kinetics on the propagation of mRNA noise to the protein level,  
149 and to the average rate of biomass synthesis. Finally, I investigate how the volume of  
150 reaction compartments and whole-cells influences the cost of noise, concluding that the  
151 cost of noise, and therefore the selection pressures favoring low gene expression noise,  
152 decays rapidly with increasing cell size, leading to a testable prediction about levels of  
153 gene expression noise across unicellular species.

154

155 **Noise-induced slowdown of Michaelis-Menten reactions.** The purpose of this section is  
156 to begin building the theory from the simplest component of a network: a single isolated  
157 reaction. We will see that input noise slows the reaction rate of a single, isolated reaction  
158 obeying Michaelis-Menten type kinetics. In the following section, the analytical results  
159 from this section are put on a more rigorous theoretical footing and then extended to  
160 networks of coupled reactions in open systems.

161 The cornerstone of biochemistry is the Michaelis-Menten (MM) reaction, which  
162 describes the enzyme catalyzed conversion of substrate,  $S$ , into product,  $P$ . The rate of  
163 product formation is described by the hyperbolic MM equation (Michaelis and Menten  
164 1913; Briggs and Haldane 1925; Ingalls 2013):

$$165 \quad \frac{dp}{dt} = V = V_{max} \frac{s}{K_m + s} \quad 1$$

166  $s$ ,  $e$  and  $p$  are reactant, enzyme and product concentration, respectively (following  
167 convention lowercase denotes concentration, which unless otherwise noted will take units  
168 of molecules per cell),  $V_{max} = k_{cat}e$  is the asymptotic reaction rate ( $k_{cat}$  is the number of  
169 substrate molecules converted to product per enzyme per unit time), and  $K_m$ , the  
170 Michaelis or half-saturation constant, is the substrate concentration at  $V = V_{max}/2$ . This  
171 equation is derived from a system of elementary reactions by applying the law of mass  
172 action and the quasi-steady state approximation (QSSA) (Briggs and Haldane 1925;  
173 Ingalls 2013) (Sup. Mat.). The QSSA is a separation of timescales approximation that  
174 eliminates fast variables by setting them to their steady-state. The law of mass action is  
175 valid in cases where the timescale of molecular diffusion is fast relative to the reaction  
176 rate, so that the system behaves as if it were well-mixed. The rate of enzymatic reactions  
177 in cells is far below the diffusion limit (Bar-Even et al 2006), so that mass action is a  
178 reasonable approximation in most cases. For stochastic reactions, wherein  $s$ ,  $e$ , and  $p$   
179 experience random number fluctuations, Eqn. 1 takes on a probabilistic interpretation as  
180 the “reaction propensity” (van Kampen 2007) giving the transition probability for  $n_p \rightarrow$   
181  $n_p + 1$  and  $n_s \rightarrow n_s - 1$ , where  $n$  denotes particle number. In general, the QSSA is valid  
182 for stochastic systems under the same conditions for which it is valid in deterministic  
183 systems (Rao and Arkin 2003; Sanft, Gillespie, and Petzold 2011), notwithstanding  
184 finite-volume corrections (Grima 2009b; Grima 2009a; Grima 2010) (see below).

185       Because Eqn. 1 is hyperbolic in  $s$ , Jensen’s inequality proves that, for random  $s$ ,  
186  $\langle V(s) \rangle \leq V(\langle s \rangle)$  (angle brackets denote ensemble average). That is, the average reaction  
187 rate is less than or equal to a noise-free (i.e., deterministic) reaction with  $s = \langle s \rangle$  (Figs. 1,  
188 S1). This is seen by expanding Eqn. 1 to second order in Taylor series about  $\langle s \rangle$  and

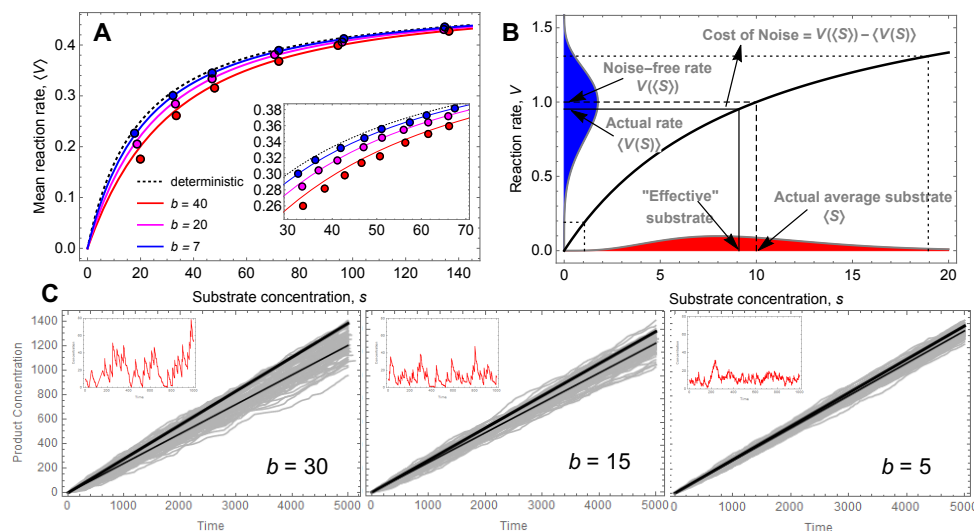


189 taking the expectation (see SI),

$$190 \quad \langle V(s) \rangle = \underbrace{V_{max} \frac{\langle s \rangle}{K_m + \langle s \rangle}}_{v(\langle s \rangle)} - \underbrace{\frac{\langle \epsilon_s^2 \rangle}{\langle s \rangle^2} V_{max} a(1-a)^2}_{\text{noise-induced slowdown}} + O(\epsilon_s^3) \quad 2$$

191  $a = K_m/(K_m + \langle s \rangle)$  is the “kinetic order” of the MM reaction evaluated at the average  
192 substrate concentration, and ( $\eta^2 = \langle \epsilon_s^2 \rangle / \langle s \rangle^2$ ) the squared coefficient of variation, which  
193 is often simply called “the noise”. The leading order correction in Eqn. 2 takes the sign of  
194 the second derivative of Eqn. 1, which is everywhere negative for hyperbolic transition  
195 rates, thus demonstrating the noise-induced slowdown of MM type reactions.

196 The kinetic order parameter,  $a$ , is a key parameter that appears repeatedly below.  
197  $a$  is identical to the steady-state fraction of unbound enzyme. In metabolic control  
198 analysis (MCA),  $a$  is equivalent to the “elasticity coefficient” (Kacser and Burns 1973).  
199 In general,  $a$  measures the nonlinearity of the reaction rate at a given mean substrate  
200 concentration. First and zeroth order reactions ( $a = 1$  and  $0$ , respectively) are both linear,  
201 and as a consequence there is no noise-induced slowdown, consistent with Jensen’s  
202 Inequality. The rate of a first order reaction is linearly dependent on the substrate  
203 concentration, whereas the rate of a zeroth order reaction is independent of substrate  
204 concentration. In a MM type reaction, the reaction is in its first order regime when  
205 substrate is scarce, and in the zeroth order regime when substrate is abundant (the flat  
206 part of the rate law curve in Fig. 1).



207

208 **Figure 1: Noise-induced slowdown of hyperbolic reactions.** Stochastic  
 209 simulations employing Gillespie's (Gillespie 1977) exact stochastic simulation algorithm  
 210 (SSA) of the microscopic system of elementary reactions for the Michaelis-Menten  
 211 reaction (see SI). Points represent the rate of product formation over time steps 100-800  
 212 (providing a burn-in more than 10 times the half life of  $P$  or  $S$ ) ensemble averaged over  
 213 1000 realizations of the stochastic process. Results fit well with the QSSA model (Eqn. 3,  
 214 solid lines in **a**, **b** and thin black line in **c**). The substrate burst input,  $b$ , is varied to alter  
 215 the substrate noise level. Substrate input rate,  $k_m$  is divided by  $b$  so that mean substrate is  
 216 not affected by bursting, and is varied to generate different substrate concentrations. All  
 217 other parameters are constant ( $e = 1$ ,  $k_{ON} = 1$ ,  $k_{OFF} = 19.5$ ,  $k_{cat} = 0.5$ ,  $\delta_S = 0.1$ ,  $\delta_P = 0.1$ ). **A**)  
 218 The rate law curve demonstrating that, for a given average substrate concentration,  
 219 substrate noise slows the reaction, but only if the reaction is not in its first or zeroth order  
 220 regime (low or high substrate, respectively). Note that the accuracy of the stochastic  
 221 QSSA requires that  $s, K_m \ll F$ , which fails for the red points in **a** and **b** (where  $K_m = 20$   
 222 and  $F = 20.5$ ). Inset is a zoomed in view. **B**) A geometric illustration of noise-induced  
 223 slowdown, showing that variance (noise) in a state variable reduces the average of a  
 224 hyperbolic function (i.e., Jensen's inequality). **C**) 100 realizations of the stochastic  
 225 process (gray) compared to the deterministic (thick black line) and stochastic (thin black  
 226 line) predictions. Red traces in inset are example outputs from a single realization of the  
 227 substrate dynamics for a short time window.

228

229 A more rigorous and general expansion method is used in the following section

230 ("Noisy Non-linear Reaction Networks in Finite Volumes"). For now, there are two main

231 points: 1) reaction kinetics depend not only on average molecule concentrations, but on

232 the mean and variance in molecule numbers, and 2) for a given average substrate

233 concentration, substrate noise reduces the rate of biochemical reactions. The difference

234 between the average stochastic rate and deterministic rate grows monotonically with  
235 increasing noise, and equality occurs only with zero noise or zero curvature of  $V(s)$ .  
236 Mechanistically, this can be understood as follows: positive fluctuations in substrate  
237 oversaturate available catalysts leaving many unbound  $S$  molecules that are not actively  
238 creating product, causing inefficiency. As a consequence, the same amount of substrate  
239 would produce a faster net reaction rate if spread out evenly over the same time interval  
240 (i.e., the same mean with zero noise).

241 If the substrate statistics are known, then the probability density of  $V$  with random  
242  $s$  can be solved explicitly for a number of common probability distributions. mRNA and  
243 protein counts are typically gamma distributed within a cell over time or among cells in a  
244 clonal population (Taniguchi et al. 2010). The gamma distribution has shape parameter  $\alpha$   
245  $= CV^{-2}$  and rate parameter  $\beta = F$ , where  $CV^2 = \sigma_x^2 / \langle x \rangle^2$  and  $F = \sigma_x^2 / \langle x \rangle$  are the squared  
246 coefficient of variation (the noise) and the Fano factor (often called “the noise strength”),  
247 respectively (Taniguchi et al. 2010).  $F$  and  $CV^2$  are common measures of noise intensity,  
248 with convenient mechanistic interpretations in gene expression:  $F$  is closely related to the  
249 burst size (e.g., number of proteins synthesized per mRNA lifetime) and  $CV^{-2}$  to the burst  
250 frequency (e.g., the number of mRNA’s synthesized per protein lifetime).  $F$  is  
251 independent of cell volume, whereas  $CV^2$  decreases with cell volume (they are intensive  
252 and extensive variables, respectively). For gamma distributed  $s$ , the probability density  
253 function of the Michaelis-Menten reaction rate (for  $V < V_{max}$ ) is exactly,

$$254 \quad Pr(V = v) = \frac{V_{max}}{k_m} \frac{v^{\alpha-1} e^{-v/z\beta}}{z(z\beta)^\alpha \Gamma(\alpha)}, \quad z = \frac{V_{max} - v}{k_m} \quad 3$$

255 Eqn. 3 is verified by Monte Carlo simulations implemented with Gillespie’s exact  
256 stochastic simulation algorithm (SSA) (Gillespie 1977) (Fig. 1). Eqn. 3 fits data from

257 stochastic simulations very well provided that  $F, K_m \ll s$ . Equation 2 gives a very good  
258 approximation to 3 under this same restriction (Fig. S2), which provides an important  
259 benchmark for the accuracy of the more rigorous expansion approximations used below.

260 Figure 1 shows that the effect of noise on MM-type reactions is to increase the  
261 observed Michaelis constant. Therefore, a stochastic formulation of Eqn. 1 that accounts  
262 for noisy substrate while retaining the classical form is obtained simply by substituting  
263 into Eqn. 1 an “effective”, stochastic Michaelis constant,  $\tilde{K}_m$ , where  $\tilde{K}_m \geq K_m$  with  
264 equality occurring only with zero noise, zero curvature of the rate law, and macroscopic  
265 reaction volume (S6).

266

267 **Noisy Non-linear Reaction Networks in Finite Volumes.** In the previous section, we  
268 showed that substrate noise slows the rate of product formation in a single hyperbolic  
269 reaction in a closed system. To apply this result, the statistics of the substrate species are  
270 required. We now take a more rigorous approach and generalize the noise-induced  
271 slowdown phenomenon to reaction networks in open systems. Further, the model allows  
272 one to predict the noise statistics of the system solely with knowledge of the network  
273 structure and reaction kinetics. Of particular interest is how network architectures may  
274 create feedbacks that buffer the system flux against noise-induced slowdown. The  
275 following theory builds on the work of van Kampen (van Kampen 2007) and Elf and  
276 Erhenberg (Elf and Ehrenberg 2003) (see also Grima (Grima 2010)).

277 The time evolution of the joint probability density of all  $N$  nodes in a chemical  
278 reaction network with  $R$  edges is given by the following “system-size expansion” (van  
279 Kampen 2007) of the multidimensional chemical master equation (derived in the SI):

$$\begin{aligned}
 & \text{linear Fokker-Planck equation} \\
 280 \quad \frac{\partial \Pi(\vec{\epsilon}, t)}{\partial t} &= - \sum_{i,k}^N J_{ik} \frac{\partial(\epsilon_k \Pi)}{\partial \epsilon_i} + \frac{1}{2} \sum_{i,k=1}^N D_{ik} \frac{\partial^2 \Pi}{\partial \epsilon_i \partial \epsilon_k} \\
 281 \quad &+ \underbrace{\frac{\Omega^{-1/2}}{2} \left( - \sum_{i,k,l=1}^N \frac{\partial J_{ik}}{\partial \bar{x}_l} \frac{\partial(\epsilon_l \epsilon_k \Pi)}{\partial \epsilon_i} + \sum_{i,k,l=1}^N \frac{\partial D_{ik}}{\partial \bar{x}_l} \frac{\partial^2(\epsilon_l \Pi)}{\partial \epsilon_i \partial \epsilon_k} \right)}_{\text{correction for non-linearity and finite volume}} + O(\Omega^{-1}) \quad 4
 \end{aligned}$$

282  $\Pi = \Pi(\vec{\epsilon}, t)$  is the joint probability density of all reacting species,  $\Omega$  is the reaction  
 283 volume,  $\epsilon_i$  is the random perturbation of the  $i^{\text{th}}$  species from  $\bar{x}_i$ , its macroscopic steady-  
 284 state (overbars will denote the macroscopic steady-state value of a variable throughout),  
 285  $S_{ij}$  is the stoichiometric coefficient of species  $i$  in the  $j^{\text{th}}$  reaction, and

$$286 \quad J_{ik} = \sum_{j=1}^R S_{ij} \frac{\partial f_j(\vec{x})}{\partial \bar{x}_k}, \quad D_{ik} = \sum_{j=1}^R S_{ij} S_{kj} f_j(\vec{x})$$

287 are entries of the Jacobian and diffusion matrices, respectively, of the continuous  
 288 deterministic system of equations. Equation 4 is a linearization of the chemical master  
 289 equation obtained by applying the ansatz that concentrations fluctuate about their  
 290 macroscopic steady-state value, with fluctuation size scaling inversely with the square  
 291 root of the reaction volume,  $\Omega$  (van Kampen 2007). This inverse square-root dependence  
 292 of noise on volume is a formal statement of the Law of Large Numbers, and is formally  
 293 equivalent to the relationship between sample size and the standard deviation of a  
 294 sampling distribution. One thus substitutes,  $x_i = \bar{x}_i + \Omega^{-1/2} \epsilon_i$ , for each species  $i$  in the  
 295 deterministic system of equations, and then expands each transition rate,  $V_j(x)$ , in powers  
 296 of  $\Omega^{-1/2}$  about  $\bar{x}_i$ , truncating the series at the desired order (van Kampen 2007). The terms  
 297 of order  $O(\Omega^0)$  in (4) give the “linear noise approximation” (LNA) (van Kampen 2007),  
 298 which is equivalent to the standard linear Fokker-Planck equation, while the terms in  $\Omega^{-1}$

299 <sup>1/2</sup> account for noise-induced deviations caused by non-linear rate laws in finite volumes,  
 300 which we show cannot be ignored in living cells.

301 To apply 4, one begins by writing a macroscopic description of a reaction network  
 302 in QSSA reduced form. Here we analyze a generalized MM scheme (S17) in an open  
 303 system obeying the following macroscopic system of ODEs (Sup. Mat.):

$$304 \quad \frac{ds}{dt} = b_S k_{in,S} - hV - \delta_S s, \quad \frac{de}{dt} = b_E k_{in,E} - \delta_E e, \quad \frac{dp}{dt} = V - \delta_P p \quad 5$$

305  $\delta_i$  and  $k_{in,i}$  are the decay and input rates, respectively, of  $i$ , and  $V = k_{cat} e s / (K_m + s)$ .  
 306 Substrate ( $S$ ) and enzyme ( $E$ ) are fed into the system via a burst process with burst size  $b_S$   
 307 and  $b_E$ , respectively ( $n_S \rightarrow n_S + b_S$  ;  $n_E \rightarrow n_E + b_E$  ).  $h = \{0,1\}$  is an indicator variable  
 308 that toggles between reactions where  $S$  is consumed ( $h = 1$ ), as in most metabolic  
 309 reactions, and reactions where  $S$  is not consumed ( $h = 0$ ), as with template-mediated  
 310 reactions such as translation where  $E$  denotes ribosomes and  $S$  mRNA, which is not  
 311 consumed in translation.

312 From System 5 and Eqn 4, we obtain the following steady-state average  
 313 concentrations of each species in a generalized, open MM reaction network (Sup. Mat.):

$$314 \quad \langle s \rangle = \bar{s} + h \underbrace{\left[ a \frac{\langle \epsilon_S^2 \rangle}{\bar{s}} - K_m \frac{\langle \epsilon_S \epsilon_E \rangle}{\bar{e} \bar{s}} \right]}_{\text{stochastic amplification}} A \Omega^{-1}, \quad \langle e \rangle = \bar{e},$$

$$315 \quad \langle p \rangle = \bar{p} - \frac{\delta_S}{\delta_P} \left[ a \frac{\langle \epsilon_S^2 \rangle}{\bar{s}} - K_m \frac{\langle \epsilon_S \epsilon_E \rangle}{\bar{e} \bar{s}} \right] A \Omega^{-1}, \quad 6$$

$$316 \quad \langle V \rangle = \bar{V} - \delta_S \underbrace{\left[ a \frac{\langle \epsilon_S^2 \rangle}{\bar{s}} - K_m \frac{\langle \epsilon_S \epsilon_E \rangle}{\bar{e} \bar{s}} \right]}_{\text{noise-induced slowdown}} A \Omega^{-1}$$

317 Where  $A = (k_{cat} \bar{e} \bar{s} / (h k_{cat} \bar{e} K_m + \delta_S (K_m + \bar{s})^2))$  and  $a = K_m / (K_m + \bar{s})$  is the MM  
 318 reaction order. Note that in the general case,  $\langle s \rangle$  is increased and  $\langle p \rangle$  is decreased by  
 319 noise compared to the macroscopic system. The non-linear mesoscopic correction terms

320 vanish when  $s \rightarrow 0$  (first order kinetics),  $s \rightarrow \text{Infinity}$  (zeroth order kinetics) or when  
321  $\Omega \rightarrow \text{Infinity}$  (macroscopic volume). Two of the second moments (Sup. Mat.) for the  
322 general MM case can be written compactly:

$$323 \quad \langle \epsilon_E^2 \rangle = \frac{1}{2} \bar{e} (b_E + 1), \quad \langle \epsilon_S \epsilon_E \rangle = -h \langle \epsilon_E^2 \rangle \frac{k_{cat} \bar{s}}{h a k_{cat} \bar{e} + (\delta_E + \delta_P) (K_m + \bar{s})} \quad 7$$

324 While the present approach builds on that derived by Grima (Grima 2010) the  
325 results and conclusions differ: Grima (Grima 2010) found that noise had no effect on  
326 product concentration in MM reactions. There are two reasons for this discrepancy. First,  
327 Grima assumed no dilution of the substrate ( $\delta_S = 0$ ). As can be seen from equation 6, this  
328 causes the finite volume correction term for  $\langle p \rangle$  to vanish. Thus, Grima's results are in  
329 agreement with the present results for this special case, and all of the results in his papers  
330 are correct. However, following (Grima 2010) exactly for  $\delta_S > 0$  also gives zero  
331 stochastic deviation of the product concentration (see SI), which is incorrect. Here is  
332 why: when the system-size expansion of the chemical master equation is applied to a  
333 system of elementary reactions, as in (Grima 2010) and most other methods, the rate of  
334 product formation,  $V$ , is proportional to the intermediate complex concentration,  $c$ .  
335 Because  $c$  is unimolecular, the method reads this rate as linear (its second derivative is  
336 zero), and so the mean cannot be affected by noise (as proven by Jensen's inequality).  
337 However, in MM reactions  $V$  is decidedly non-linear. Indeed, it is hyperbolic. Thus, the  
338 elementary reaction conceals an essential non-linearity of the system, and will give  
339 incorrect results in the general case. To rectify this, one must apply Eqn. 4 to a QSSA  
340 reduced system of biochemical equations as shown in the SI. Note that the use of  
341 elementary or QSSA reactions is not up to choice: the elementary reaction scheme will

342 give the wrong answer. Note also that applying the method of (Grima 2010) to a QSSA  
343 reduced system also gives incorrect results, for reasons discussed in the SI.

344

345 **Noise propagation in gene expression with hyperbolic filtering.** Because gene  
346 expression noise is costly, there will ostensibly be selection pressure for noise control.  
347 But how can a cell suppress gene expression noise? That is, what are the evolutionary  
348 targets of noise-control? To answer these questions, we must model the process of gene  
349 expression, from transcription through translation, to identify how noise is generated and  
350 filtered.

351 We deviate from previous models of gene expression noise by assuming that  
352 translation obeys non-linear reaction kinetics. Previous models of gene expression noise  
353 have considered translation as a first-order process (Thattai and van Oudenaarden 2001;  
354 Shahrezaei and Swain 2008; Paulsson 2004; Bar-Even et al. 2006; Pedraza and Paulsson  
355 2008; McAdams and Arkin 1997). However, translation in living cells may actually be  
356 closer to zeroth-order than first-order, with a reaction order of  $a = 0.1-0.2$  in log-phase  
357 growth in *E. coli* and budding yeast, based on the fraction of unoccupied ribosomes  
358 (Arava et al. 2003; Zenklusen, Larson, and Singer 2008; Ingolia 2014). We now  
359 investigate the effects of hyperbolic translation kinetics on gene expression noise.

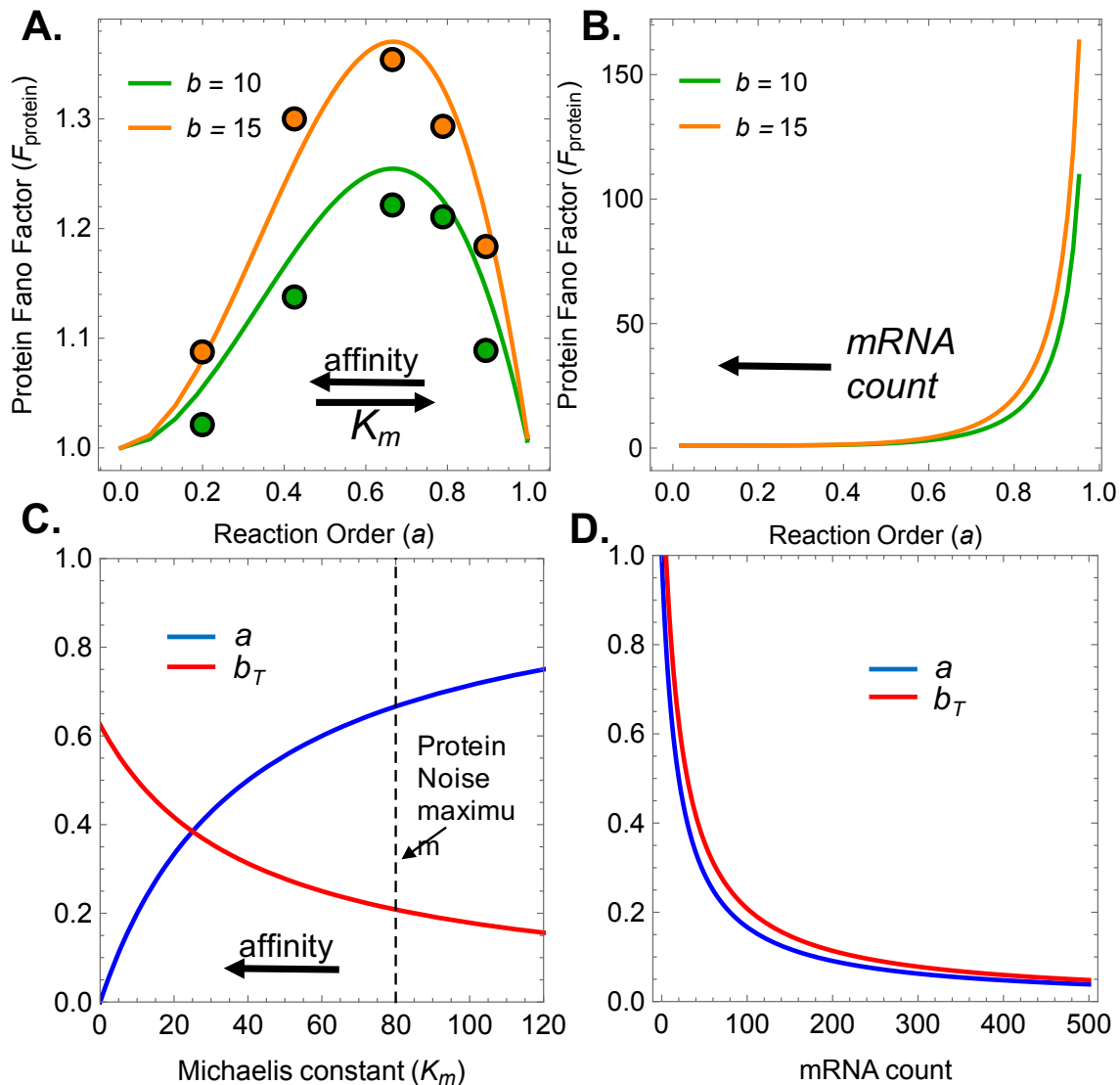
360 Following previous models of gene expression noise (Thattai and van  
361 Oudenaarden 2001; Shahrezaei and Swain 2008; Paulsson 2004; Bar-Even et al. 2006;  
362 Pedraza and Paulsson 2008; McAdams and Arkin 1997), we assume for now that  
363 translation is initiation limited (see “Cell fitness” section below for generalization), such  
364 that it relies only on the concentrations of ribosomes,  $r$ , and mRNA,  $m$ . If mRNA is



365 transcribed in bursts of size  $b_t$ , then the steady-state variance and Fano factor of mRNA  
 366 are (Sup. Mat.),

$$367 \quad \sigma_{mRNA}^2 = \frac{1}{2} \bar{m} (b_t + 1), \quad F_{mRNA} = \frac{1}{2} (b_t + 1) \quad 8$$

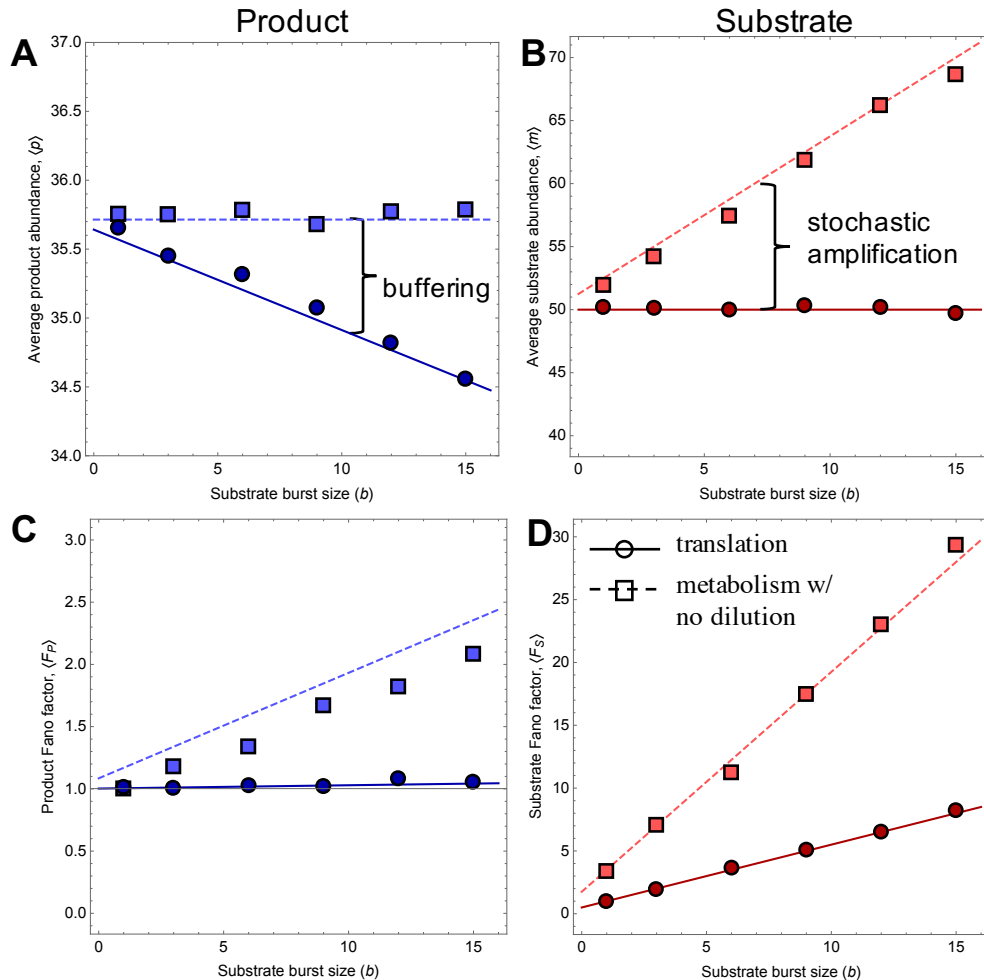
368 which is a special case of (Elf and Ehrenberg 2003) and which is in excellent agreement  
 369 with simulations (Figs 2, 3).



370

371 **Figure 2: Protein noise with hyperbolic filtering.** Translation propagates mRNA noise  
 372 to the protein level, but hyperbolic reaction kinetics ( $0 < a < 1$ ) can substantially ameliorate  
 373 this noise and cause highly non-linear relationships between noise and underlying

374 parameters. Points represent ensemble variance of protein number from the last of 110  
 375 time steps over 10,000 realizations of the SSA (see SI), lines represent theoretical  
 376 prediction (from Eqns 9 and S21). **A, C**) Increasing ribosomal binding affinity (i.e.,  
 377 reducing  $K_m$ ) has two opposing effects on protein noise: it reduces reaction order but  
 378 increases translational bursting. Which effect dominates depends on the population's  
 379 starting point relative to the protein noise peak at  $a^* = 2/3$ . Most species have  
 380 translational  $a < 2/3$ , suggesting that increasing ribosomal binding affinity will actually  
 381 reduce protein noise, even though it increases bursting. **B, D**) Increasing mRNA  
 382 abundance, holding mRNA noise constant, monotonically decreases protein noise for all  
 383 parameter values, as both the reaction order and translational burst parameters decrease  
 384 with increasing mRNA (Sup. Mat. for theoretical derivation of this result). Parameters:  $e$   
 385  $= 5$ ,  $k_{ON} = 1$ ,  $k_{OFF} = 19.5$  (varied in **A**),  $k_{cat} = 0.5$ ,  $\delta_S = 0.1$ ,  $\delta_P = 0.1$ ,  $k_{in} = 4/b$  (varied in **B**).  
 386  $K_m$  was tuned by varying  $k_{OFF}$ ,  $s$  was tuned by varying  $k_{in}$ , while mRNA noise was held  
 387 constant with a burst input of either  $b = 10$  or  $15$ .



388  
 389

390 **Figure 3: Effect of network architecture on the cost of noise.** Theoretical predictions  
 391 (Eqns. 6-8, S20, S21) for metabolic reactions (dashed lines;  $h = 1$ ,  $\delta_S = 0$ ) and translation  
 392 (solid lines;  $h = 0$ ,  $\delta_S = 0.1$ ) compared to ensemble average (blue) and variance (red) of

393 1,000-5,000 SSA realizations (points). **A**) Noise-induced slowdown causes an  
 394 accumulation of substrate (“stochastic amplification”) which completely buffers reaction  
 395 flux for metabolic reactions when substrate dilution/degradation is zero ( $\delta_S = 0$ ) **(B)**.  
 396 Because mRNA is not consumed in translation ( $h = 0$ ), its concentration is independent of  
 397 reaction rate so does not experience stochastic amplification. Thus, translation is not  
 398 buffered by network feedback. **C,D**) The effect of network architecture on substrate (C)  
 399 and product (D) noise.  $e = 1$ ,  $k_{ON} = 1$ ,  $k_{OFF} = 19.5$ ,  $k_{cat} = 0.5$ ,  $\delta_P = 0.01$ , and  $k_{in} = 5/b$   
 400 (solid lines) or  $0.357/b$  (dashed lines).  
 401

402 The variance in protein numbers at steady-state is,  $\langle \epsilon_{protein}^2 \rangle \cong \bar{p}(1 +$   
 403  $a^2 b_T F_{mRNA} + CV_R^2 \bar{p}/2)$ , where the parameter  $a^2 = [K_m/(K_m + m)]^2$  is the squared reaction  
 404 order and  $CV_{ribo}^2 = \sigma_{ribo}^2/\bar{r}^2$  is the ribosomal noise (Sup. Mat.). The translational burst  
 405 parameter,  $b_T$ , is the number of proteins synthesized by a single mRNA molecule over its  
 406 lifetime, which with hyperbolic translation kinetics is,  $b_T = k_T r / ((K_m + m)\delta_S)$ .  
 407 Previous work assumed first-order translation, such that  $b_T = k_T r / (K_m \delta_S)$  (Thattai and  
 408 van Oudenaarden 2001; Shahrezaei and Swain 2008). Introducing non-linearity causes  
 409 the translational burst size to become a non-linear, decreasing function of mRNA  
 410 abundance: all else equal, an mRNA molecule competing with fewer other mRNA’s will  
 411 be translated more often in its lifetime, increasing its burst size. Alternatively, in the  
 412 zeroth order regime ( $m \gg K_m$ ), translational bursting goes to zero because ribosomes are  
 413 so oversaturated with mRNA’s that any given mRNA is more likely to be degraded than  
 414 translated. The protein noise strength and noise, then, are, respectively,

415 
$$F_{protein} \cong 1 + a^2 b_T F_{mRNA} + CV_{ribo}^2 \bar{p}/2$$

416

$$CV_{protein}^2 \cong \underbrace{\frac{1}{\bar{p}} + \frac{a^2 b_T F_{mRNA}}{\bar{p}}}_{\text{intrinsic noise}} + \underbrace{\frac{CV_{ribo}^2}{2}}_{\text{extrinsic noise}}$$

9

417 Eqn. 9 agrees well with simulation results (Figs. 2, 3) and retrieves the result of  
 418 Thattai and van Oudenaarden (Thattai and van Oudenaarden 2001) ( $F_{protein} \cong 1 + b_T$ )  
 419 as a special case when there is no transcriptional bursting ( $b_t = 1 \Rightarrow F_{mRNA} = 1$ ), no  
 420 ribosome fluctuations ( $CV_{ribo}^2 = 0$ ), and translation is first order (i.e., linear) in mRNA  
 421 abundance ( $a = 1$ ). The first terms in  $F_{protein}$  and  $CV_{protein}^2$  (1 and  $1/\bar{p}$ , respectively) are  
 422 simply Poisson terms reflecting low copy number fluctuations (Paulsson 2004; Pedraza  
 423 and Paulsson 2008), while the second and third terms introduce factors that increase noise  
 424 above Poisson levels. The second terms ( $a^2 b_T F_{mRNA}$  and  $a^2 b_T F_{mRNA}/\bar{p}$ , respectively)  
 425 incorporate the consequences of translational and transcriptional bursting. These terms  
 426 propagate mRNA fluctuations,  $F_{mRNA}$ , to the protein level. Importantly, hyperbolic  
 427 translation kinetics ( $0 < a < 1$ ) dampen the effect of mRNA fluctuations on protein noise,  
 428 thus attenuating the propagation of mRNA noise via a process that might be called  
 429 “hyperbolic filtering”. Together the first and second terms are considered “intrinsic  
 430 noise” (Swain, Elowitz, and Siggia 2002; Raser and O’Shea 2005; Paulsson 2004; Bar-  
 431 Even et al. 2006; Sanchez and Golding 2013). The last term ( $CV_{ribo}^2 \bar{p}/2$  and  $CV_{ribo}^2/2$ )  
 432 accounts for fluctuations in ribosome numbers, which create positively correlated  
 433 fluctuations among protein species called “extrinsic noise” (Swain, Elowitz, and Siggia  
 434 2002; Raser and O’Shea 2005; Paulsson 2004; Bar-Even et al. 2006; Sanchez and  
 435 Golding 2013). This extrinsic noise term dominates for abundant proteins (large  $p$ ),  
 436 consistent with the observation that protein noise scales inversely with protein abundance

437 for low copy number proteins, but is independent of abundance for high copy number  
438 proteins (Blake et al. 2003; Raser and O’Shea 2004; Bar-Even et al. 2006).

439 Translation initiation in *E. coli* and yeast has a reaction order of  $a \sim 0.1-0.2$  (3,  
440 49), which is comfortably to the left of the protein noise peak at  $a^* = 2/3$  (see Sup. Mat.;  
441 Fig 2a). This implies that, contrary to predictions based on first-order translation  
442 initiation models (Thattai and van Oudenaarden 2001; Raser and O’Shea 2005), evolving  
443 or engineering increased binding affinity of mRNA to ribosomes (i.e., reduced  
444 translational  $K_m$ ) will actually lower protein noise, not increase it. This occurs because  
445 reducing  $K_m$  has two complementary effects on protein noise: it increases translational  
446 bursting,  $b_T$ , but reduces reaction order,  $a$ , which attenuates noise via a phenomenon we  
447 term “zeroth-order insensitivity”. The latter effect is not accounted for in previous gene  
448 expression models, but has implications for the evolutionary and synthetic targets of  
449 noise-amelioration (see Discussion). Hyperbolic filtering also has practical implications  
450 when making inferences from data. For example, translational burst size,  $b_T$ , can be  
451 estimated from Eqn 9 with knowledge of the mRNA and protein statistics; however,  
452 assuming first order reaction kinetics ( $a = 1$ ) can lead to large underestimates of this  
453 parameter from data. For example, for  $a \sim 1/5$ , assuming first order translation kinetics  
454 can lead to a 25-fold underestimate of translational burst size.

455

#### 456 **Correlation between mRNA and protein noise with hyperbolic translational**

457 **filtering:** We now quantify the effect of hyperbolic translation kinetics on the correlation  
458 between mRNA and protein noise. Empirical work in yeast, *E. coli* and human cell lines  
459 have found that there is no correlation between mRNA and protein levels in the cell, or

460 between mRNA and protein noise (Chen et al. 2002; Ghaemmaghami et al. 2003;  
 461 Taniguchi et al. 2010; Gygi et al. 1999). Here I show that this lack of correlation can be  
 462 explained, in part, by modeling translation as a hyperbolic reaction.

463 Assuming as before that mRNA lifetimes are much shorter than protein lifetimes  
 464 ( $\delta_{mRNA} \gg \delta_{protein}$ ), the covariance between mRNA and protein noise,  $\sigma_{mp}$ , is found to be,  
 465  $\sigma_{mp} \cong ab_T \sigma_{mRNA}^2$ . The linear regression coefficient,  $\beta_{mp} \equiv \sigma_{mp} / \sigma_{mRNA}^2$  is then,  
 466  $\beta_{mp} \cong ab_T$ . Both of these show that hyperbolic translation kinetics ( $a < 1$ ) diminish the  
 467 statistical association between mRNA and proteins. This is most usefully illustrated by  
 468 the Pearson correlation coefficient,  $\rho \equiv \sigma_{mp} / \sigma_{mRNA} \sigma_{protein}$ , which is approximately  
 469 equal to,

$$470 \quad \rho \cong a \sqrt{\left(\frac{\delta_{protein}}{\delta_{mRNA}}\right) \left(\frac{b_T F_{mRNA}}{1 + a^2 b_T F_{mRNA} + CV_{ribo}^2 \bar{p} / 2}\right)} \quad 10$$

471 Consistent with the theoretical predictions of Tanaguchi et al (3), Eqn. 10 shows that the  
 472 mRNA-protein correlation is increased with increasing transcriptional noise strength,  
 473  $F_{mRNA}$ , is diminished with larger differences in lifetimes of mRNA and protein  
 474 ( $\delta_{protein} / \delta_{mRNA} < 1$ ), and is diminished by extrinsic noise ( $CV_{ribo}^2 \bar{p} / 2$ ). In addition, the  
 475 covariance, linear regression coefficient, and correlation between mRNA and protein  
 476 fluctuations all increase with increasing translational bursting,  $b_T$ . The most important  
 477 novel result here is that the correlation between mRNA noise and protein noise is  
 478 diminished by the hyperbolic filtering of translation (that is, when  $a < 1$ ).

479

480 **Network architecture, feedback, and buffering:** While noise slows a reaction for a  
 481 given substrate concentration, network feedbacks cause adjustment to substrate levels

482 that may in turn alter reaction rates. Specifically, if a reaction proceeds more slowly, then  
483 substrate levels accumulate, which then cause the reaction to speed up again. Accounting  
484 for this feedback is crucial in quantifying the costs of noise. We now ask: How exactly  
485 does this feedback influence the total flux through noisy reaction networks? To address  
486 this problem, we now analyze how network architecture creates feedbacks that alter  
487 chemical flux, potentially compensating for noise-induced slowdown. We treat three  
488 limiting cases: 1)  $h = 1, \delta_S = 0$ ; 2)  $h = 1, \delta_S > 0$ ; and 3)  $h = 0$ .

489         *Cases 1 ( $h = 1, \delta_S = 0$ ): Stochastic substrate amplification and buffering:* One  
490 important feedback in noisy networks is the amplification of substrate concentration by  
491 noise (Eqn. 6), which has been studied in a process called “stochastic focusing”  
492 (Paulsson, Berg, and Ehrenberg 2000). Eqn. 6 shows that stochastic amplification will  
493 occur whenever substrate is consumed by the reaction ( $h = 1$ ), as in metabolic reactions.  
494 Consequently, this steady-state adjustment of substrate creates intrinsic buffering of  
495 reaction flux that ameliorates noise-induced slowdown: the reaction is slowed by noise,  
496 allowing substrate to accumulate, which then speeds up the reaction. The extent of  
497 buffering depends on the substrate dilution parameter,  $\delta_S$ . In the MM network studied by  
498 Grima (Grima 2010; Grima 2009b), for example, the substrate entered the reaction  
499 volume at rate  $k_{in}$  and exited only once consumed by the focal reaction at rate  $V$ . Flux  
500 balance requires that  $V = k_{in}$  in such a system, so that the reaction rate in steady-state is  
501 unaffected by noise, as found by Grima (Grima 2010; Grima 2009b). Thus, under the flux  
502 constraint  $V = k_{in}$ , stochastic amplification completely buffers the system against noise-  
503 induced slowdown (Fig 3).

504           *Case 2 ( $h = 1, \delta_S > 0$ ): Incomplete amplification and buffering.* However, the  
505 precise balancing of  $V$  with  $k_{in}$  does not apply in living cells, where substrate exits the  
506 system via multiple alternative channels ( $\delta_S > 0$ ), primarily dilution by cell division, but  
507 also through excretion, titration by other reactions and enzymatic degradation.  
508 Extracellular substrate also experiences non-zero  $\delta_S$ , as it is diluted by bulk flow (such as  
509 in a chemostat) or by uptake from neighboring cells. Accounting for these effects  
510 modifies the flux constraint to:  $k_{in} = V + \delta_S s$ . With this network architecture, stochastic  
511 amplification compensates for some of the diminution of  $V$  by noise, but not all.  
512 Consequently, stochastic amplification fails to completely buffer metabolic reactions in  
513 living cells against noise-induced slowdown.

514           *Case 3 ( $h = 0, \delta_S > 0$ ): Gene expression is especially sensitive to noise-induced*  
515 *slow-down; no stochastic amplification, no buffering.* The  $V = k_{in}$  flux constraint is  
516 entirely abolished in template-mediated reactions ( $h = 0$ ), such as protein synthesis via  
517 translation of mRNA. Because mRNA is not consumed by translation, steady-state  
518 mRNA concentration is set by the flux constraint  $k_{in} = \delta_S s$ , which is independent of  
519 translation rate,  $V$ . Thus, stochastic amplification does not occur in translation and so  
520 cannot buffer translation from mRNA noise (Fig. 3). As a consequence, gene expression  
521 is unusually sensitive to noise.

522

523 **Cell fitness.** Noise slows biochemical reactions, but how does it affect fitness? We now  
524 consider a reaction network representing a coarse-grained model of cell growth. We then  
525 analyze this reaction network in the stochastic case using the method developed above in  
526 order to quantify the effects of gene expression noise on the average rate of cell division.



527 During balanced exponential growth, it has been shown empirically that fitness,  
528 here defined as the division rate of a cell, is equivalent to the total rate of protein  
529 synthesis by translation in a cell (Shahrezaei and Marguerat 2015; Scott et al. 2010; Scott  
530 et al. 2014). We denote the rate of cell division by the parameter  $\lambda$ . In general, biomass  
531 synthesis involves the import of extracellular nutrients followed, ultimately, by their  
532 conversion into new biomass via translation. Translation initiation involves the binding of  
533 mRNA, in abundance  $m$ , to a ribosome, in abundance  $r$ , followed by elongation via the  
534 polymerization addition of amino acids (aa) to a growing protein chain. Elongation rate is  
535 limited by the intracellular nutrient, in abundance  $s_I$ , in scarcest supply (i.e., von Leibig's  
536 Law of the Minimum (Droop 1974; De Baar 1994; Tilman 1982)), which is typically  
537 either ATP or aa-charged tRNA's. In eutrophic conditions, where nutrients are saturating,  
538 translation is initiation-limited (Shah et al. 2013), whereas in oligotrophic conditions,  
539 where nutrients are scarce, translation is elongation-limited (Tuller et al. 2010).

540 This coarse-grained cell model follows a QSSA system of ODE's (S22) with  
541 growth rate,

$$542 \quad \frac{dp}{dt} = k_T r \frac{s_I m}{K_1 + K_2 s_I + K_3 m + s_I m} - \lambda p \quad 11$$

543 which is equivalent to a compulsory-order two-substrate enzymatic reaction with  
544 independent (i.e., non-cooperative) binding events (Ingalls 2013).  $k_T$  is the translation  
545 elongation rate per mRNA per ribosome (which absorbs the number of  $S_I$  per protein),  
546  $K_1/ K_2/ K_3$  are affinity constants, and  $\lambda$  is the protein dilution rate, which during steady-  
547 state growth is equivalent to the rate of cell division. Eutrophic and oligotrophic regimes,  
548 corresponding to initiation- and elongation-limited translation, respectively, are found  
549 from this equation in the limit as  $s_I$  or  $m$ , respectively, goes to infinity. Eqn. 11 can be

550 rearranged to form a Monod equation, which maps growth rate to extracellular nutrient  
 551 supply with a number of practical and conceptual benefits (S24).

552 The expected steady-state rate of cell division (cells/time) is found by setting Eqn.  
 553 11 to zero and solving for  $\lambda$ ,

$$554 \quad \langle \lambda \rangle = \frac{\langle V \rangle}{p_{ss}} \quad 12$$

555  $p_{ss}$  is defined as the species-specific total protein content of a cell (protein/cell), which is  
 556 assumed constant for a given species,  $V$  is the synthesis rate of proteins (protein/time).

557

558 **Fitness with noise.** Following the method above, we write down a QSSA reduced system  
 559 of equations for growth rate according to our coarse-grained cell model (S22), giving a  
 560 system in five equations describing the time evolution of mRNA, uptake proteins,  
 561 extracellular substrate, intracellular substrate and protein biomass. We assume that  
 562 mRNA and extracellular nutrients undergo a linear birth-death process, supplied with  
 563 burst inputs of size  $b_t$  and  $b_{SE}$  respectively, and ignore ribosome fluctuations.  
 564 Intracellular nutrients are consumed by protein synthesis, but mRNA is not. Expected  
 565 fitness, then, is (Sup. Mat.):

$$566 \quad \langle \lambda \rangle \cong \frac{\bar{V}}{p_{ss}} \left[ 1 - \underbrace{\frac{\overbrace{\langle \epsilon_m^2 \rangle}^{mRNA\ noise} A_1}{\bar{m} \Omega}}_{\text{cost of gene expression noise}} - \frac{\overbrace{\langle \epsilon_{PU}^2 \rangle}^{protein\ noise} A_2}{\bar{p}_U \Omega} - \frac{\overbrace{\langle \epsilon_{SE}^2 \rangle}^{extracellular\ noise} A_3}{\bar{s}_E \Omega} \right] \quad 13$$

567 Where the  $A_i$ 's are positive constants given in the SI. In reality, the protein noise term  
 568 will contain contributions from all proteins involved in the pathway importing and  
 569 processing nutrients to their final form utilized in translation. The MCA model of Wang  
 570 and Zhang (Wang and Zhang 2011) then fits within this term. The extracellular noise

571 term quantifies the strength of selection favoring phenotypic “bet hedging” (Starrfelt and  
 572 Kokko 2012) strategies in response to environmental unpredictability.

573 Surprisingly, transcriptional noise directly reduces fitness. This occurs because  
 574 fluctuations in mRNA levels reduce the average translation rate. This is counter to the  
 575 typical view, which holds that transcriptional noise impinges on cell function only to the  
 576 extent that it propagates to generate protein noise. Unlike the cost of protein or  
 577 extracellular nutrient noise, which require a number of parameter estimates that are  
 578 highly context-dependent, the cost of mRNA noise is easily estimated from available data  
 579 if we assume that cells are growing in their eutrophic regime ( $s_I \rightarrow \infty$ ), where  
 580 growth rate is limited by translation initiation. Taking the eutrophic limit of Eqn 11  
 581 gives,  $dp/dt \cong k_T r m / (K_m + m) - \lambda p$ , which is simply a MM reaction with mRNA  
 582 serving as the substrate and ribosomes as the catalyst. Noting that fluctuations in total  
 583 ribosome number and mRNA number are uncorrelated, we find from equations 4, 7 and  
 584 9,

$$585 \quad \langle \lambda \rangle \cong \frac{\bar{V}}{p_{ss}} \left[ 1 - \underbrace{\frac{(b_t+1)}{2\bar{m}} a(1-a)\Omega^{-1}}_{\text{cost of transcriptional bursting}} \right] \quad 14$$

586 Where  $\bar{V} = k_T r m / (K_m + m)$ . The cost of transcriptional bursting vanishes with zeroth-  
 587 or first-order translation ( $a = 0$  and  $1$ , respectively) and/or as cell volume and/or mRNA  
 588 abundance become large.

589 During balanced exponential growth of *E. coli* in rich medium:  $m \sim 1400$   
 590 molecules,  $r \sim 21,000$  molecules, and  $k_T \sim 21aa/s \sim 4.6$  proteins/min (given a mean  
 591 protein length of  $275aa$ ). The total protein count of a cell is  $p_{ss} = 2.4 \times 10^6$  molecules.  
 592 During growth in nutrient rich medium,  $(1 - a) \sim 0.85$  (Arava et al. 2003; Ingolia 2014;

593 Zenklusen, Larson, and Singer 2008). Applying these parameter estimates to  $\bar{V}/p_{ss}$  gives  
594 a doubling time for *E. coli* of ~30 minutes. Similar calculations with parameters for the  
595 budding yeast *Saccharomyces cerevisiae* ( $m \sim 15,000-60,000$ ,  $r \sim 243,000$ ,  $k_T \sim 15aa/s \sim$   
596  $2$  proteins/min (average protein length  $450aa$ ),  $p_{ss} = 5 \times 10^7$ ,  $a = .85$ ) gives a doubling  
597 time of ~120 minutes. Both values are in close agreement with experiments in nutrient  
598 rich media, providing an important benchmark.

599 In an *E. coli* cell, according to Eqn. 14, a mutation causing one unit increase in the  
600 total mRNA Fano factor results in an average loss of 7.5 proteins/min, 223 proteins/cell  
601 division, and  $1.5 \times 10^{-4}$  cells/generation, such that selection against mRNA noise is more  
602 than 10,000 times stronger than random drift in this species (the effective population size  
603 determines the strength of random allele frequency fluctuations and is  $N_e = 2 \times 10^8$  in *E.*  
604 *coli*). In contrast, the same mutation in yeast, which has 40-fold higher  $m$ , has a cost of  
605  $2.1 \times 10^{-6}$  cells/generation, which is only about 20 times stronger than random drift ( $N_e =$   
606  $1 \times 10^7$  in *S. cerevisiae*). This illustrates how fitness magnifies small differences in  
607 reaction rates by integrating over the whole cell and cell lifetime, and highlights the  
608 importance of cell volume (or more directly, molecule number) in influencing the  
609 efficacy of selection against noise.

610

611 **Rapid decay of selection efficacy with cell volume.** As we just showed, the cost of  
612 noise depends on reaction volume. Because unicellular species differ over several orders  
613 of magnitude in their cell volumes, it may be possible to use these two facts in order to  
614 make predictions about taxonomic patterns of gene expression noise.

615 The fate of a mutation that modifies cellular noise depends on the relative  
616 magnitude of selection and random genetic drift. Selection efficiently overcomes random  
617 drift provided that  $2N_e|s| > 1$  (Kimura 1957; Kimura 1962). The boundary at  $2N_e|s| = 1$   
618 defines the "drift barrier" separating efficient and inefficient selection regimes (Sung et  
619 al. 2012; Lynch 2007b). The drift barrier acts as a kind of evolutionary attractor  
620 preventing organismal perfection: recurrent deleterious mutations with effects smaller  
621 than  $2N_e|s|$  push species away from their adaptive optima, and selection is too weak to  
622 push them back. Thus, the balance between recurrent mutation, drift and selection  
623 establishes a suboptimal phenotypic steady-state. We now derive a scaling relationship  
624 between cell volume and selection efficacy demonstrating that the drift boundary for  
625 noise attenuation scales superlinearly with cell volume.

626

627 **Selection strength scales inversely with cell size.** Selection strength,  $s$ , is defined as the  
628 difference in absolute fitness of mutant ( $m$ ) and resident ( $r$ ) alleles scaled by the mean  
629 fitness of the population. For a rare mutant we have:  $s = (\lambda_m - \lambda_r)/\lambda_r$ . This is simply  
630 the difference in the absolute number of daughter cells produced by the mutant per  
631 resident doubling. For simplicity, we assume a haploid population and that the mutant  
632 and resident genotypes differ only in transcriptional noise. When the only source of noise  
633 is intrinsic mRNA fluctuations, the selection coefficient is approximately (see SI for  
634 derivation),

$$635 \quad s \cong \Delta F_{mRNA} \frac{K_m}{(K_m + m)^2} \Omega_c^{-1} = \Delta CV_{mRNA}^2 a(1 - a) \Omega_c^{-1} \quad 15$$

636  $\Omega_c$  denotes the volume of the reaction compartment. The sign of selection is determined  
637 by  $\Delta F = (F_r - F_m)$  or  $\Delta CV^2 = (CV_r^2 - CV_m^2)$ , the difference in Fano factors (noise

638 strength) or squared coefficients of variation (noise), respectively, between resident and  
639 mutant alleles. A mutation that increases transcriptional noise will be deleterious ( $s < 0$ ),  
640 while a decrease in noise is advantageous ( $s > 0$ ). More generally,  $s \cong \sum_i^n (\Delta F_i A_i) \Omega_{ci}^{-1}$ ,  
641 where the sum is over all  $n$  sources of noise affected by the mutation, the  $A_i$  are constants,  
642 and  $\Omega_{ci}$  is the reaction volume of the  $i^{th}$  noise source. Note that selection cannot act on  
643 noise levels in first or zeroth order reactions ( $a = 1$  or  $0$ , respectively), since noise does  
644 not alter the mean in these cases. Noise will, however, alter the heritability (e.g.,  
645 penetrance) of a trait (Wang and Zhang 2011), which will reduce the response to  
646 selection, but we do not consider this effect here. Importantly, the strength of selection  
647 against noise explicitly depends on reaction volume.

648

649 **Drift strength scales positively with cell size.** A species population size tends to scale  
650 inversely with its body size: a species must divide a limited resource pool amongst its  
651 members, resulting either in a large number of small individuals or a small number of  
652 large ones (Damuth 1981; Damuth 1987). For multicellular species the empirical scaling  
653 relationship between body mass and population size, known as “Damuth’s Law”, follows  
654 a power law with a scaling exponent of  $-3/4$ , likely reflecting the  $3/4$  power scaling of body  
655 size with energetic demands (Damuth 1981; Damuth 1987). However, in unicellular  
656 species’ metabolic rate scales either linearly (eukaryotes) or superlinearly (prokaryotes)  
657 with cell mass (DeLong et al. 2010). Assuming that  $N_e$  scales linearly with  $N$ , that mass  
658 scales linearly with cell volume (i.e., cells have constant density), then  $N_e = B\Omega_w^{-g_{drift}}$ ,  
659 where the subscript on  $\Omega_w$  denotes whole cell volume and  $g_{drift} = [3/4, 5/4]$  is the scaling  
660 exponent. The constant,  $B$ , is proportional to the population size of a species with unit

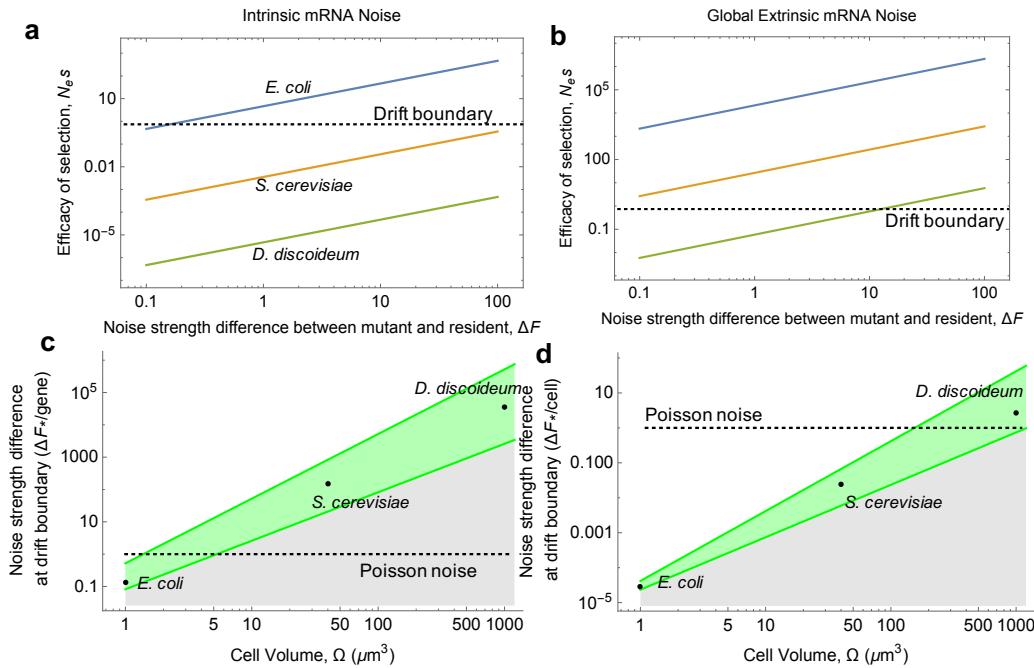
661 volume, and absorbs cell density as the proportionality constant between mass and  
 662 volume.  
 663  
 664 **Combining forces: Selection efficacy scales superlinearly with cell size.** Combining  
 665 these results implies that the efficacy of selection against noise is proportional to  
 666  $(\Omega_w^{-g_{drift}} \Omega_c^{-1})$ . For translation, it is possible to establish a relationship between  $\Omega_c$  and  
 667  $\Omega_w$ , allowing us to reduce the number of independent variables.

668 Protein synthesis occurs in the cytoplasm/cytosol, suggesting that reaction-  
 669 compartment volume may simply equal whole-cell volume ( $\Omega_c \sim \Omega_w$ ) for translation. This  
 670 relationship appears to hold between *E. coli* and yeast ( $m = 1500$  and  $60,000$ ,  
 671 respectively, with a 40 fold difference in cell volume) and within mammalian cells as  
 672 they change volume throughout the cell cycle (Kempe et al. 2015; Padovan-Merhar et al.  
 673 2015). However, Gillooly et al (Gillooly et al. 2005) found that RNA concentration scales  
 674 as body mass to the  $-1/4$  power over a broad range of eukaryotes, corresponding to RNA  
 675 number scaling with mass to the  $3/4$  power. Although unicells were underrepresented in  
 676 their sample, this agrees well with the scaling of mRNA between *E. coli* and the amoeba  
 677 *Dictyostelium discoideum* where  $m = 155,000$  and cell size is approximately 500-1000  
 678 fold larger than *E. coli*, such that  $\Omega_c \sim \Omega_w^{2/3}$  to  $\Omega_w^{3/4}$ .

679 Thus, for translation, we assume that reaction volume scales with whole cell  
 680 volume according to,  $\Omega_c = \Omega_w^{-g_{sel}}$ , where  $g_{sel} = [3/4, 1]$ . Therefore, the efficacy of  
 681 selection against mRNA noise can be written as an allometric scaling law:

$$682 \quad N_e |s| \cong \Delta CV_{mRNA}^2 a(1 - a) B \Omega_w^{-g} \quad 16$$

683 The superlinear scaling exponent,  $g = g_{sel} + g_{drift} = [3/2, 2^{1/4}]$ , makes selection against  
 684 noise extremely weak in large cells. For example, all else equal, selection against a  
 685 mutation that increases expression noise is thirty thousand to one million times stronger  
 686 in *E. coli* than in a species, such as the amoeba *Dictyostelium discoideum*, which is about  
 687 1000 times larger (Fig. 4).



688

689 **Figure 4: Power-law scaling of selection efficacy with cell volume.** **A,B)** The predicted  
 690 efficacy of selection,  $N_e|s|$ , for *E. coli*, *S. cerevisiae* and *Dictyostelium discoideum*, as a  
 691 function of the jump size in noise caused by a mutation,  $|\Delta F|$ , for mutations acting at a  
 692 single locus (left-hand column) and mutations acting on the noise of the global  
 693 transcriptome (right-hand column). The mRNA concentrations and effective population  
 694 size ( $N_e$ ) estimates used for the figure are found in Table S1 ( $m = 1500, 60,000$  and  
 695  $155,000$ , respectively, and ploidy-adjusted  $N_e = 2 \times 10^8, 1 \times 10^7$  and  $1.1 \times 10^5$ , respectively).  
 696 Reaction order was assumed to be  $a = 0.8$  for all three species. The “drift boundary” is  
 697  $2N_e|s| = 1$ . **C,D)** The predicted noise jump size at the drift boundary as a function of cell  
 698 volume. The green region represents a range of values for the scaling exponent (-2 to -  
 699  $3/2$ ), reaction order (0.75-0.9), and in **C)** the gene number (3250-12,500). Above the  
 700 green region (or individual points) selection is strong, below mutations are effectively  
 701 neutral.  
 702



703 Solving for  $|\Delta F|$  at the drift boundary ( $2N_e|s| = 1$ ), gives the jump size in noise  
704 strength “visible” to natural selection. If the only cost of transcriptional noise is reduced  
705 translation, then the selectively impermissible jump sizes in transcriptional noise follows  
706 the scaling relationship,

$$707 \quad |\Delta F_{mRNA}^*| > b\Omega_w^g \quad 17$$

708 where  $b = (K_m + m)^2 / BK_m$ . Mutations causing jumps smaller than this critical size are  
709 “effectively neutral”, and thus invisible to selection acting on translation rate. Adaptive  
710 mutations that decrease noise by an amount less than this jump size cannot efficiently  
711 spread, while deleterious mutations with effect size below this threshold cannot be  
712 efficiently purged by selection.

713 Parameterizing our model with literature estimates of cell volume, effective  
714 population size, genome size, ribosomal occupancy, and translation kinetics, we predict  
715 that *E. coli* should be strongly selected to minimize intrinsic transcriptional noise simply  
716 because this noise slows translation. Whereas the cost of transcriptional noise is too weak  
717 per gene (on average) in *S. cerevisiae* or *Dictyostelium* for selection to maintain intrinsic  
718 noise at the Poisson minimum if noise-induced slowdown of translation is the only cost  
719 of mRNA noise. Fig 4 shows that in cells below about 5fL (most prokaryotes), selection  
720 favors suppression of intrinsic transcriptional noise to its theoretical minimum defined by  
721 Poisson statistics, even if the only cost of transcription noise is reduced translation rate.  
722 But because of the superlinear scaling of selection efficacy, cells only 100 times larger  
723 (most eukaryotes) cannot efficiently oppose intrinsic noise jumps even as large as  $\Delta F =$   
724 1000 (again, if this is the only selection pressure for noise-minimization acting).  
725 Alternatively, jumps in global transcriptional noise (i.e., extrinsic noise) have a much

726 greater selective effect, meaning that traits affecting extrinsic noise, such as ribosome  
727 fluctuations (Eqn 9), should be under strong selection for noise-amelioration even in  
728 moderately sized species.

729

## 730 **DISCUSSION**

731 Biochemical reactions are non-linear, but are typically modeled as linear in stochastic  
732 models of gene expression due to technical limitations (van Kampen 2007). Here we have  
733 introduced a novel procedure for easily and more accurately incorporating non-linearity  
734 into models of coupled stochastic chemical reactions. We have found that hyperbolic  
735 reaction kinetics filter input noise in such a way that the average reaction rate is slowed  
736 by noise, while attenuating noise propagation. These results have a number of  
737 implications.

738

739 **A drift barrier hypothesis for gene expression noise and robustness.** Cell size places  
740 severe constraints on the evolution of noise suppression. The magnitude of noise-induced  
741 slowdown of Michaelis-Menten reactions is inversely proportional to the reaction  
742 volume, which follows intuitively from the fact that, for a given molecular density,  
743 greater volumes contain more molecules and thus generate less noise. Lynch (Lynch  
744 2007b) has previously noted the “genomic perils of evolving large body size”. Large  
745 bodied species tend to be less efficient at purging deleterious mutations due to the  
746 positive scaling of genetic drift strength with body size, leading to a number of broad  
747 taxonomic patterns including greater genomic mutation rate, transposon abundance and  
748 intron number in unicellular eukaryotes and metazoans relative to prokaryotes and

749 viruses (Lynch 2007b; Lynch 2007a; Sung et al. 2012). Unique to gene expression noise,  
750 cell size affects not only drift but also fitness, resulting in a negative superlinear scaling  
751 of selection efficacy with cell size (Eqns. 15-17). Unicellular species span more than four  
752 orders of magnitude in cell volume ( $\sim 0.1 \mu\text{m}^3$  Mycoplasma to more than  $1000 \mu\text{m}^3$  in  
753 amoebazoans and other protists), with a sharp taxonomic divide separating prokaryotes  
754 and eukaryotes. “Typical” cell diameters are 0.5-5 $\mu\text{m}$  for prokaryotes and 10-100 $\mu\text{m}$  for  
755 unicellular eukaryotes. Taking a conservative value of a 100-fold average volume  
756 difference between prokaryotes and eukaryotes, this translates to a 1000-10,000-fold  
757 difference in the average efficacy of selection against noise for non-compartmentalized  
758 processes (Eqn. 16).

759         Of course, numerous other fitness effects of gene expression noise occur (Eqn. 13  
760 and (Fraser et al. 2004; Wang and Zhang 2011; Lehner 2008; Raj and van Oudenaarden  
761 2008; Balázsi, van Oudenaarden, and Collins 2011; Raser and O’Shea 2005)), but these  
762 must all scale with cell (or compartment) volume, such that selection for noise-  
763 amelioration will often be weaker in eukaryotes than prokaryotes. It is possible that the  
764 greater transcriptional complexity in eukaryotic cells, which employ multisubunit  
765 regulatory complexes, enhancer regions, heterochromatic gene silencing, nuclear  
766 localization into “transcription factories” and other features that greatly increase the  
767 burstiness of transcription, evolved only once greater cell size freed populations from the  
768 selective constraint of noise minimization. The cause of the size disparity between  
769 prokaryotes and eukaryotes is thought to be due to differences in cellular respiration:  
770 prokaryotes synthesize ATP via the electron transport chain embedded in the cell’s  
771 plasma membrane, which requires a high surface area to volume ratio for the whole cell

772 in order to balance energy production and demand. Respiration via mitochondria free  
773 eukaryotes from this constraint, allowing them to achieve large cell sizes without  
774 bankrupting the cell's energy budget. Thus, the evolution of mitochondria may have  
775 paved the way for the greater transcriptional sophistication of eukaryotes.

776

777 **Translation attenuates the propagation of transcriptional noise:** Gene expression  
778 noise is generated at two levels: transcription and translation. While the molecular  
779 mechanisms controlling noise at each of these levels are becoming increasingly well  
780 understood, much remains unclear about the propagation of noise between levels. I have  
781 shown here that hyperbolic translation kinetics greatly attenuate the propagation of  
782 mRNA noise to the protein level. Given parameter estimates of translation kinetics in *E.*  
783 *coli* and yeast, the results here imply that hyperbolic translation will reduce mRNA noise  
784 by between 25 and 100 fold, substantially attenuating the propagation of mRNA noise to  
785 the protein level. The hyperbolic filtering of mRNA noise by translation also manifests as  
786 a substantial reduction in the correlation between mRNA and protein abundance (Eqn.  
787 10). This result is consistent with data from *E. coli* showing that there is zero correlation  
788 between a gene's mRNA noise and its protein noise (Taniguchi et al 2010). We have  
789 shown that hyperbolic translation kinetics can help explain this lack of correlation, since  
790 the correlation coefficient between mRNA and protein noise is proportional to the  
791 reaction order parameter,  $\alpha$ , which is between 0.1 and 0.2 in *E. coli*, implying a large  
792 reduction in the correlation. Geometrically, one can visualize this diminution of  
793 correlation by observing the Michaelis-Menten rate curve: on the flat part of the curve, all

794 mRNA concentrations lead to the same rate of protein synthesis, thus creating a zero  
795 correlation between mRNA and protein.

796 A major implication of these results is that the vast majority of all intrinsic protein  
797 noise comes from translation, not transcription. Consequently, if protein noise, rather  
798 than mRNA noise, determines the functional consequences of gene expression noise, then  
799 natural selection on transcriptional noise will be extremely weak, and synthetic systems  
800 desiring control of protein noise should target translation, not transcription. We have thus  
801 come to two conclusions about transcriptional noise: 1) mRNA noise directly inhibits  
802 growth by slowing translation and 2) very little mRNA noise propagates to the protein  
803 level. These two predictions cast new light on the functional consequences of  
804 transcriptional noise in living cells.

805

806 **Survival of the flattest.** We have shown that hyperbolic reactions, by filtering out high-  
807 amplitude input fluctuations, cause a reduction in both signal (average output) and noise  
808 (relative output variance) compared to first-order reactions. The strength of selection  
809 favoring organismal robustness is determined by the former, while the target of selection  
810 for increased robustness is determined by the latter. Both the strength of noise and its cost  
811 are mediated by the reaction order,  $a = K_m / (K_m + s)$  (Eqns. 6, 9), making this the obvious  
812 target for noise-attenuation. Noise-induced slowdown is minimal for first and zeroth  
813 order reactions ( $a = 1, 0$ , respectively), while noise itself is minimal only in zeroth order  
814 ( $a = 0$ ). Indeed, zeroth order kinetics completely filter out super-Poisson intrinsic input  
815 noise (Eqn. 9; Fig. 2), making the output entirely insensitive to input noise. We might call  
816 this phenomenon “zeroth order insensitivity”. Selection for robustness, then, should push

817 reactions towards their zeroth order regime. Geometrically, the zeroth order regime  
818 corresponds to the flat part of the rate curve, such that adaptive robustness by this  
819 mechanism promotes “survival of flattest”, a term previously invoked to describe  
820 selection for robustness via very different mechanisms (Wilke et al. 2001; Codoñer et al.  
821 2006).

822         There are two targets of modification for reaction order minimization: increasing  
823 substrate abundance,  $s$ , which moves the substrate distribution to the flatter part of the  
824 rate curve, or reducing  $K_m$ , which shifts the rate curve itself. Importantly, these strategies  
825 are not equivalent (Fig. 2). Decreasing  $K_m$  (while holding  $s$  constant) causes a unimodal  
826 change in product noise because  $K_m$  has two complementary effects: it reduces reaction  
827 order,  $a$ , which filters noise, but it also increases product burst size,  $b_T$  (Fig. 2).

828 Alternatively, increasing mRNA concentration (with  $K_m$  held constant) monotonically  
829 decreases protein noise because both reaction order,  $a$ , and translational bursting,  $b_T$ ,  
830 decline. This strategy increases fitness monotonically, thus avoiding the mRNA/protein  
831 noise trade-off. Importantly, increasing robustness by increasing substrate concentration  
832 is different from increasing robustness by decreasing substrate noise. To highlight this,  
833 one could force noise strength to stay constant while increasing mean substrate. The  
834 result would be increased robustness without a decrease in noise strength. This strategy  
835 undoubtedly generates other costs not accounted for in the current model, such as the  
836 energetic cost of increased mRNA synthesis. Future work should explicitly examine the  
837 evolutionary feasibility of cellular robustness via this mechanism in light of its potential  
838 trade-offs.

839

840 **Robustness via compartmentalization:** The present theory provides new insight into the  
841 recent argument that compartmentalization of reactions in eukaryotic cells is a strategy  
842 for noise control (Stoeger, Battich, and Pelkmans 2016). If cell compartments are filled  
843 with reactants via active transport by membrane transporters, then the resulting  
844 hyperbolic filtering (described above) will attenuate copy number fluctuations of  
845 reactants potentially down to the Poisson minimum, though not lower (e.g., Eqns 9 and  
846 S21). But hyperbolic filtering also slows the average rate of transport, creating  
847 inefficiency. In addition, reducing the reaction volume increases the relative size of  
848 fluctuations (the noise), which slows the absolute average rate of reactions within the  
849 compartment. Future work is necessary to determine exactly when the conflicting costs  
850 and benefits of hyperbolic filtering favor reaction compartmentalization as a strategy for  
851 cellular robustness. In eukaryotes, this takes on another dimension, because mRNA is  
852 actively transported out of the nucleus and into the cytoplasm before translation. The  
853 hyperbolic filtering of transcriptional noise at this level will greatly ameliorate  
854 transcriptional noise even before mRNA is filtered by translation. Future work should  
855 model the consequences of this serial transcriptional noise filtering in eukaryotes.

856

857 **Metabolic flux.** Wang and Zhang (Wang and Zhang 2011) used metabolic control  
858 analysis (MCA) to study the effect of enzyme fluctuations on metabolic flux, concluding  
859 that enzyme noise diminishes fitness in pathways with more than one enzyme. This result  
860 arises from the hyperbolic relationship between enzyme concentration and steady state  
861 flux in multi-enzyme metabolic pathways (Kacser and Burns 1973), and so, like the  
862 present paper, is a straightforward consequence of nonlinear averaging. However, MCA

863 models assume first-order enzyme kinetics (Kacser and Burns 1973). In light of the  
864 present results, it is not immediately clear if this assumption makes the results of (Wang  
865 and Zhang 2011) an underestimate of the true cost of gene expression noise, because  
866 hyperbolic rate laws cause noise-induced slowdown at each pathway step, or an  
867 overestimate of cost, because hyperbolic rate laws attenuate noise at each step, thus  
868 damping noise propagation through the pathway. Future work should distinguish between  
869 these two possibilities.

870

871 **The crucial assumption:** All of the major results of this paper rely on the validity of  
872 hyperbolic reaction kinetics as an accurate description of *in vivo* biochemical reactions.  
873 This assumption is not guaranteed, especially for translation. Translation is a unique  
874 biochemical reaction because a single substrate molecule (mRNA) is simultaneously  
875 bound to multiple enzymes (ribosomes), and it is therefore easy to imagine that  
876 translation kinetics may not be well-described by Michaelis-Menten reaction kinetics.  
877 Therefore, the present theoretical results must remain in question until future empirical  
878 work carefully measures the *in vivo* rate law of protein translation.

879

## 880 **METHODS**

881 All Monte Carlo simulations were run on the elementary (microscopic) system of  
882 reactions (S17) using Gillespie's Exact Stochastic Simulation Algorithm (SSA) (Gillespie  
883 1977) implemented in *Mathematica* v10.3.0.0 with the open source xSSALite package,  
884 freely available from the xCellerator project (<http://www.xlr8r.info/SSA/>). Each  
885 realization of the simulation is a Markov Chain with irregular (exponentially distributed)  
886 time intervals. For statistical analysis, each realization was regularized using the



887 TimeSeriesAggregate[] command. Mean[] and Variance[] commands of the regularized  
888 time series were used to compute the moments of the Markov Chain. Data points in  
889 figures 1,2 and 3 are ensemble averages or variances over 1000-10,000 simulated  
890 realizations of the stochastic process; each realization was run for at least 10 times the  
891 half life of S or P (half-life =  $\ln(2)/\delta$ ), whichever was longer, and the value for that  
892 realization was  $s$  or  $p$  during the final second of the simulation run.

893

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